

§1. Computation of stable stems

Goal: $\pi_{\leq 14, *}(v_{\mathbb{F}_2} S_2^{\wedge})$. Abbreviate by $\pi_{*,*}(S)$

Input: $\pi_{\leq 15, *}(C\tau)$. This is algebraic data: Ext groups of \mathcal{A}_* -comodules.

Prop: $\pi_{0,*}(S) \cong \mathbb{Z}_2\langle \tau \rangle / (\tau^2 = 2)$, and $\tilde{2} \mapsto h_0$ under $S \rightarrow C\tau$

Proof: First check that there are no differentials into zero-stem.

Suppose $d_r(h_1) = h_0^{r+1}$. $\xrightarrow{\text{Leibniz}}$ $d_r(h_0 h_1) = h_0^{r+2}$ \Downarrow
 $d_r(0) = 0$

Leibniz rule is: $\partial : C\tau \rightarrow \Sigma^{1,-2} C\tau$ is a derivation

By omnibus: $\pi_{0,*}(S)$ is τ -power torsion-free.

$\Rightarrow \pi_{0,*} S \hookrightarrow \pi_{0,*} S[\tau^{-1}] \cong \mathbb{Z}_2[\tau^{\pm 1}]$.

is precisely $\xrightarrow{\text{mmap}}$ the quotient by the τ -multiples \otimes

\downarrow
 $\pi_{0,*} C\tau \cong \mathbb{F}_2[h_0]$

Can now determine $\pi_{0,*} S$ as a $\mathbb{Z}[\tau]$ -submodule of $\mathbb{Z}_2[\tau^{\pm 1}]$.

We know $\pi_{0, \leq 0}(S) \cong \mathbb{Z}_2\langle \tau \rangle$. Only need to determine positive filtrations

\otimes implies:

something in $\pi_{0,*} S$ is τ -div. \Leftrightarrow maps to zero in $C\tau$.

Now we use: $\pi_{0,*} C\tau \cong \mathbb{F}_2[h_0]$ ($|h_0| = (0,1)$)

$1 \mapsto 1 \neq 0, \quad 2 \mapsto 2 = 0, \quad \text{so } \boxed{\frac{2}{\tau} \in \pi_{0,*} S.} \leftarrow =: \tilde{2}$

Likewise, $\tilde{4} := \frac{4}{\tau^2}, \quad \tilde{6} := \frac{6}{\tau}, \dots$ in general, $\tilde{n} = \frac{n}{\tau^{\nu_2(n)}}$.

Then must have $\tilde{2} \mapsto h_0$ by surjectivity.

This determines everything: $\pi_{0,*} \cong \mathbb{Z}_2\langle \tau, \tilde{2} \rangle / (\tau^2 = 2)$. \square

Observe: All $\pi_{n,*}$ are modules over $\pi_{0,*} S$.

Prop: $\pi_{1,*} \cong \mathbb{Z}/2[\tau, \tilde{\tau}] \langle \eta \rangle / (\tilde{\tau}\eta = 0)$. $|\eta| = (1,1)$, $\eta \mapsto h_1$.

Proof: Again τ -torsion free.

h_1 supports no differentials $\xrightarrow{\text{omnibus}}$ lifts to $\pi_{1,1} \mathbb{S}$.

Lift is unique: no elements of higher filtration. Call this η .

Remaining: $\tilde{\tau} \cdot \eta = ?$

$\tilde{\tau} \cdot \eta \in \pi_{1,2}$ but $\pi_{1,2} = 0$ because $E_2^{1,5} = 0$ for $s \geq 2$.

Lastly, $2 \cdot \eta = \tau \cdot \tilde{\tau} \eta = 0$, so this is indeed a $\mathbb{Z}/2$ -module. \square

Remark: Write η instead of $\tilde{\eta}$: we are not reconstructing $\pi_{*,*} \mathbb{S}$ from $\tau \mathbb{S}$, but building $\pi_{*,*} \mathbb{S}$ from the ground up. (But $\tilde{\tau}$ is forced on us: there we can only think about it as lifting the classical element.)

Prop: $\pi_{3,*} \cong \mathbb{Z}/8[\tau, \tilde{\tau}] \langle v \rangle / (\tilde{\tau}^3 \cdot v = \eta^3, \tilde{\tau}^3 \cdot v = 0)$.

Proof: τ -torsion-free, and h_2 lifts to $\pi_{3,1}$.

Choose a lift v of h_2 . Or, you can lift your preferred $v \in \pi_3 \mathbb{S}$.

$\tilde{\tau}^3 \cdot v - \eta^3 \mapsto h_0^3 h_2 - h_1^3 \mapsto 0$ under $\mathbb{S} \rightarrow \mathbb{C}\tau$.

$$\begin{array}{ccccc} \pi_{3,4} \mathbb{S} & \xrightarrow{\tau} & \pi_{3,3} \mathbb{S} & \longrightarrow & \pi_{3,3} \mathbb{C}\tau \\ \times & \longmapsto & \tilde{\tau}^3 \cdot v - \eta^3 & \longmapsto & 0 \end{array}$$

But $\pi_{3,4} \mathbb{S} = 0$.

As before, $8v = \tau^3 \cdot \tilde{\tau}^3 \cdot v = 0$. \square

One continues in the same way, until we get to stem 14, where differentials could (and do indeed) appear.

Prop: $\pi_{7,*} \cong \mathbb{Z}/16[\tau, \tilde{\tau}] \langle \sigma \rangle / (\tilde{\tau}^4 \cdot \sigma = 0)$ $|\sigma| = (7,1)$.

Then $\sigma \cdot \sigma = -\sigma \cdot \sigma$, so $2\sigma^2 = 0$.

$$\rightarrow 0 = 2\sigma^2 = \tau \cdot \tilde{\tau} \sigma^2.$$

Omnibus: $\tilde{\tau} \sigma^2$ is hit by a diff of length at most 2^n .
 $h_0 h_3$.

Only one option:

Prop: $d_2(h_4) = h_0 h_3^2$ first Hopf invariant 1 differential.

Rmk: While this time we didn't need it, it's a very general & powerful method to have equations put bounds on differentials. Also, we got more than the differential: the specific lift $\tilde{\sigma}^2$ is τ -torsion. (In general, only have: $d_r \Rightarrow \exists \tau^{r-1}$ -torsion lift)

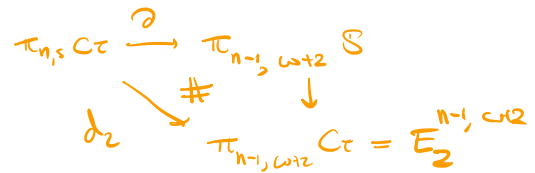
Next goal: $d_3(h_0 h_4) = h_0 d_0$.

Cannot follow Rom Leibniz rule ∇
We'll use a beefed-up version.

Recall: ∂ is boundary map $S^{s-1} \xrightarrow{\tau} S \rightarrow C\tau \xrightarrow{\partial} \Sigma^{s-2} S$.

Prop: $\partial(h_4) = \tilde{\sigma}^2$.

Proof: Know that $\tau \cdot \tilde{\sigma}^2 = 0$. LES:



$$\mathbb{Z}/2\langle h_4 \rangle \cong \pi_{15,1}(C\tau) \rightarrow \pi_{14,3} S \xrightarrow{\tau} \pi_{14,2} S \quad \square$$

only option \rightarrow $\tilde{\sigma}^2 \mapsto 0$

Prop: $d_3(h_0 h_4) = h_0 d_0$.

Proof:

$$\begin{aligned} \partial(h_0 h_4) &= \partial(\tilde{\sigma} \cdot h_4) = \tilde{\sigma} \cdot \partial(h_4) \\ &= \tilde{\sigma} \cdot \tilde{\sigma} \cdot \sigma^2 \stackrel{!}{=} \tau \cdot \tilde{\sigma} \kappa. \end{aligned}$$

In part, $\tau^2 \tilde{\sigma} \kappa = 0$

This implies $d_3(h_0 h_4) = \overline{\tilde{\sigma} \kappa} = h_0 d_0$. \square

$\tilde{\sigma}^2$ contributes to $\pi_{14,3}$ not to $\pi_{14,2}$

Prop: $\pi_{14,*} = \mathbb{Z}/8\langle \tilde{\sigma}, \tau \rangle \langle \kappa, \sigma^2 \rangle / (\tilde{\sigma}^2 = \tau \tilde{\sigma} \kappa, \tau \cdot (\tilde{\sigma}^2) = 0, \tilde{\sigma} \kappa = 0)$.

The relation $\tilde{\sigma}^2 = \tau \cdot \tilde{\sigma} \kappa$ has one τ in it, causing a differential stretch by one

In a precise way, these "hidden extensions" are equivalent to differential stretching.

Eg. Suppose we know $\partial(h_4) = \tilde{\sigma}^2$ & $\partial(h_0 h_4) = \tau \tilde{\sigma} \kappa$.

Observe: $h_0 h_4 = \tilde{\sigma} \cdot h_4$. So

$$\tilde{\sigma} \cdot (\tilde{\sigma}^2) = \tilde{\sigma} \partial(h_4) = \partial(\tilde{\sigma} h_4) = \partial(h_0 h_4) = \tau \tilde{\sigma} \kappa. \quad \square$$

§2. Keeping track of filtration

In S_p , it is not always clear what a construction will do to the Adams filtration of a map.
 In $S_{\mathbb{Z}_p}$, the Ad \mathbb{Q} manifests geometrically, which makes studying this easier.

$$f, g, h \in \pi_* S \text{ with } fg=0 \text{ \& } gh=0 \rightarrow \langle f, g, h \rangle \subseteq \pi_* S.$$

in degree $|f| + |g| + |h| - 1$.
└
due to suspension

But there's indeterminacy, and we have no control over Adams filtration.

Ex. $\langle \overset{1}{\eta}, \overset{0}{2}, \overset{6}{\nu^2} \rangle = \{ \varepsilon, \varepsilon + \tau\sigma \} \subseteq \pi_8 S,$

If we have synthetic lifts $\tilde{f}, \tilde{g}, \tilde{h}$ and $\tilde{f}\tilde{g}=0$ & $\tilde{g}\tilde{h}=0$ still, then get

$$\langle \tilde{f}, \tilde{g}, \tilde{h} \rangle \subseteq \pi_{**} S \text{ in stem } |f| + |g| + |h| - 1 \text{ and filtr. } |f|_S + |g|_S + |h|_S + 1.$$

Ex. $\langle \underset{(1,1)}{\eta}, \underset{(0,1)}{\tilde{2}}, \underset{(6,2)}{\nu^2} \rangle = \{ \varepsilon \} \subseteq \pi_{8,3} S.$

So one can now define ε as $\langle \eta, \tilde{2}, \nu \rangle$.

Sometimes, we need to insert τ 's to make synthetic brackets be zero.

Ex. $\langle \sigma^2, \underset{\uparrow}{2}, \nu \rangle \subseteq \pi_{16,2} S.$ This detects $[h_1, h_4]$.
 $= \tau \tilde{2}$
 $\tilde{2}\sigma^2 \neq 0$ while $2\sigma^2 = 0$.

Have a procedure for evaluating Toda brackets with τ 's in them: Moss' theorem.

"Keeping track of filtration" (also in more involved ways) is at the heart of Burklund - Hahn - Singer.

§4. Importing algebraic differentials.

Warning: anachronistic perspective!

So far, we only looked internally at the sseq coming from a synthetic spectrum, and then used beefed up methods to compute this. Now, we'll take a different approach: use Syn_{MU} to create a zigzag

$$ANSS(\mathbb{S}) \leftarrow \{E_r^{x,*}\} \rightarrow \text{"algebraic Novikov sseq"}$$

at which point you can import differentials.

\Rightarrow Isaksen - Wang - Xu stem go.

Important: previous things (total diff, S/\mathbb{S} , etc) could, with some extra work, have been done with filtered spectra instead. This application is truly synthetic.

Idea: You can run the \mathbb{F}_p -Adams sseq in MU-synthetic spectra.

This will result in an E_2 -term with an additional grading.

$$E_2^{h,w,s} = \text{Ext}_{A_{*,*}}^{S, h+s, w-s}(\mathbb{F}_p[\tau], \mathbb{F}_p[\tau]) \Rightarrow \pi_{h,w} S_p^\wedge$$

The w is new: it measures Adams - Novikov filtration.

ms Another case of "synthetic spectra help you keep track of the Adams filtration."

Observe: when invert τ , we obtain ordinary \mathbb{F}_p -ASS. So τ^{-1} gives a map

$$v\mathbb{F}_p\text{-ASS}(v\mathbb{S}) \longrightarrow \mathbb{F}_p\text{-ASS}.$$

Can compute E_2 -page of sASS: looks similar, but with τ 's adjoined. The additional grading is new information however.

Isaksen: used the additional grading as a restraint on differentials.

Diff'1 has to τ -invert to a classical one, and all classical differentials must arise in some way. But sometimes, the third grading excludes differentials for degree reasons.

Deeper is to then also consider $C\tau \otimes -$. Then we get

$$(v\mathbb{F}_p/\tau)\text{-ASS}(C\tau) \text{ in } \text{Stable}_{MU, \mu} \cong \text{algebraic Novikov sseq.}$$

Have two maps: $v\mathbb{S} \rightarrow C\tau$ and $C\tau \rightarrow S^{1,-2}$. Can use both to obtain diff's in $v\mathbb{S}$. Then τ -invert.

