

## 22a. Case $0 < m_0(j_1) < p$ , and $\varepsilon = 1$

Now  $r_0$  is active in restricting the components. We consider the two lowest components of F.

```
In[1]:= Clear[f, fp]
F = tht[0] f[r0, t] Phi[h, p, r0, p] + tht[1] f[r0 + 2, t] Phi[h, p, r0 + 2, p]
Out[1]= f[r0, t] Phi[h, p, r0, p] tht[0] + f[2 + r0, t] Phi[h, p, 2 + r0, p] tht[1]
```

We use that one shift operator gives zero :

```
In[2]:= sh[3, -1, F, subnab] /. eps → 1 // Simplify(* should be zero *)
Coefficient[%, Phi[3 + h, -1 + p, 1 + r0, -1 + p]]
fr0p2 = f[2 + r0, t] /. Solve[% == 0, f[2 + r0, t]][1] // Simplify
Out[2]= -1/(4 (1 + p)) p ((4 - h + 2 p + r0 - 4 ell π t^2) f[r0, t] Phi[3 + h, -1 + p, 1 + r0, -1 + p] tht[0] +
4 i √(2 π) t √Abs[ell] f[2 + r0, t] Phi[3 + h, -1 + p, 1 + r0, -1 + p] tht[0] +
6 f[2 + r0, t] Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] -
h f[2 + r0, t] Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] +
2 p f[2 + r0, t] Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] +
r0 f[2 + r0, t] Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] -
4 ell π t^2 f[2 + r0, t] Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] -
2 t Phi[3 + h, -1 + p, 1 + r0, -1 + p] tht[0] f^(0,1)[r0, t] -
2 t Phi[3 + h, -1 + p, 3 + r0, -1 + p] tht[1] f^(0,1)[2 + r0, t])

Out[3]= -1/(4 (1 + p)) p ((4 - h + 2 p + r0 - 4 ell π t^2) f[r0, t] tht[0] +
4 i √(2 π) t √Abs[ell] f[2 + r0, t] tht[0] - 2 t tht[0] f^(0,1)[r0, t])

Out[4]= -(i ((-4 + h - 2 p - r0 + 4 ell π t^2) f[r0, t] + 2 t f^(0,1)[r0, t])) / (4 √(2 π) t √Abs[ell])
Out[5]= -(i ((-4 + h - 2 p - r0 + 4 ell π t^2) f[r0, t] + 2 t f^(0,1)[r0, t])) / (4 √(2 π) t √Abs[ell])
```

We get  $F^+$  by the other downward shift operator.

```
In[ 0]:= Fp = sh[-3, -1, F, subnab] /. eps → 1 // Simplify
fp[r0 - 1, t] = Coefficient[Fp, Phi[-3 + h, -1 + p, -1 + r0, -1 + p]] / tht[0] // Simplify
Out[ 0]= - $\frac{1}{4 \times (1 + p)}$ 
p (f[r0, t] ((4 + h + 2 p - r0 + 4 ell π t2) Phi[-3 + h, -1 + p, -1 + r0, -1 + p] × tht[0] - 4 i √2 π
t √Abs[ell] Phi[-3 + h, -1 + p, 1 + r0, -1 + p] × tht[1]) +
f[2 + r0, t] ((2 + h + 2 p - r0 + 4 ell π t2) Phi[-3 + h, -1 + p, 1 + r0, -1 + p] × tht[1] -
8 i √π t √Abs[ell] Phi[-3 + h, -1 + p, 3 + r0, -1 + p] × tht[2]) -
2 t (Phi[-3 + h, -1 + p, -1 + r0, -1 + p] × tht[0] f(0,1)[r0, t] +
Phi[-3 + h, -1 + p, 1 + r0, -1 + p] × tht[1] f(0,1)[2 + r0, t]))
```

$$\frac{-p(4 + h + 2 p - r0 + 4 \text{ell} \pi t^2) f[r0, t] + 2 p t f^{(0,1)}[r0, t]}{4 \times (1 + p)}$$

We know the lowest order term of  $b_u^+$ .

Hence we have with a non-zero factor **cnz**:

```
In[ 0]:= Clear[cnz]
eq = (- (4 + h + 2 p - r0 + 4 ell π t2) f[r0, t] + 2 t f(0,1)[r0, t]) ==
cnz t^(0 + 1) WhittakerV[-p + 1 - (j1 + 1)/2, nu1, 2 Pi Abs[ell] t^2] /.
Abs[ell] → ell /. Abs[h + p] → h + p // Simplify
(* Use that V is even in the second parameter. *)
Out[ 0]= 2 t f(0,1)[r0, t] ==
```

$$(4 + h + 2 p - r0 + 4 \text{ell} \pi t^2) f[r0, t] + cnz t \text{WhittakerV}\left[\frac{1}{4} \times (2 - h - p), \frac{h + p}{2}, 2 \text{ell} \pi t^2\right]$$

We solve for the derivative of  $f_{r_0}$

```
In[ 0]:= fr0d = f(0,1)[r0, t] /. Solve[eq, f(0,1)[r0, t]][[1]] // Simplify
Out[ 0]=  $\frac{1}{2 t} \left( (4 + h + 2 p - r0 + 4 \text{ell} \pi t^2) f[r0, t] + cnz t \text{WhittakerV}\left[\frac{1}{4} \times (2 - h - p), \frac{h + p}{2}, 2 \text{ell} \pi t^2\right] \right)$ 
```

This can be inserted in the eigenfunction equations for  $r = r_0$ , together with the expressions for  $f_{r_0+2}$

```
In[ = e i = efeqn[h, p, r0, f, ell, 0, 1] /. j → j1 /. nu → nu1 /. f[r0 + 2, t] → fr0p2 /.
      f^(0,1)[2 + r0, t] → D[fr0p2, t] /. f^(0,2)[r0, t] → D[fr0d, t] /.
      f^(0,1)[r0, t] → fr0d // Whrel /. hpsub /. Abs[h + p]^2 → (h + p)^2 // Simplify

Out[ = {1/16 cnz t (4 (-4 + 2 h + p + r0 + 8 ell π t^2) WhittakerV[(1/4 (2 - h - p), (h + p)/2, 2 ell π t^2)] -
      (-16 + 3 h^2 + 8 p + 3 p^2 + h (8 + 6 p)) WhittakerV[(1/4 (6 - h - p), (h + p)/2, 2 ell π t^2)]),
      3/16 cnz (p - r0) t (-4 (-4 + 2 h + p + r0 + 8 ell π t^2) WhittakerV[(1/4 (2 - h - p), (h + p)/2, 2 ell π t^2)] +
      (-16 + 3 h^2 + 8 p + 3 p^2 + h (8 + 6 p)) WhittakerV[(1/4 (6 - h - p), (h + p)/2, 2 ell π t^2)])}
```

These expressions should be zero. We insert the asymptotic behavior of the V-Whittaker function.

```
In[ = Clear[tau]
eia = ei /. t → Sqrt[tau]/Sqrt[2 Pi Abs[ell]] /. Abs[ell] → ell /.
      WhittakerV[kap_, s_, tau_] → -E^(-Pi I kap) tau^(-kap) E^(tau/2) // Simplify

Out[ = {1/(16 Sqrt[ell] Sqrt[2 π]) cnz e^(tau/2) tau^(1/4 (-4+h+p)) (e^(1/4 i (-6+h+p) π) (-16 + 3 h^2 + 8 p + 3 p^2 + h (8 + 6 p)) - 4 e^(1/4 i (-2+h+p) π) tau (-4 + 2 h + p + r0 + 4 tau)),
      1/(16 Sqrt[ell] Sqrt[2 π]) 3 cnz e^(tau/2) (p - r0) tau^(1/4 (-4+h+p)) (-e^(1/4 i (-6+h+p) π) (-16 + 3 h^2 + 8 p + 3 p^2 + h (8 + 6 p)) + 4 e^(1/4 i (-2+h+p) π) tau (-4 + 2 h + p + r0 + 4 tau))}
```

```
In[ = eia E^(-tau/2) tau^(-1 - (h + p)/4) // Simplify
Limit[% , tau → Infinity]

Out[ = {((i cnz e^(1/4 i (h+p) π) (3 h^2 + 3 p^2 + 4 p (2 + tau) + h (8 + 6 p + 8 tau) + 4 (-4 + (-4 + r0) tau + 4 tau^2))) / (16 Sqrt[ell] Sqrt[2 π] tau^2), -(3 i cnz e^(1/4 i (h+p) π) (p - r0) (3 h^2 + 3 p^2 + 4 p (2 + tau) + h (8 + 6 p + 8 tau) + 4 (-4 + (-4 + r0) tau + 4 tau^2))) / (16 Sqrt[ell] Sqrt[2 π] tau^2)} / ((16 Sqrt[ell] Sqrt[2 π] tau^2))

Out[ = {i cnz e^(1/4 i (h+p) π) / (16 Sqrt[ell] Sqrt[2 π]), -48 i cnz e^(1/4 i (h+p) π) p + 48 i cnz e^(1/4 i (h+p) π) r0 / (16 Sqrt[ell] Sqrt[2 π])}
```

So **cnz** has to vanish.