Answers final exam UCU SCI 211, December 18, 2003

1) According to the theory, for instance of Section 8.2 in the Guide Book wit $\gamma = 1$ and L = 1, the solution is of the form

$$u(x, t) = \sum_{k=0}^{\infty} a_k e^{-\pi^2 k^2 t} \cos(k\pi x),$$

in which substitution of t = 0 yields

$$1 + \cos(\pi x) = u(x, 0) = \sum_{k=0}^{\infty} a_k \cos(k\pi x).$$

Here the left hand side is a cosine series of the same form as appears in the right hand side, with $a_0 = 1$, $a_1 = 1$ and $a_k = 0$ for every k > 1. (One can also use the integral formulas for the Fourier cofficients a_k , which is a little more work.) The conclusion is that

$$u(x, t) = 1 + e^{-\pi^2 t} \cos(\pi x).$$

It follows that

$$u(x, t) - 1| = \left| e^{-\pi^2 t} \cos(\pi x) \right| = e^{-\pi^2 t} \left| \cos(\pi x) \right| \le e^{-\pi^2 t},$$

because $|\cos(\pi x)| \leq 1$. Substitution of t = 1 yields the desired inequality.

2) Let f(x) be a twice differentiable function on the real axis which is periodic with period 2L and has the property that f(x) = 0 when $-L \le x \le 0$. Define

$$u(x, t) = f(t+x) - f(t-x).$$

a) Differentiating u(x, t) twice with respect to t we obtain f''(t+x) - f''(t-x). Differentiating u(x, t) with respect to x we obtain f'(t+x) + f'(t-x), and differentiating this again with respect to x we obtain f''(t+x) - f''(t-x), which is equal to the second order derivative of u(x, t) with respect to t. This shows that u(x, t) satisfies the wave equation.

Putting x = 0 we get u(0, t) = f(t) - f(t) = 0. Putting x = L we obtain f(t + L) - f(t - L) = 0, because x + L = (x - L) + 2L and f was assumed to be periodic with period 2L. This shows that u(x, t) satisfies the boundary conditions.

Finally, if 0 < x < 3 then u(x, 0) = f(x) - f(-x) = f(x), because x > 0 implies -x < 0 and therefore f(-x) = 0. If we put t = 0 in $\partial u(x, t)/\partial t = f'(t+x) - f'(t-x)$ we obtain that the initial velocity is equal to f'(x) - f'(-x) = f'(x), because f(y) = 0 for all y < 0 implies that f'(y) = 0 for all y < 0. This shows that u(x, t) satisfies the initial conditions.

b) The pictures should look like the following.



In the interval $0 \le x \le 3$, the bump is translated to the left with speed 1 until it hits the left end point of the interval, whereas after time t = 2 it has reappeared upside down and is shifting to the right with speed equal to 1.

c) We have u(x, 1.5) = f(1.5 + x) - f(1.5 - x) = 0, in which the second equation just expresses the symmetry of the function f with repect to the reflection $1.5 + x \mapsto 1.5 - x$. This is not such a paradox as it may seem, because at time t = 1.5 the velocity $\partial u(x, t)/\partial t$ is nonzero, the string just moves through the horizontal position at time t = 1.5.

3) If we substitute u(x, y) = f(x) g(y) in the partial differential equation, we obtain

$$f''(x) g(y) + f(x) g''(y) + 2f'(x) g(y) + f(x) g(y) = 0$$

Assuming that u(x, y) = f(x) g(y) is nonzero we can divide by f(x) g(y) and obtain the equivalent equation

$$\frac{f''(x)}{f(x)} + \frac{g''(y)}{g(y)} + 2\frac{f'(x)}{f(x)} + 1 = 0,$$

or equivalently

$$\frac{f''(x)}{f(x)} + 2\frac{f'(x)}{f(x)} + 1 = -\frac{g''(y)}{g(y)}.$$

The left hand side does not depend on y, and is equal to the right hand side which does not depend on x, and therefore it is equal to a constant λ neither depending on y nor on x. This implies that our equation is equivalent to the two equations

$$\frac{f''(x)}{f(x)} + 2\frac{f'(x)}{f(x)} + 1 = \lambda = -\frac{g''(y)}{g(y)},$$

or equivalently

$$f''(x) + 2f'(x) + (1 - \lambda)f(x) = 0, \quad g''(y) + \lambda g(y) = 0$$

in which λ is an arbitrary constant. The exponential function e^{px} is a solution for the equation for f(x) if and only if $p^2 + 2p + (1 - \lambda) = 0$, or $p = -1 \pm \sqrt{\lambda}$. If $\lambda \neq 0$ then, using the two-dimensionality of the space of solutions of the differential equation for f(x), we obtain that the general solution is

$$f(x) = c_+ e^{(-1+\sqrt{\lambda})x} + c_- e^{(-1-\sqrt{\lambda})x},$$

in which c_+ and c_- are arbitrary constants. Similarly e^{qy} is a solution of the differential equation for g(y) if and only if $q^2 + \lambda = 0$, and we have

$$g(y) = d_+ e^{\sqrt{-\lambda}y} + d_- e^{-\sqrt{-\lambda}y},$$

in which d_+ and d_- are arbitrary constants, and therefore the answer is

$$u(x, y) = \left(c_{+} e^{(-1+\sqrt{\lambda})x} + c_{-} e^{(-1-\sqrt{\lambda})x}\right) \left(d_{+} e^{\sqrt{-\lambda}y} + d_{-} e^{-\sqrt{-\lambda}y}\right),$$

in which λ , c_- , c_+ , d_- , d_+ are arbitrary complex numbers. The only exceptional case occurs for $\lambda = 0$, in which case this should be replaced by

$$u(x, y) = (a e^{-x} + b x e^{-x}) (c + d y),$$

in which a, b, c, d are arbitrary constants.

4)

a) We have

$$d(x(t)^{2} - y(t)^{2}) / dt = 2x'(t) x(t) - 2y'(t) y(t) = 2y(t) x(t) - 2x(t) y(t) = 0.$$

This implies that $x(t)^2 - y(t)^2$ is constant as a function of t and therefore equal to $x(0)^2 - y(0)^2 = 1^2 - 0^2 = 1$.

b) The Euler approximation for this system of ordinary differential equations is

$$x_{n+1} = x_n + h y_n$$
, $y_{n+1} = y_n + h x_n$, $x_0 = 1$, $y_0 = 0$.

The identity for $(x_n)^2 - (y_n)^2$ obviously holds when n = 0. Now assume that it holds for some given n. Then

$$(x_{n+1})^{2} - (y_{n+1})^{2} = (x_{n} + h y_{n})^{2} - (y_{n} + h x_{n})^{2} = (1 - h^{2}) ((x_{n})^{2} - (y_{n})^{2})$$
$$= (1 - h^{2}) (1 - h^{2})^{n} = (1 - h^{2})^{n+1},$$

which is the identity for $(x_n)^2 - (y_n)^2$ with everywhere *n* replaced by n+1. Here we have used the Euler procedure in the first equality, a straightforward calculation in the second equality, and the induction hypothesis in the third equality. This completes the proof of the identity for $(x_n)^2 - (y_n)^2$ for every $n \ge 0$.

c) (bonus) The first identity is obtained by the substitution of n = N, h = t/N, where we subsequently write

$$\left(1 - \frac{t^2}{N^2}\right)^N = \mathrm{e}^{N \ln\left(1 - \frac{t^2}{N^2}\right)} \,.$$

For the second and the third identity we use the first order Taylor expansion which says that for any differentiable function f(x) we have f(x) = f(a) + f'(y)(x - a), in which y is some point between a and x. If we do this for the function $f(x) = \ln x$, with a = 1, $x = 1 - t^2/N^2$ and y = p, we obtain the second identity. We obtain the third identity if we do this for $f(x) = e^x$, with a = 0,

$$x = -\frac{1}{p} \frac{t^2}{N}$$

and y = q.

The last statement follows if we write $\theta_N = e^q / p$ and observe that the inequalities for p and q imply that $p \to 1$ and $q \to 0$ as $N \to \infty$.