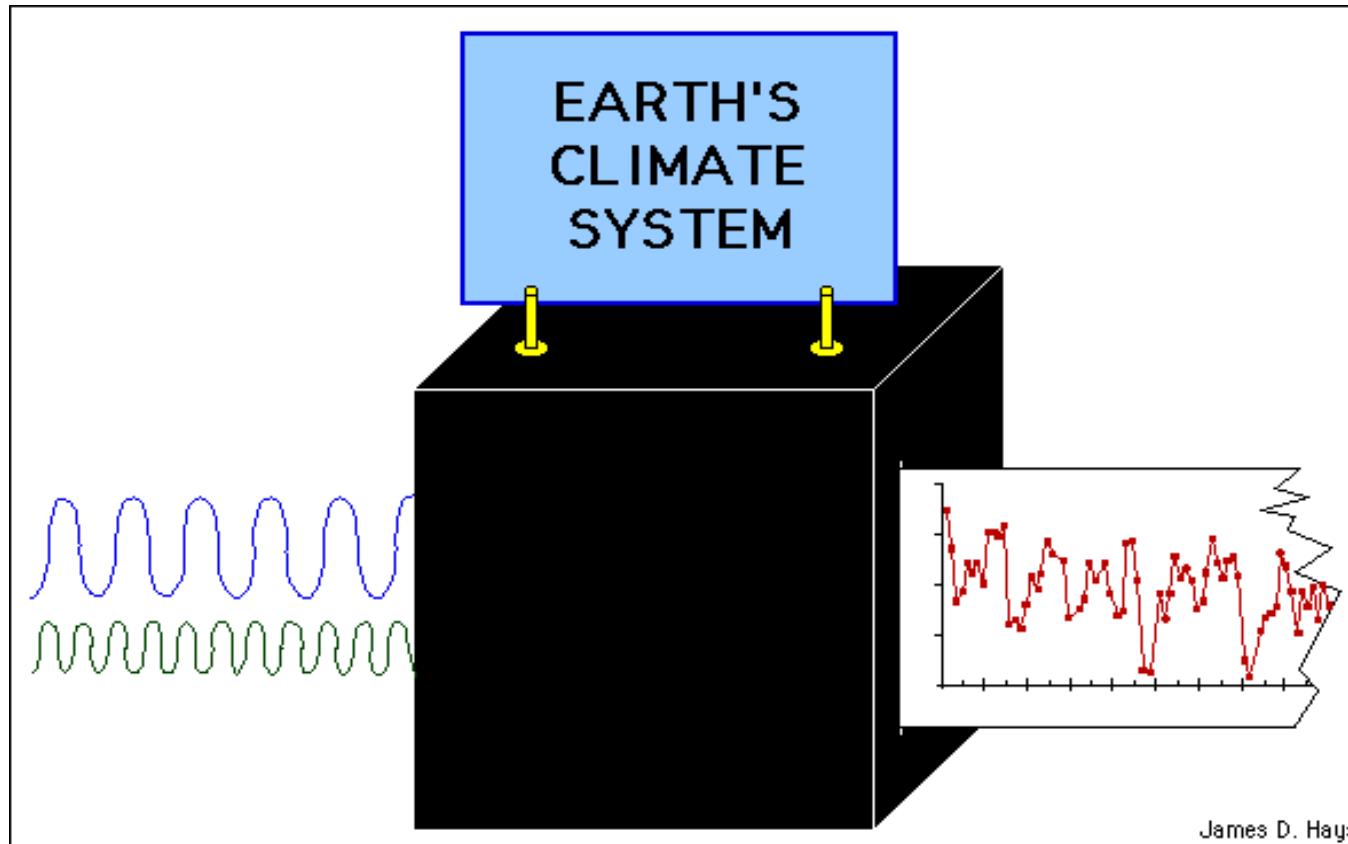
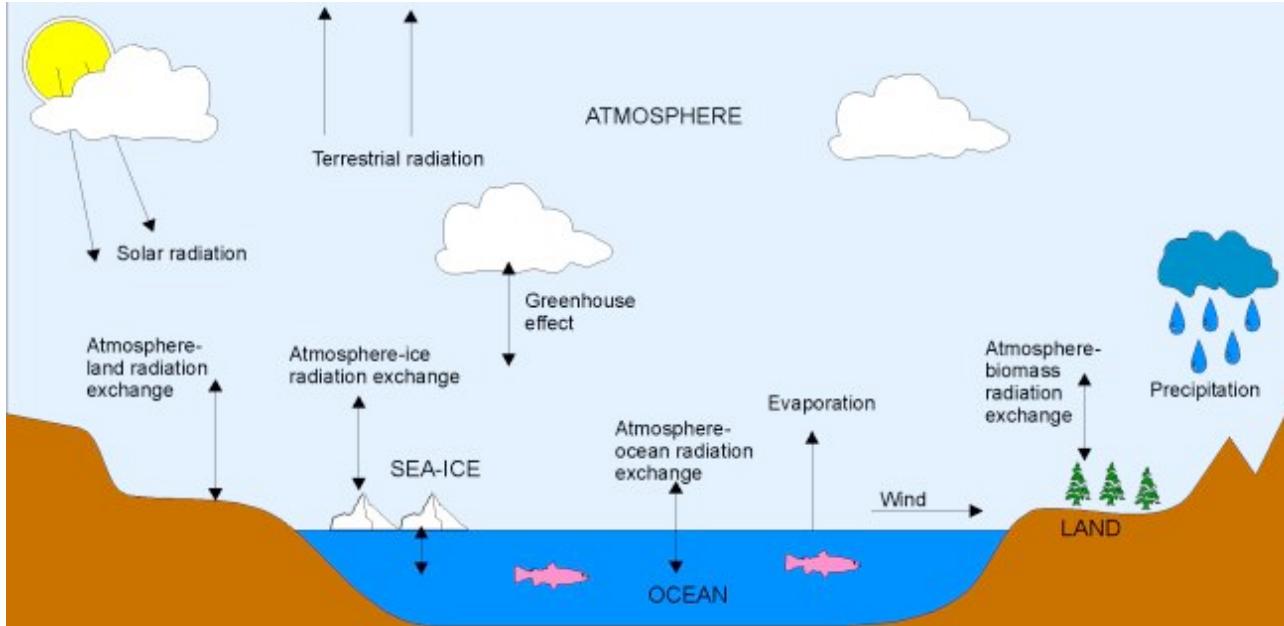


Stochastic Climate Dynamics



Henk Dijkstra, IMAU & CCSS
Physics Department, Utrecht University, Utrecht, Netherlands

Climate System



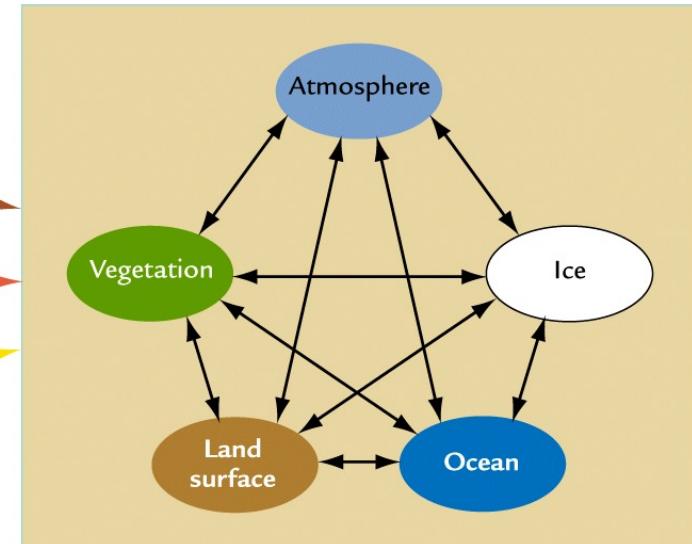
CAUSES
(external forcing)

Changes in plate tectonics

Changes in Earth's orbit

Changes in Sun's strength

CLIMATE SYSTEM
(internal interactions)



CLIMATE VARIATIONS
(internal responses)

Changes in Atmosphere

Changes in Ice

Changes in vegetation

Changes in Ocean

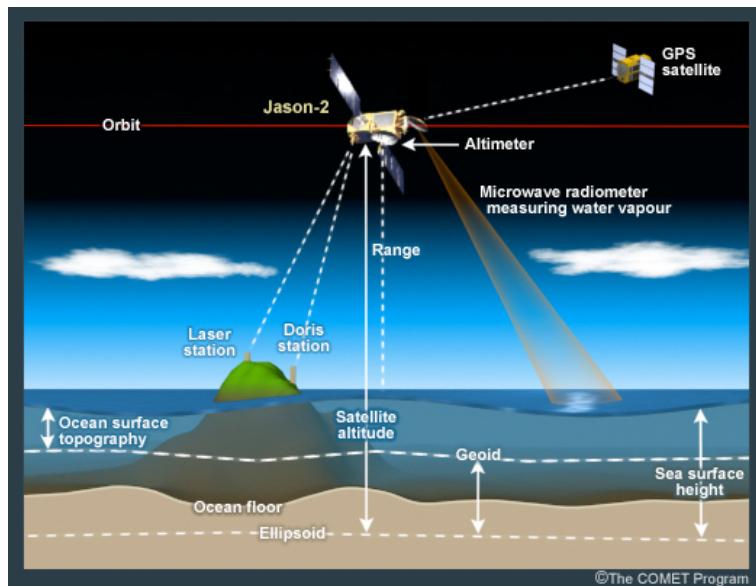
Changes in land surface

Observations

Instrumental data
(~ 1880 -)



(1992 -)

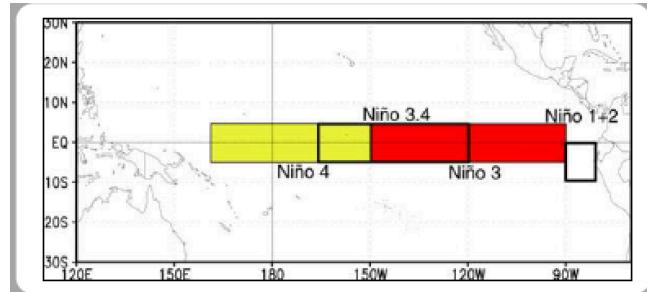


Proxy data
(geological past)



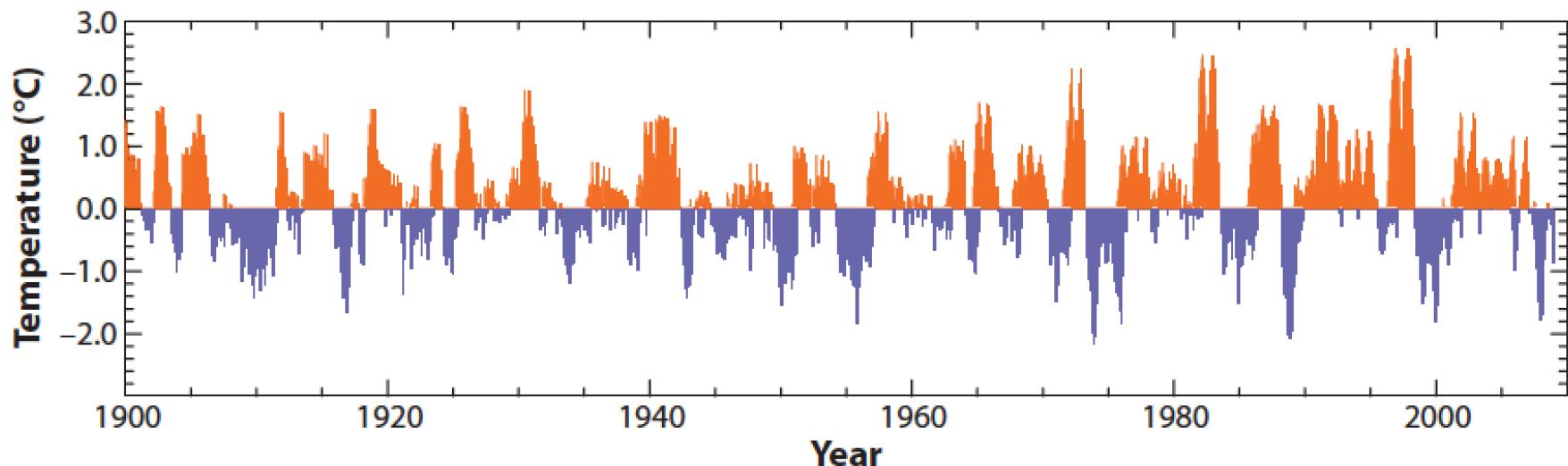
Very limited for hypotheses/theory falsification

Example instrumental data: El Niño/Southern Oscillation (ENSO)

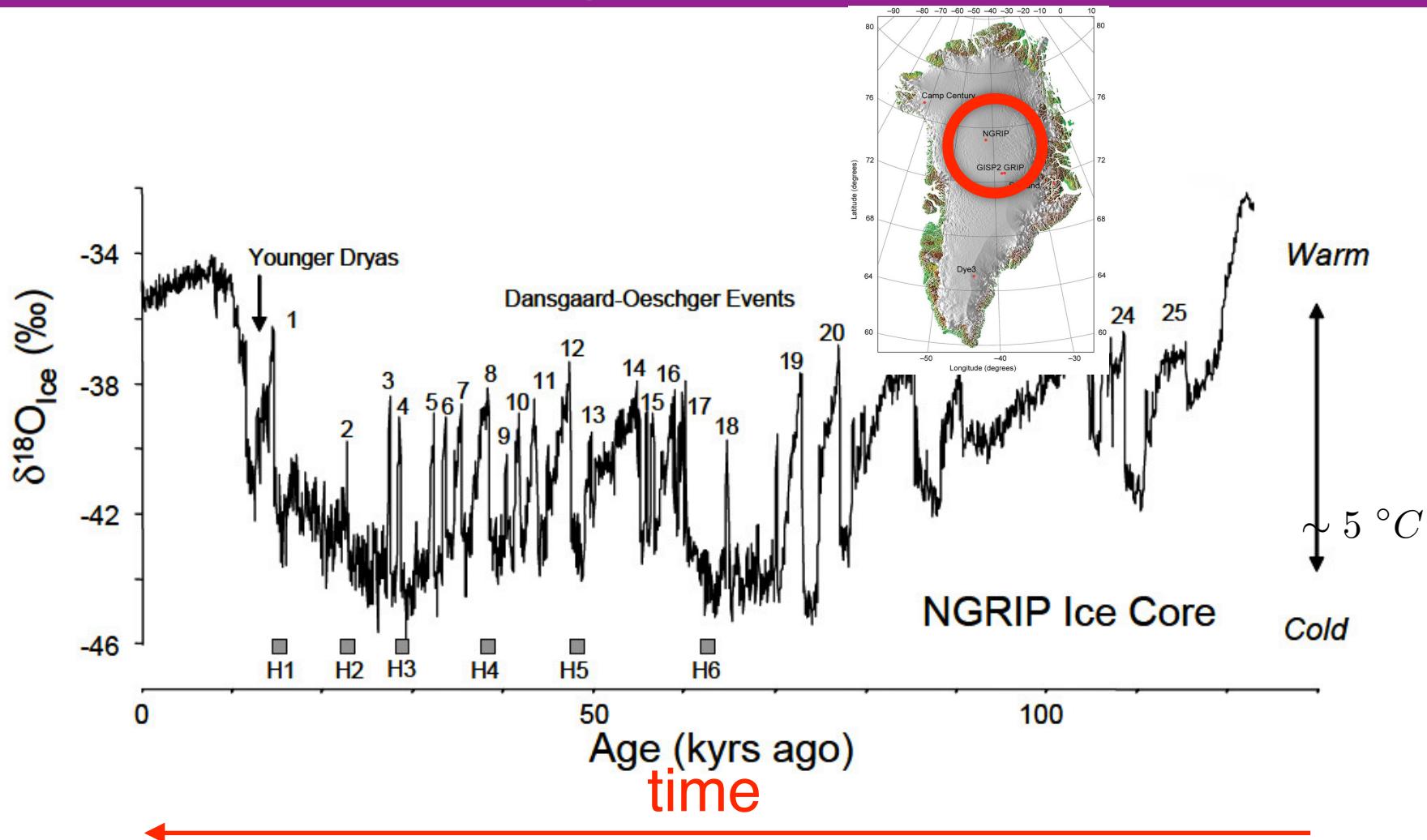


NINO3.4 index

sea surface temperature
anomaly equatorial Pacific



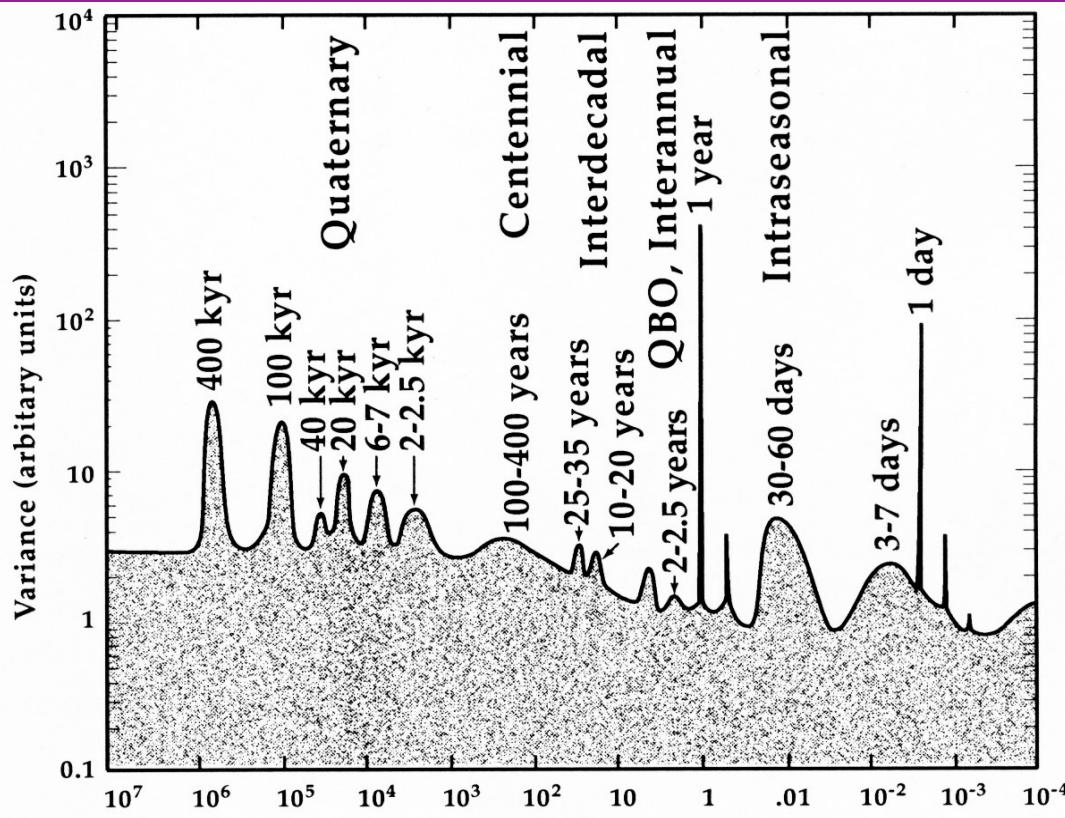
Example proxy data: Ice core oxygen isotope record



<https://www.ncdc.noaa.gov/data-access/paleoclimatology-data/datasets>

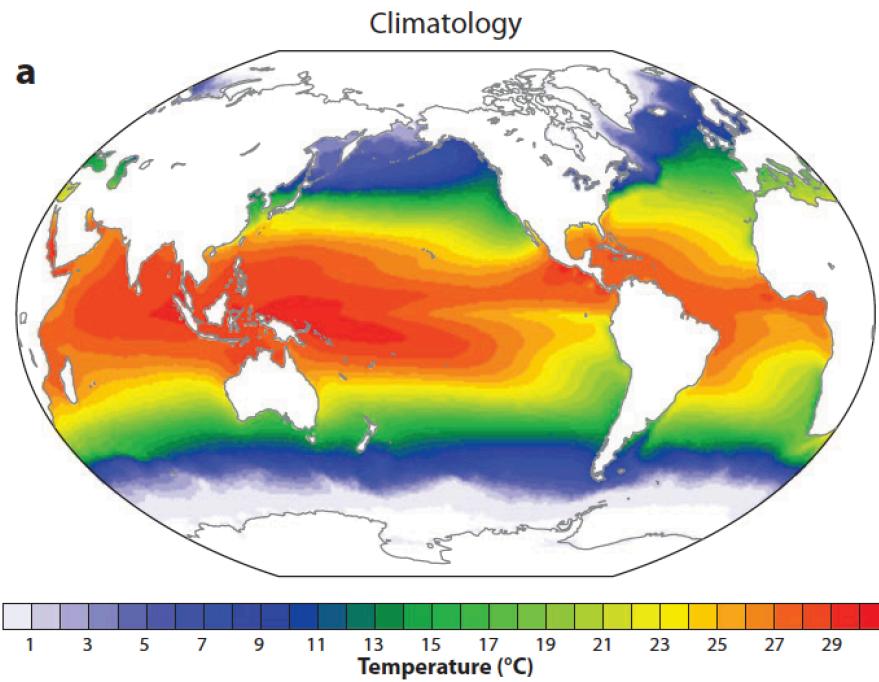
Traditional (artist's) view of Variability in the Climate System

‘Energy’

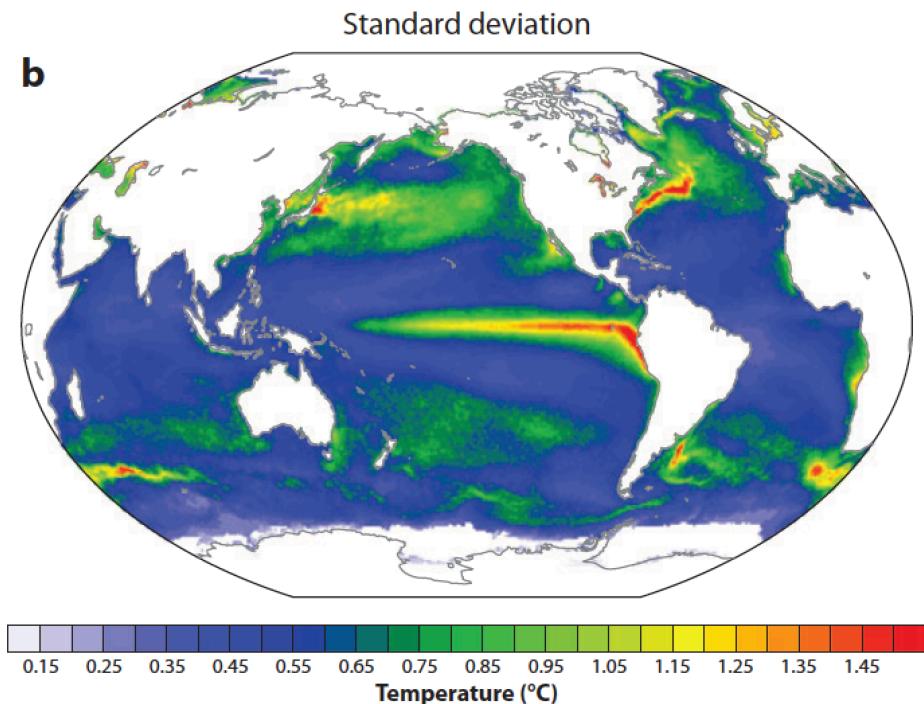


1. Sharp peaks
2. Continuous background
3. Elevated energy in frequency bands

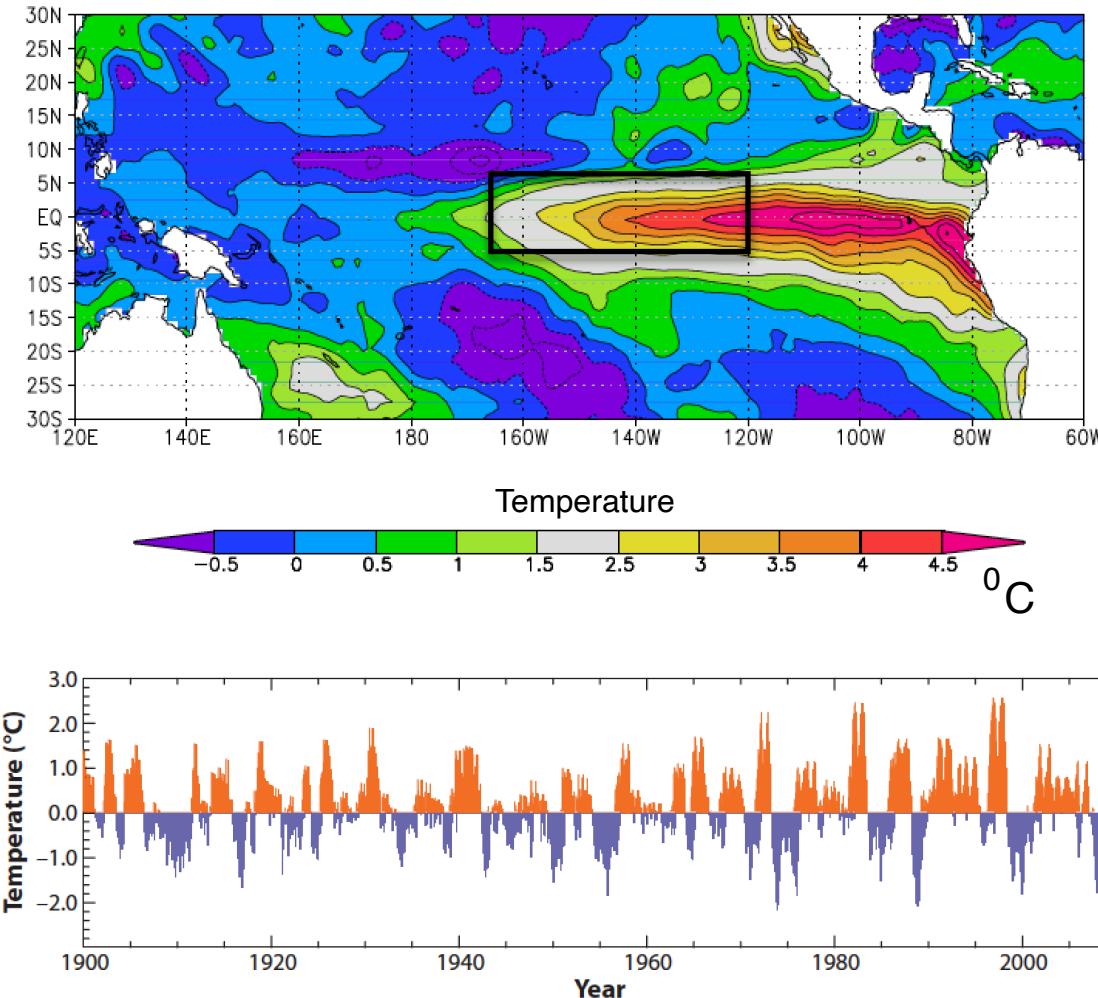
Sea surface temperature (SST) variability



1982-2008

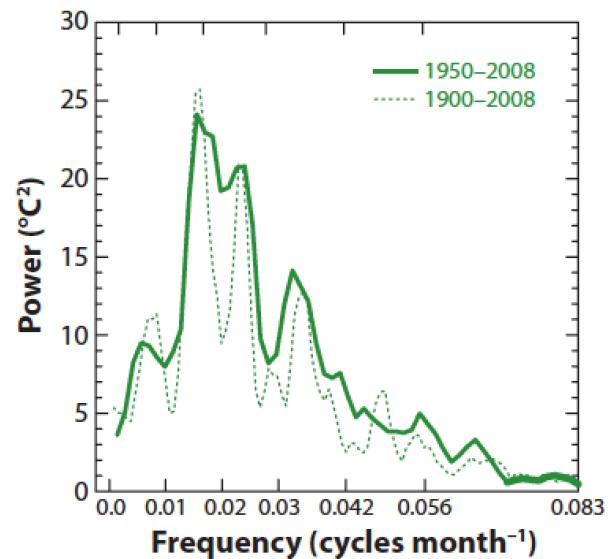


Patterns of variability: El Nino/Southern Oscillation (ENSO)



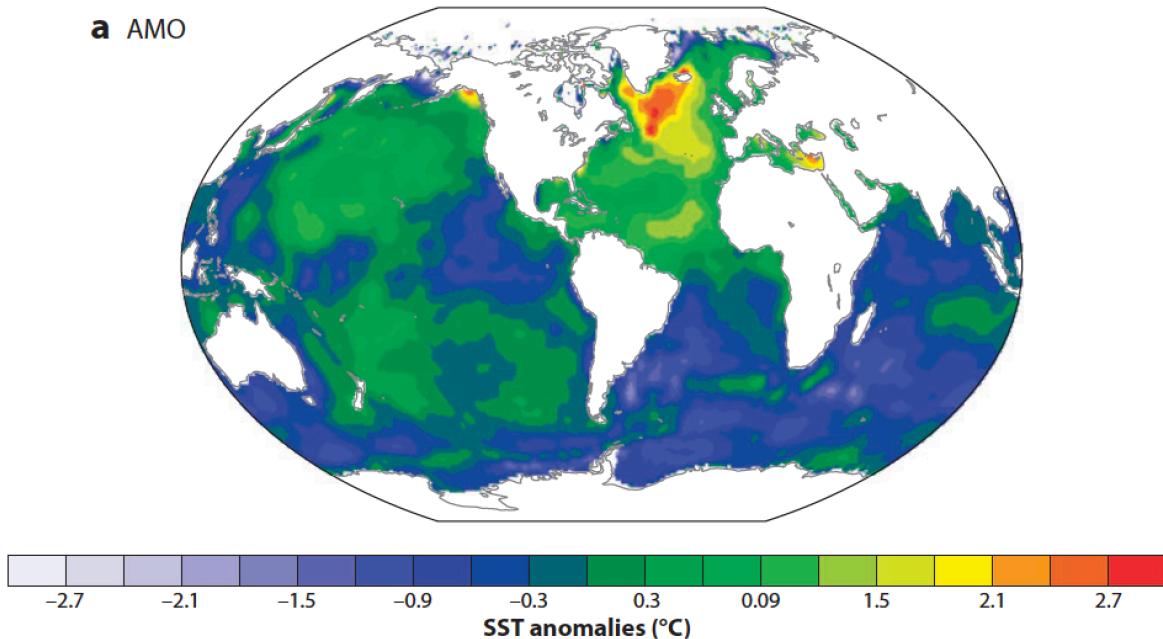
NINO3.4 index

SST anomaly
(1982-2010 mean)
of December 1997



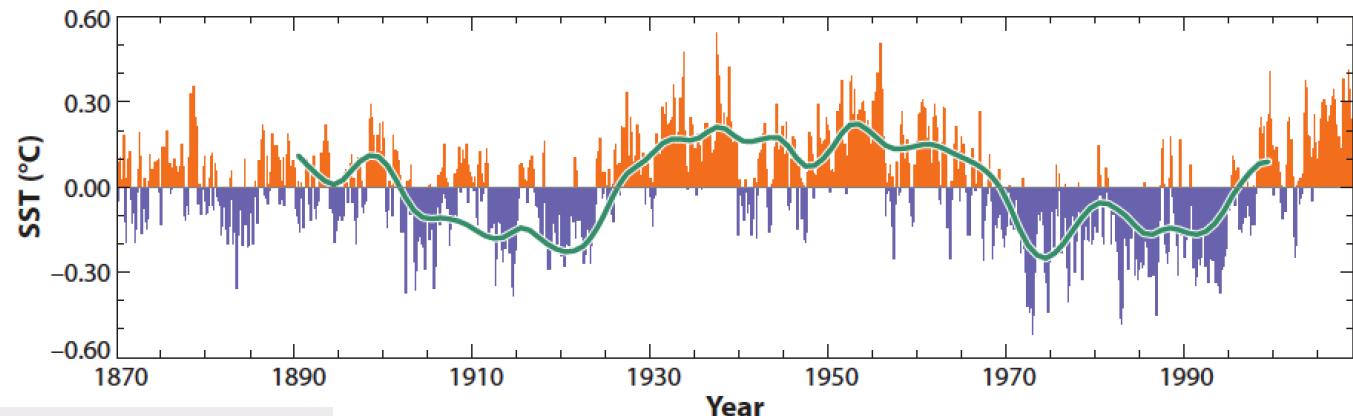
Patterns of variability: Atlantic Multidecadal Variability (AMV)

a AMO



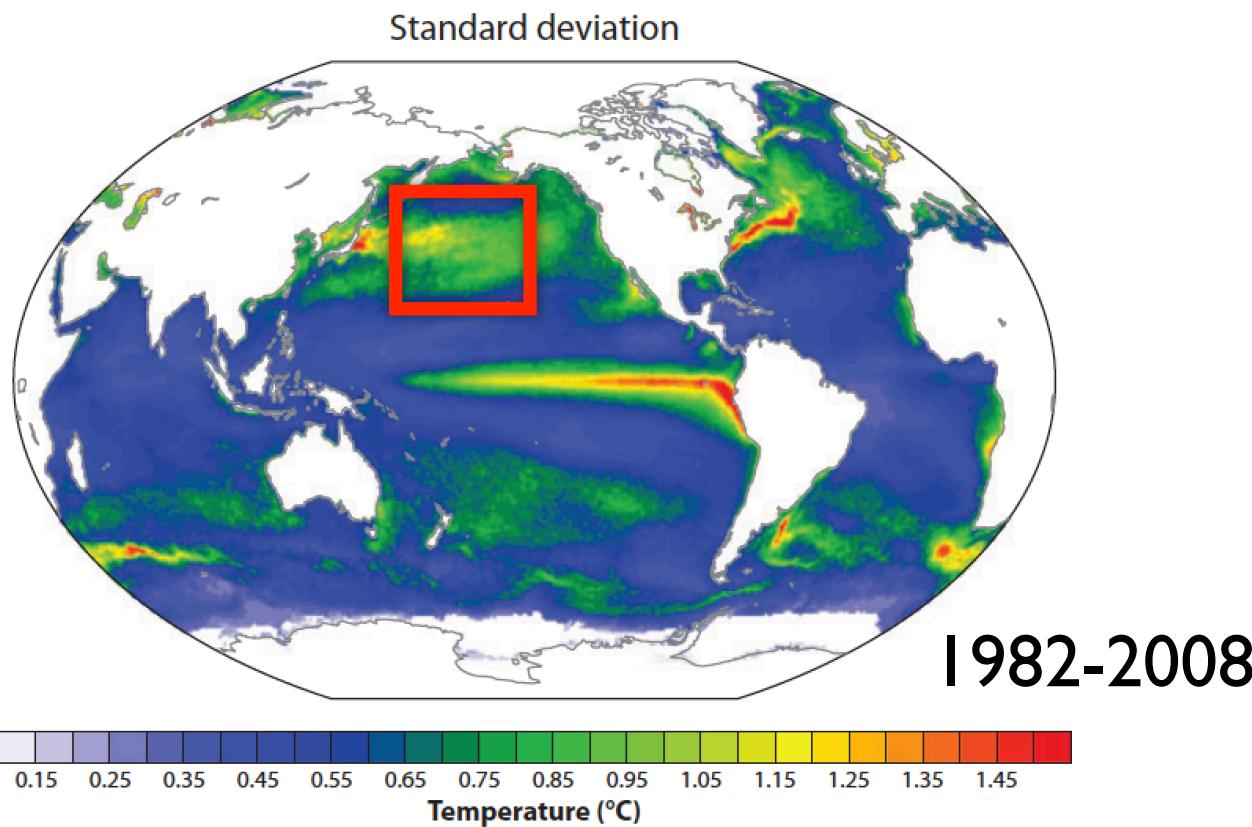
0-70N

b North Atlantic SST index



Question time

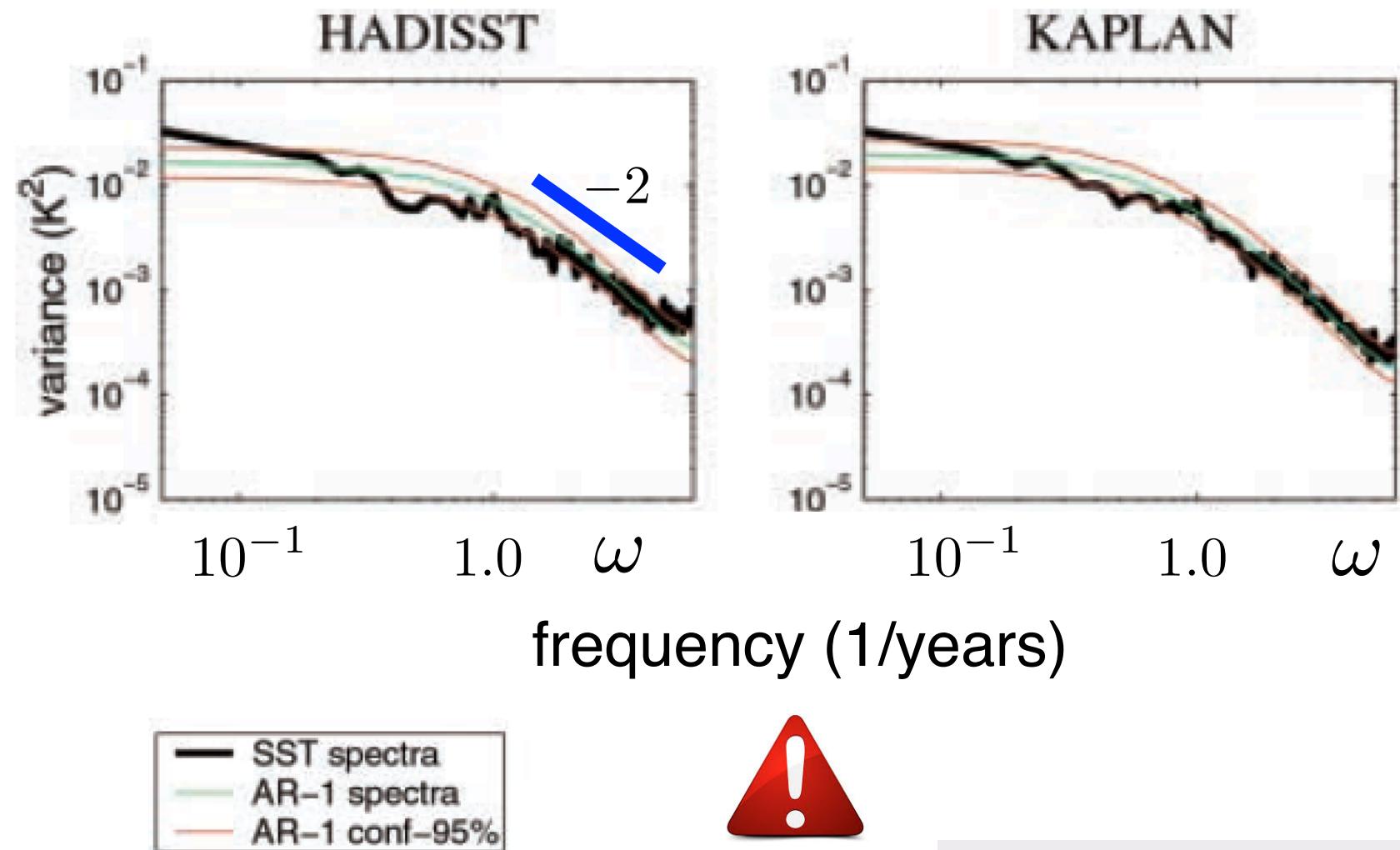
Sea surface temperature variability



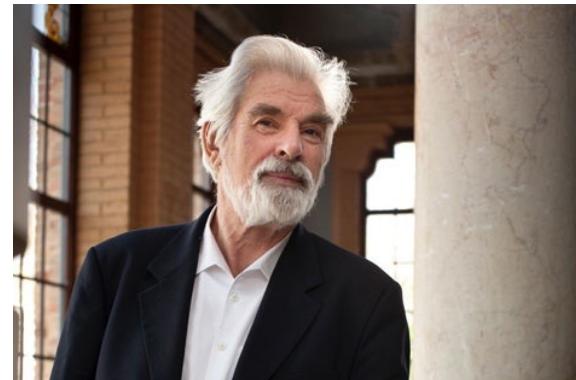
What is the null-hypothesis of sea surface temperature variability?

Midlatitude SST Spectra

25N – 50N Pacific (1903-1994)

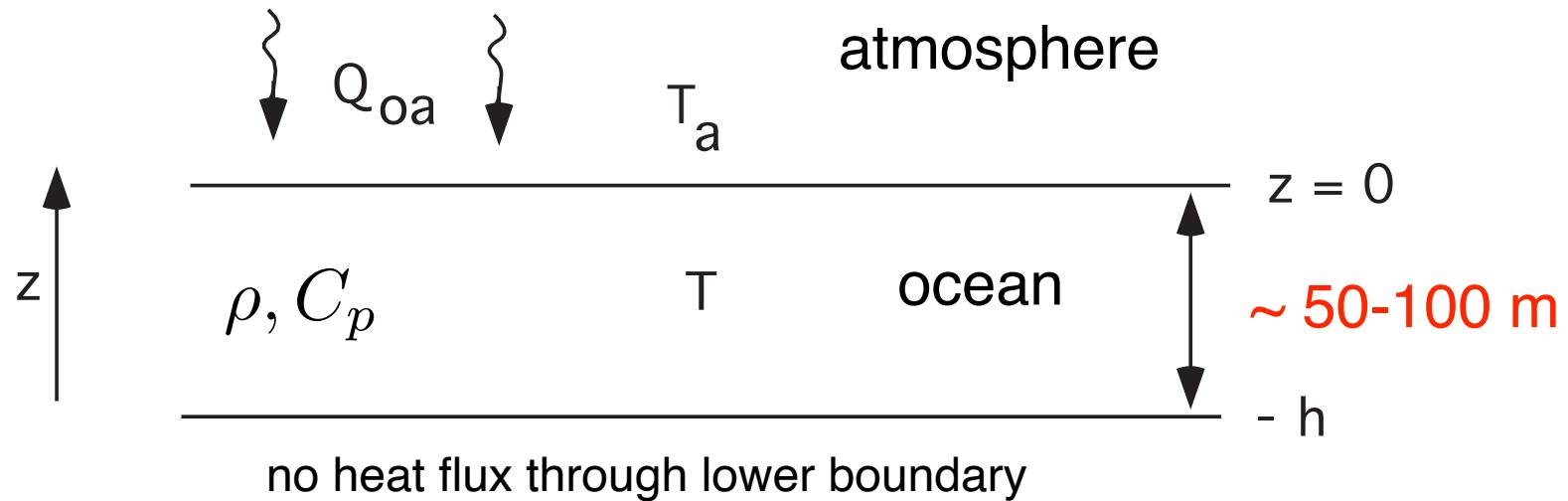


Hasselmann approach (1976)



Hasselmann K. (1976), "Stochastic climate models,
Part 1: Theory", *Tellus*, 28: 473-485.

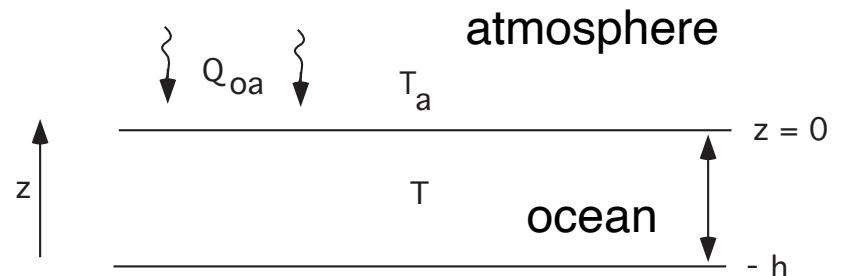
Klaus Hasselmann (1931-)



Ocean Mixed-layer Model

Equation:

$$\rho C_p \frac{\partial T}{\partial t} = K \frac{\partial^2 T}{\partial z^2}$$



Boundary conditions:

$$z = 0 : K \frac{\partial T}{\partial z} = Q_{oa}$$

$$z = -h : \frac{\partial T}{\partial z} = 0$$



Solution

Use: $\bar{T} = \frac{1}{h} \int_{-h}^0 T \ dz$ and $Q_{oa} = \alpha(T_a - \bar{T})$

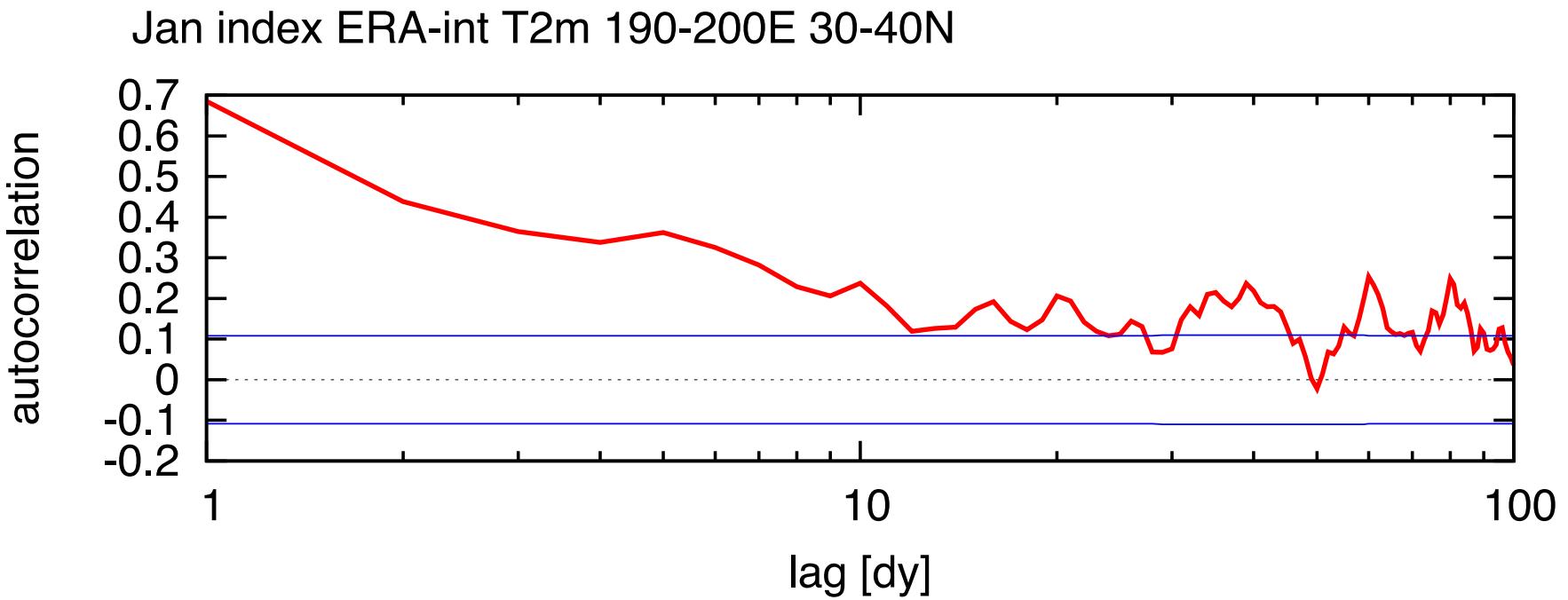
$$\bar{T} = \langle \bar{T} \rangle + \tilde{T}$$
$$T_a = \langle T_a \rangle + \tilde{T}_a$$

Result:

$$\frac{d\tilde{T}}{dt} = \frac{\alpha}{\rho C_p h} (\tilde{T}_a - \tilde{T})$$

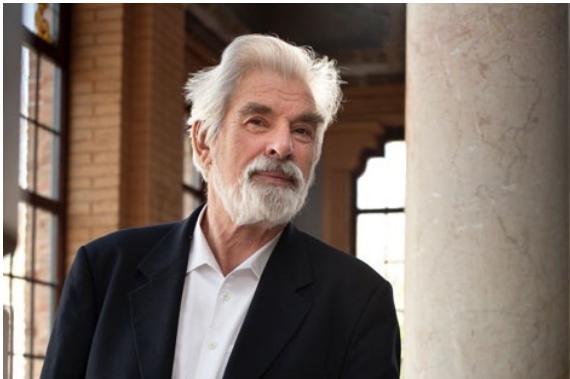
$$\gamma = \frac{\alpha}{\rho C_p h} \sim 1/(100 \text{ days})$$

Example: Autocorrelation Pacific atmospheric surface temperatures



Decorrelation time scale atmospheric forcing << ocean damping time scale

The Hasselmann (1976) stochastic climate model



Hasselmann K. (1976), "Stochastic climate models,
Part 1: Theory", *Tellus*, 28: 473-485.

$$\frac{d\tilde{T}}{dt} = -\gamma \tilde{T} + \sigma \xi$$

$$\gamma = \frac{\alpha}{\rho C_p h}$$

white noise

$$E[\xi(t)] = 0$$

$$E[\xi(t)\xi(s)] = \delta(t - s)$$

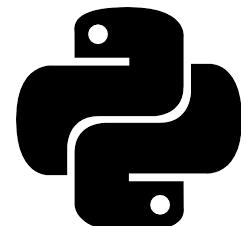
“The choice between a deterministic and a stochastic formulation of the equations [is] dictated by convenience” Lorenz (1987)

Question time

Stochastic Differential Equations (SDEs)

from

$$\frac{d\tilde{T}}{dt} = -\gamma\tilde{T} + \sigma\xi$$



Stochastic process: $X_t = \tilde{T}$

Wiener process: W_t $N(0, t)$ distributed
 $E[(dW_t)^2] = dt$

to

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

or

$$X_t = X_0 + \int_0^t (-\gamma X_s) ds + \int_0^t \sigma dW_s$$

Stochastic integrals

Kiyoshi $\hat{It\ddot{o}}$ (1915-2008)



$$\int_0^T h(t) \, dW_t = \lim_{N \rightarrow \infty} \sum_{j=0}^{N-1} h(t_j)(W(t_{j+1}) - W(t_j))$$

‘left end point’

Ruslan Stratonovich (1930-1997) ‘mid point’



$$\int_0^T h(t) \circ dW_t = \lim_{N \rightarrow \infty} \sum_{j=0}^{N-1} h\left(\frac{t_j + t_{j+1}}{2}\right)(W(t_{j+1}) - W(t_j))$$

both in the mean-square sense

$$\lim_{N \rightarrow \infty} E[(I - I_N)^2] \rightarrow 0$$

Stochastic integrals: evoluation

Problem: Evaluate $\int_0^T W_t dW_t$

Itô


$$f(W_T) - f(W_0) = \int_0^T f'(W_t) dW_t + \int_0^T \frac{1}{2} f''(W_t) dt$$

Take: $f(t) = t^2$

Answer: $\int_0^T W_t dW_t = \frac{1}{2}(W_T^2 - T)$

Solution of SDEs

$$f(t + dt, X_t + dX_t) - f(t, X_t) = f_1 dt + f_2 dX_t + \frac{1}{2} (f_{11}(dt)^2 + 2f_{12}dtdX_t + f_{22}(dX_t)^2) + \dots$$

Itô



$$dX_t = A_t dt + B_t dW_t$$

$$f(t, X_t) - f(0, X_0) = \int_0^t (f_1 + f_2 A_s + \frac{1}{2} f_{22} B_s^2) ds + \int_0^t f_2 B_s dW_s$$

Example: Hasselmann

$$dX_t = -\gamma X_t dt + \sigma dW_t$$



$$A_t = -\gamma X_t ; \quad B_t = \sigma$$

$$f(t, X_t) - f(0, X_0) = \int_0^t (f_1 + f_2 A_s + \frac{1}{2} f_{22} B_s^2) ds + \int_0^t f_2 B_s dW_s$$

Take: $f(t, x) = x e^{\gamma t}$

Result: $X_t = e^{-\gamma t} (X_0 + \sigma \int_0^t e^{\gamma s} dW_s)$

Ornstein-Uhlenbeck process

Numerical solution of SDEs

$$X(t) = X(0) + \int_0^t f(X(s))ds + \int_0^t g(X(s))dW(s)$$

$$\tau_j = j\Delta t, j = 0, \dots, n \text{ on } [0, T]$$

$$\Delta t = T/n$$

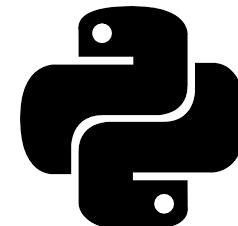


Euler-Maruyama scheme:

Gisiro Maruyama (1916-1986)

$$X_j - X_{j-1} = f(X_{j-1})\Delta t + g(X_{j-1})(W(\tau_j) - W(\tau_{j-1}))$$

Ornstein-Uhlenbeck process



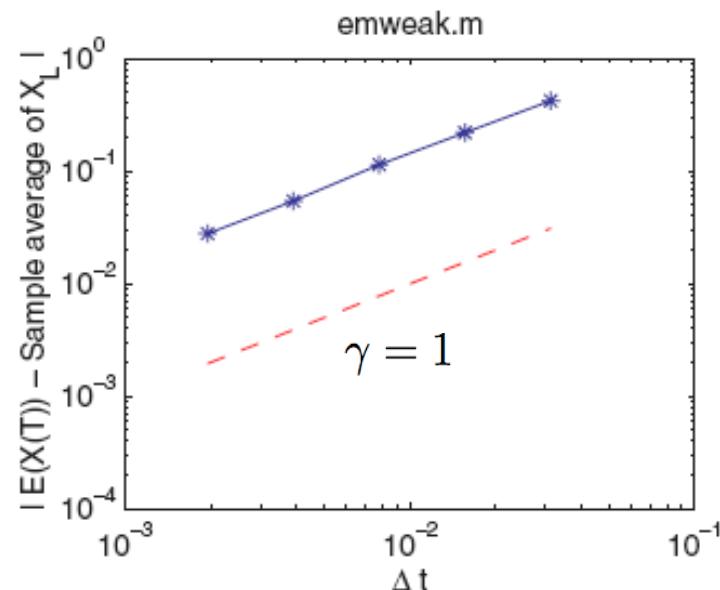
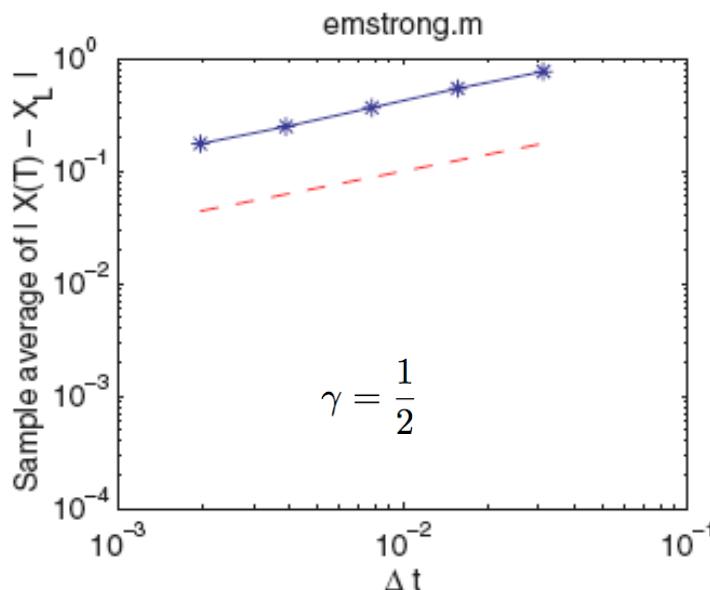
Convergence of EM-method

Strong convergence:

$$E[|X_k - X(\tau_k)|] \leq (\Delta t)^\eta$$

Weak convergence:

$$|E[X_k] - E[X(\tau_k)]| \leq (\Delta t)^\eta$$



Question time

Probability density

$$dX_t = a(X_t, t)dt + b(X_t, t)dW_t$$

$$E[f(X_t)] = \int f(x)p(x, t)dx$$

Probability Density Function (PDF)



Fokker-Planck Equation

$$\frac{\partial p}{\partial t} = - \frac{\partial(ap)}{\partial x} + \frac{1}{2} \frac{\partial^2(b^2 p)}{\partial x^2}$$

+ BC's + IC

Example

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

$$a = -\gamma x ; b = \sigma$$



$$\frac{\partial p}{\partial t} = -\frac{\partial(ap)}{\partial x} + \frac{1}{2} \frac{\partial^2(b^2 p)}{\partial x^2} = 0$$

equilibrium

$$p_e(x) = C e^{-\frac{\gamma x^2}{\sigma^2}}$$

Result:

$$\int_{-\infty}^{\infty} p_e(x) = 1 \rightarrow C = \frac{1}{\sigma} \sqrt{\frac{\gamma}{\pi}}$$

Statistics Ornstein-Uhlenbeck process

$$dX_t = -\gamma X_t dt + \sigma dW_t$$

Solution: $X_t = e^{-\gamma t} (X_0 + \sigma \int_0^t e^{\gamma s} dW_s)$



Wanted: $E[X_t X_{t+s}], s > 0$

Result: $E[X_t X_{t+s}] = e^{-\gamma(2t+s)} (X_0^2 + \sigma^2 \frac{e^{2\gamma t} - 1}{2\gamma})$

$$E[X_t X_{t+s}] \rightarrow \frac{\sigma^2}{2\gamma} e^{-\gamma s}, \quad t \rightarrow \infty$$

Spectrum: $S(\omega) = \frac{\sigma^2}{2\gamma} \mathcal{F}(e^{-\gamma s}) = \frac{\sigma^2}{\omega^2 + \gamma^2}$

Summary: Ornstein-Uhlenbeck process



Leonard Ornstein
(1880-1941)



George Uhlenbeck
(1900-1988)

Fokker-Planck equation:

$$\frac{\partial p}{\partial t} = \frac{\partial(\gamma xp)}{\partial x} + \frac{\sigma^2}{2} \frac{\partial^2 p}{\partial x^2}$$

$$x \rightarrow \pm\infty : p \rightarrow 0$$

Stationary density:

$$p_{stat} = \sqrt{\frac{\gamma}{\pi}} \frac{1}{\sigma} e^{\frac{-\gamma x^2}{\sigma^2}}$$

Stationary autocorrelation:

$$E[X_t X_s] = \frac{\sigma^2}{2\gamma} e^{-\gamma|\tau|}, \tau = t - s$$

Stationary spectrum:

$$S(\omega) = \frac{\sigma^2}{\gamma^2 + \omega^2}$$

The Red Noise (AR(1)) process

$$X_t = e^{-\gamma t} \left(X_0 + \sigma \int_0^t e^{\gamma s} dW_s \right)$$

Discrete time:

$$t_n = n\Delta t ; \quad X_n = X_{t_n}$$



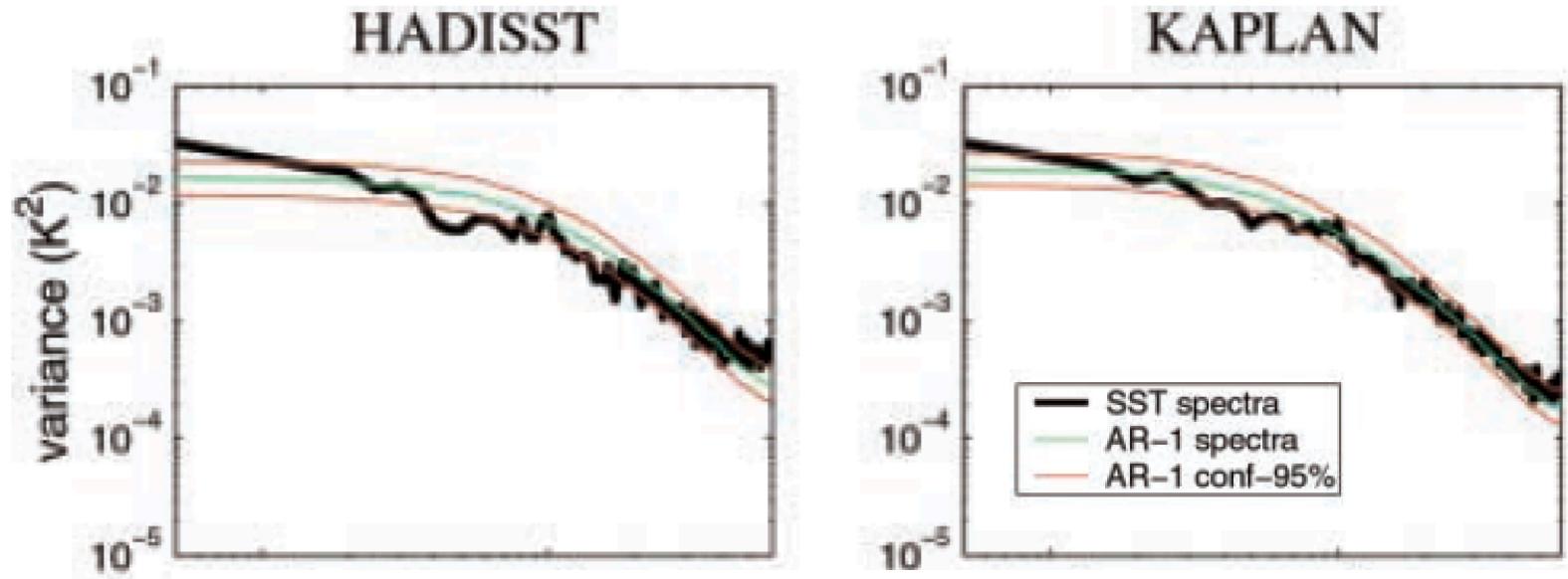
$$X_{n+1} = \alpha X_n + Z_{n+1}$$

$$\alpha = e^{-\gamma \Delta t}, \quad Z_n = \sigma dW_n$$

α can be estimated from the autocorrelation of the time series



Red noise representation



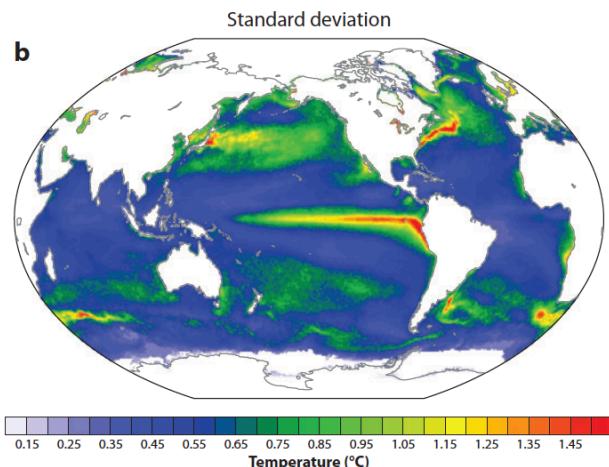
The red noise spectrum serves as a null-hypothesis for climate variability!

Question time

Summary: SST variability

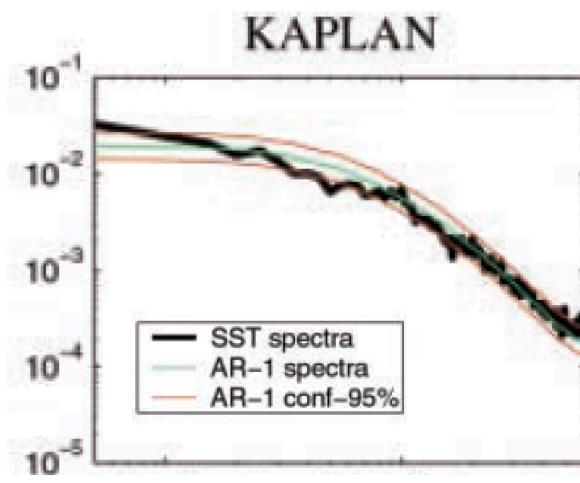
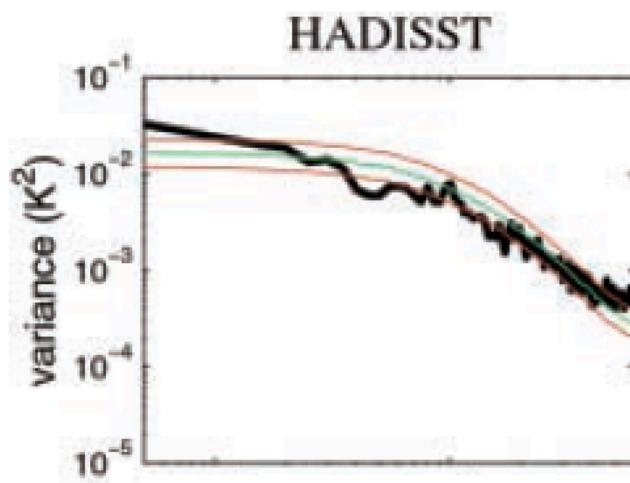
Ocean mixed layer temperature:

$$dX_t = -\gamma X_t dt + \sigma dW_t$$



Wiener process

Continuous: Ornstein-Uhlenbeck process
Discrete: red noise or AR(1) process



The red noise spectrum serves as a null-hypothesis for climate variability!

End of Lecture 1