

BIFURCATION PHENOMENA

Lecture 3: Two-parameter bifurcations of planar ODEs

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Literature

1. Yu.A. Kuznetsov *Elements of Applied Bifurcation Theory*, 3rd ed. Applied Mathematical Sciences 112, Springer-Verlag, New York, 2004
2. L.P. Shilnikov, A.L. Shilnikov, D.V. Turaev, and L.O. Chua *Methods of Qualitative Theory in Nonlinear Dynamics*, Part II, World Scientific, Singapore, 2001
3. V.I. Arnol'd, V.S. Afraimovich, Yu.S. Il'yashenko, and L.P. Shil'nikov *Bifurcation theory*, In: V.I. Arnol'd (ed), *Dynamical Systems V. Encyclopaedia of Mathematical Sciences*, Springer-Verlag, New York, 1994
4. A.D. Bazykin *Nonlinear Dynamics of Interacting Populations*, World Scientific Publishing, River Edge, NJ, 1998

1. LOCAL CODIM 2 BIFURCATIONS IN 2D

Consider a smooth 2D system depending on two parameters

$$\dot{X} = f(X, \alpha), \quad X \in \mathbb{R}^2, \quad \alpha \in \mathbb{R}^2.$$

Curves of codim 1 bifurcations:

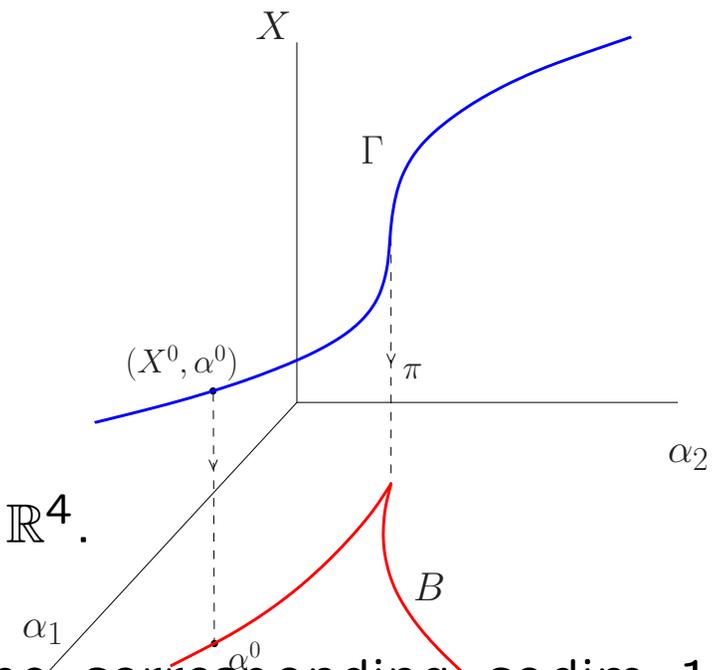
$$\text{Fold} : \begin{cases} f(X, \alpha) = 0, \\ \det f_X(X, \alpha) = 0. \end{cases}$$

$$\text{Hopf} : \begin{cases} f(X, \alpha) = 0, \\ \text{Sp } f_X(X, \alpha) = 0. \end{cases}$$

In both cases, we have $3=2+1$ equations in \mathbb{R}^4 .

When we cross $B = \pi\Gamma$ in the α -plane, the corresponding codim 1 bifurcation occurs.

One has to check that $\lambda_{1,2} = \pm i\omega$ along the Hopf curve.



Local codim 2 cases in the plane:

$$\text{Fold : } \lambda_1 = 0 \quad \begin{cases} \dot{x} = ax^2 + O(3) \\ \dot{y} = \lambda_2 y + O(2) \end{cases} \begin{matrix} \nearrow \\ \searrow \end{matrix}$$

$$\textcircled{1} \lambda_1 = 0, a = 0$$

$$\textcircled{2} \lambda_1 = 0, \lambda_2 = 0$$

$$\text{Hopf : } \lambda_{1,2} = \pm i\omega \quad \begin{cases} \dot{\rho} = l_1 \rho^3 + O(4) \\ \dot{\varphi} = \omega + O(1) \end{cases} \begin{matrix} \nearrow \\ \searrow \end{matrix}$$

$$\textcircled{3} \lambda_{1,2} = \pm i\omega, l_1 = 0$$

To meet each case, we need to “tune” two parameters while following Γ (or B) \Rightarrow codim 2.

Cusp bifurcation: $\lambda_1 = 0, a = 0$

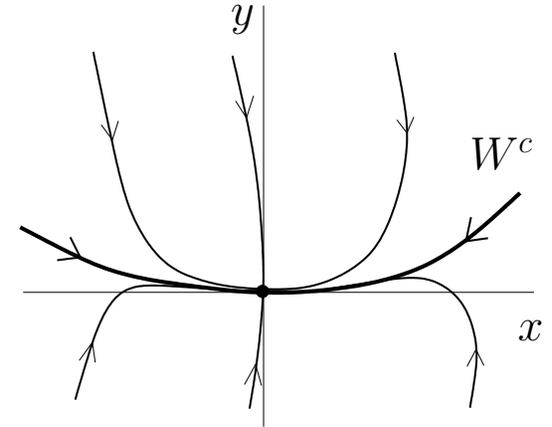
The critical system $\dot{X} = f(X, 0)$ can be transformed by a linear diffeomorphism to

$$\begin{cases} \dot{x} = p_{11}xy + \frac{1}{2}p_{02}y^2 + \frac{1}{6}p_{30}x^3 + \dots, \\ \dot{y} = \lambda_2y + \frac{1}{2}q_{20}x^2 + q_{11}xy + \frac{1}{2}q_{02}y^2 + O(3). \end{cases}$$

It has an invariant 1D **center manifold** $W^c = \{(x, y) : y = W(x)\}$:

$$y = W(x) = \frac{1}{2}w_2x^2 + O(3)$$

where $w_2 = -\frac{q_{20}}{\lambda_2}$.



Thus, the restriction of $\dot{X} = f(X, 0)$ to W^c is

$$\dot{x} = cx^3 + O(4), \quad \text{where } c = \frac{1}{6} \left(p_{30} - \frac{3}{\lambda_2} q_{20} p_{11} \right).$$

Cusp normal form

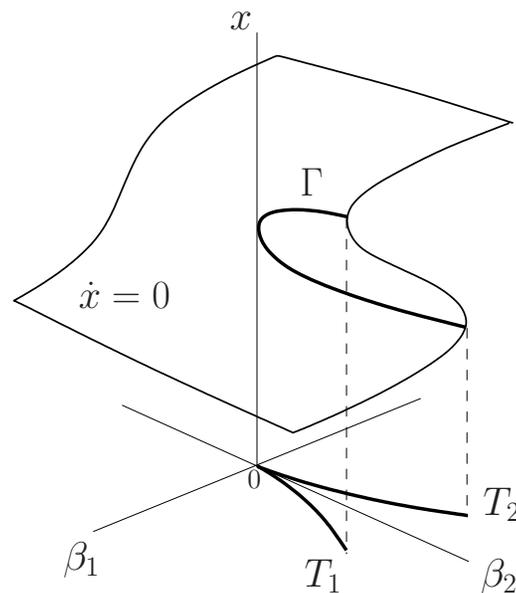
Theorem 1 *If $c \neq 0$, then $\dot{X} = f(X, \alpha)$ is locally topologically equivalent near the cusp bifurcation to*

$$\begin{cases} \dot{x} = \beta_1(\alpha) + \beta_2(\alpha)x + sx^3, \\ \dot{y} = \lambda_2 y, \end{cases}$$

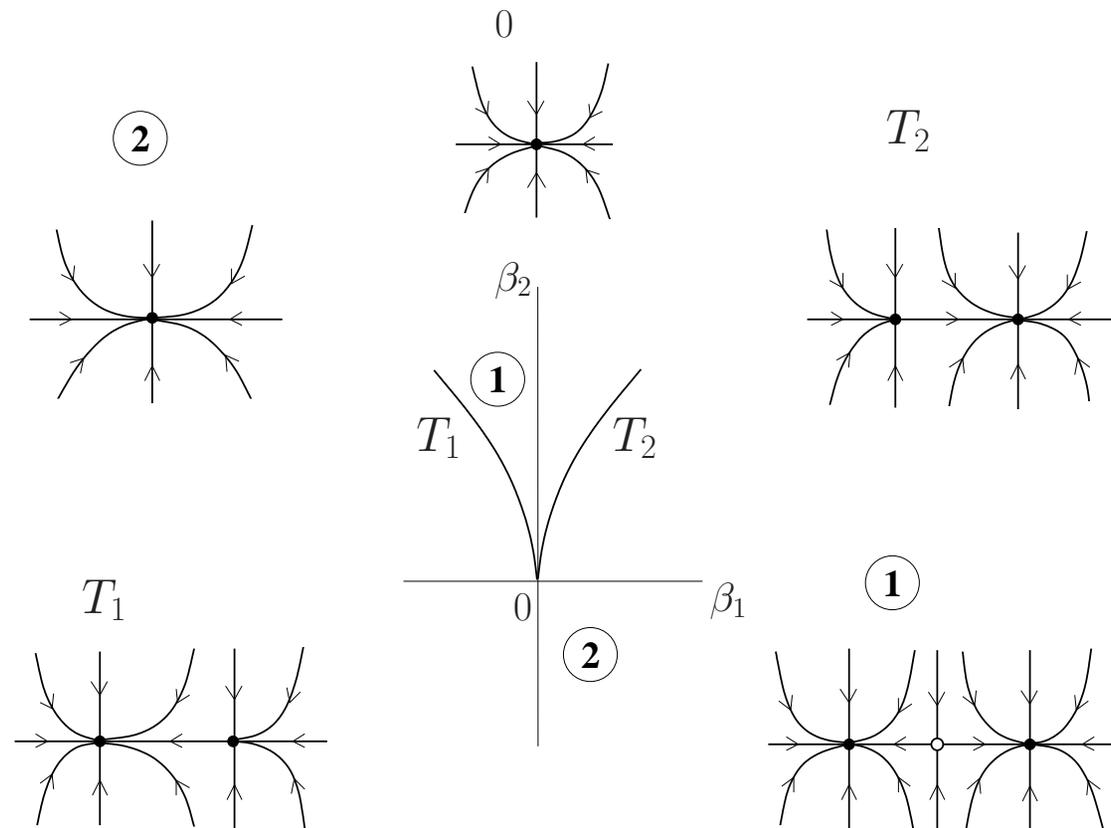
where $\beta_1(0) = \beta_2(0) = 0$ and $s = \text{sign}(c) = \pm 1$.

Fold curve(s) $4\beta_2^3 + 27s\beta_1^2 = 0$

Equilibrium manifold:



Cusp bifurcation diagram ($c < 0, \lambda_2 < 0$)



Three equilibria exist inside the wedge, pairwise colliding at its borders $T_{1,2}$ and leaving one equilibrium outside.

Bogdanov-Takens bifurcation: $\lambda_1 = \lambda_2 = 0$ ($A = f_X(0,0) \neq 0$)

The critical system $\dot{X} = f(X,0)$ can be transformed by a linear diffeomorphism to

$$\begin{cases} \dot{x} = y + \frac{1}{2}p_{20}x^2 + p_{11}xy + \frac{1}{2}p_{02}y^2 + O(3) \equiv P(x,y), \\ \dot{y} = \frac{1}{2}q_{20}x^2 + q_{11}xy + \frac{1}{2}q_{02}y^2 + \frac{1}{6}q_{03}x^2 + O(3). \end{cases}$$

By a nonlinear local diffeomorphism (change of variables)

$$\begin{cases} \xi = x, \\ \eta = P(x,y), \end{cases}$$

this system can be reduced near the origin to

$$\begin{cases} \dot{\xi} = \eta, \\ \dot{\eta} = a\xi^2 + b\xi\eta + \dots, \end{cases}$$

where

$$a = \frac{1}{2}q_{20}, \quad b = p_{20} + q_{11}.$$

Bogdanov-Takens normal form

Theorem 2 *If $ab \neq 0$, then $\dot{X} = f(X, \alpha)$ is locally topologically equivalent near the BT-bifurcation to*

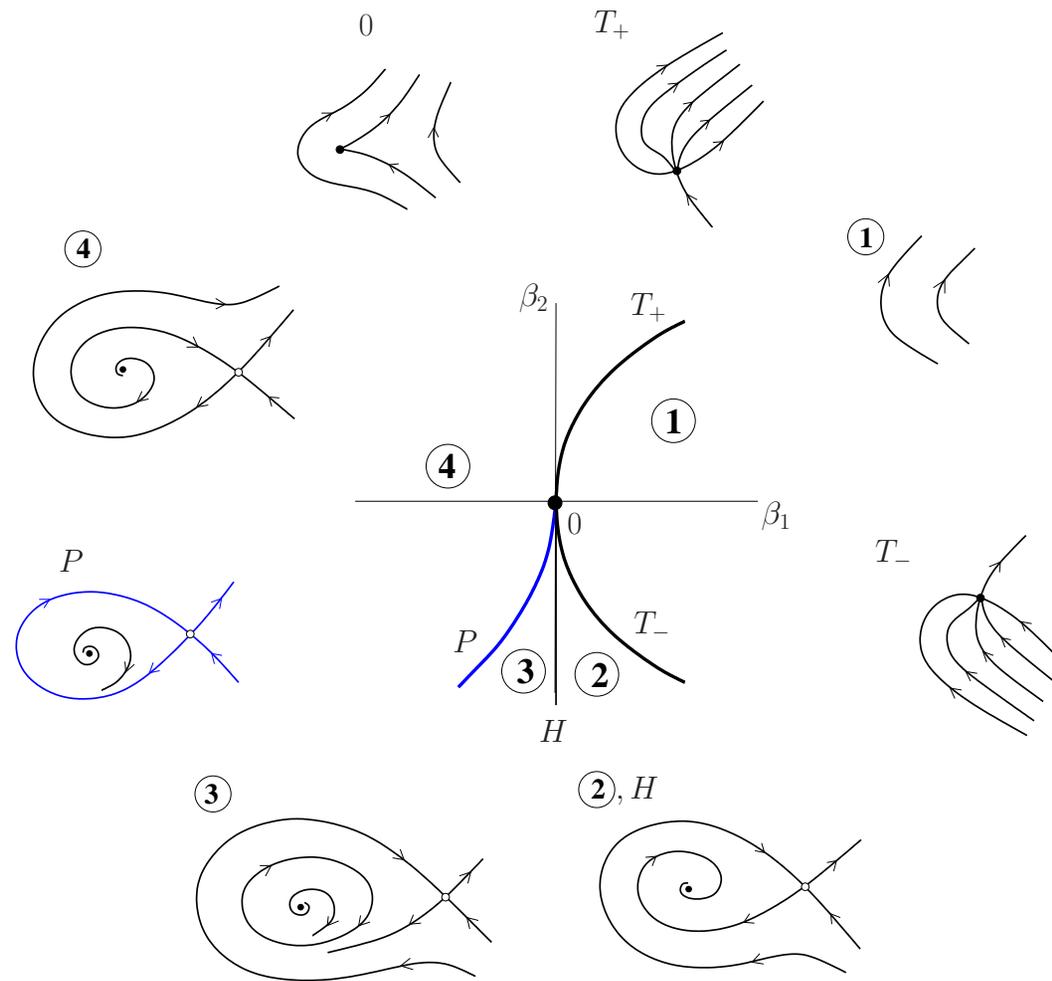
$$\begin{cases} \dot{x} = y, \\ \dot{y} = \beta_1(\alpha) + \beta_2(\alpha)x + x^2 + sxy, \end{cases}$$

where $\beta_1(0) = \beta_2(0) = 0$ and $s = \text{sign}(ab) = \pm 1$.

Bifurcation curves ($ab < 0$):

- **fold** $T : \beta_1 = \frac{1}{4}\beta_2^2$
- **Andronov-Hopf** $H : \beta_1 = 0, \beta_2 < 0$
- **saddle homoclinic** $P : \beta_1 = -\frac{6}{25}\beta_2^2 + O(3), \beta_2 < 0$ (**global bifurcation**)

BT bifurcation diagram ($ab < 0$)



A unique limit cycle appears at Andronov-Hopf bifurcation curve H and disappears via the saddle homoclinic orbit at the curve P .

Bautin (“generalized Hopf”) bifurcation: $\lambda_{1,2} = \pm i\omega$, $l_1 = 0$

The critical system $\dot{X} = f(X, 0)$ can be transformed by a linear diffeomorphism to the complex form

$$\dot{z} = i\omega z + \sum_{2 \leq j+k \leq 5} \frac{1}{j!k!} g_{jk} z^j \bar{z}^k + O(6),$$

which is locally smoothly equivalent to the **Poincaré normal form**

$$\dot{w} = i\omega w + c_1 w |w|^2 + c_2 w |w|^4 + O(6),$$

where the **Lyapunov coefficients**

$$l_j = \frac{1}{\omega} \Re(c_j)$$

satisfy

$$2l_1 = \frac{1}{\omega} \left(\Re(g_{21}) - \frac{1}{\omega} \Im(g_{20}g_{11}) \right) \Rightarrow l_1 = \frac{1}{2\omega^2} \Re(ig_{20}g_{11} + \omega g_{21})$$

If $l_1 = 0$ then

$$\begin{aligned}
12l_2(0) &= \frac{1}{\omega} \Re(g_{32}) \\
&+ \frac{1}{\omega^2} \Im [g_{20}\bar{g}_{31} - g_{11}(4g_{31} + 3\bar{g}_{22}) - \frac{1}{3}g_{02}(g_{40} + \bar{g}_{13}) - g_{30}g_{12}] \\
&+ \frac{1}{\omega^3} \left\{ \Re [g_{20}(\bar{g}_{11}(3g_{12} - \bar{g}_{30}) + g_{02}(\bar{g}_{12} - \frac{1}{3}g_{30}) + \frac{1}{3}\bar{g}_{02}g_{03}) \right. \\
&\quad \left. + g_{11}(\bar{g}_{02}(\frac{5}{3}\bar{g}_{30} + 3g_{12}) + \frac{1}{3}g_{02}\bar{g}_{03} - 4g_{11}g_{30}) \right] \\
&\quad \left. + 3\Im(g_{20}g_{11}) \Im(g_{21}) \right\} \\
&+ \frac{1}{\omega^4} \left\{ \Im [g_{11}\bar{g}_{02}(\bar{g}_{20}^2 - 3\bar{g}_{20}g_{11} - 4g_{11}^2)] \right. \\
&\quad \left. + \Im(g_{20}g_{11}) [3\Re(g_{20}g_{11}) - 2|g_{02}|^2] \right\}
\end{aligned}$$

Bautin normal form

Theorem 3 *If $l_2 \neq 0$ and $\omega \neq 0$, then $\dot{X} = f(X, \alpha)$ is locally topologically equivalent near Bautin bifurcation to the normal form in the polar coordinates:*

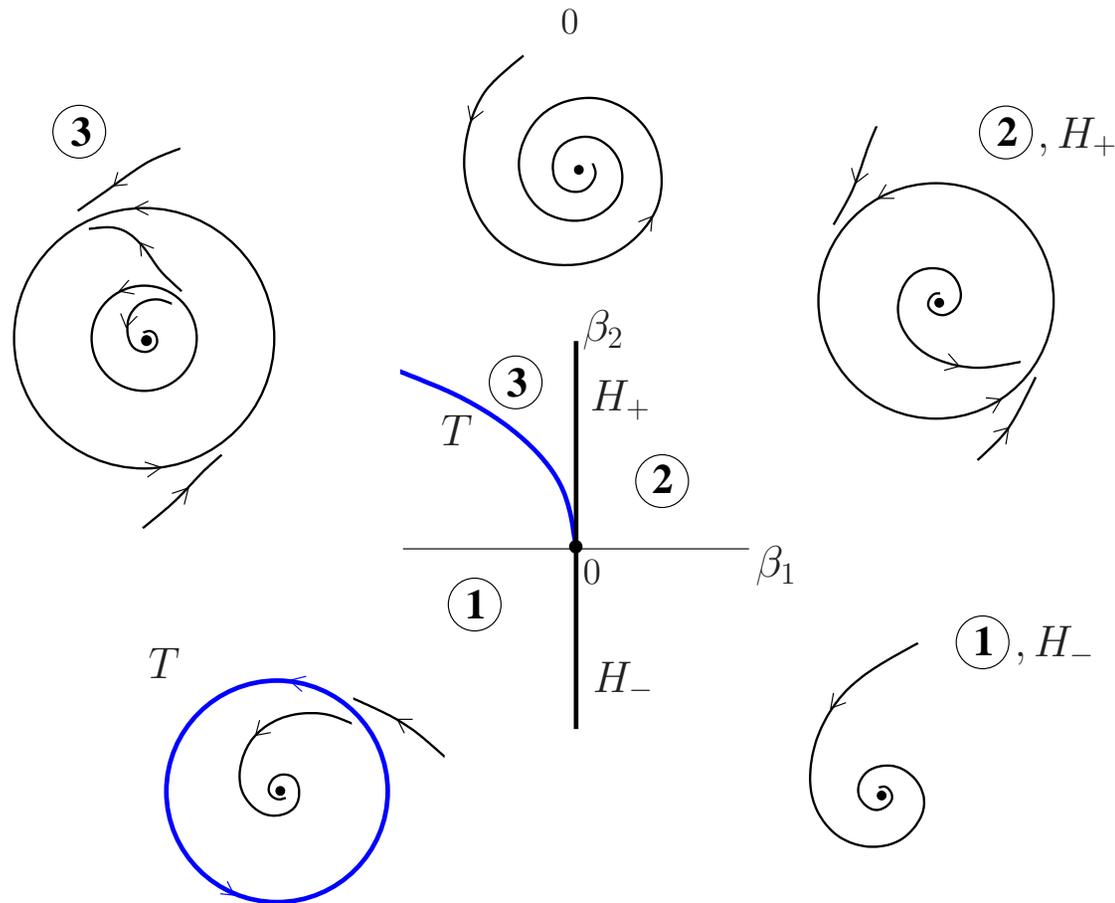
$$\begin{cases} \dot{\rho} = \rho(\beta_1(\alpha) + \beta_2(\alpha)\rho^2 + s\rho^4), \\ \dot{\varphi} = 1, \end{cases}$$

where $\beta_1(0) = \beta_2(0) = 0$ and $s = \text{sign}(l_2) = \pm 1$.

Bifurcation curves ($l_2 < 0$):

- **supercritical Andronov-Hopf H^- :** $\beta_1 = 0, \beta_2 < 0$
- **subcritical Andronov-Hopf H^+ :** $\beta_1 = 0, \beta_2 > 0$
- **cyclic fold T_c :** $\beta_1 = -\frac{1}{4}\beta_2^2, \beta_2 > 0$ (global bifurcation)

Bautin bifurcation diagram ($l_2 < 0$)



In the wedge between H^+ and T_c there exist two limit cycles born via different Andronov-Hopf bifurcations, which merge and disappear at the cyclic fold curve T_c .

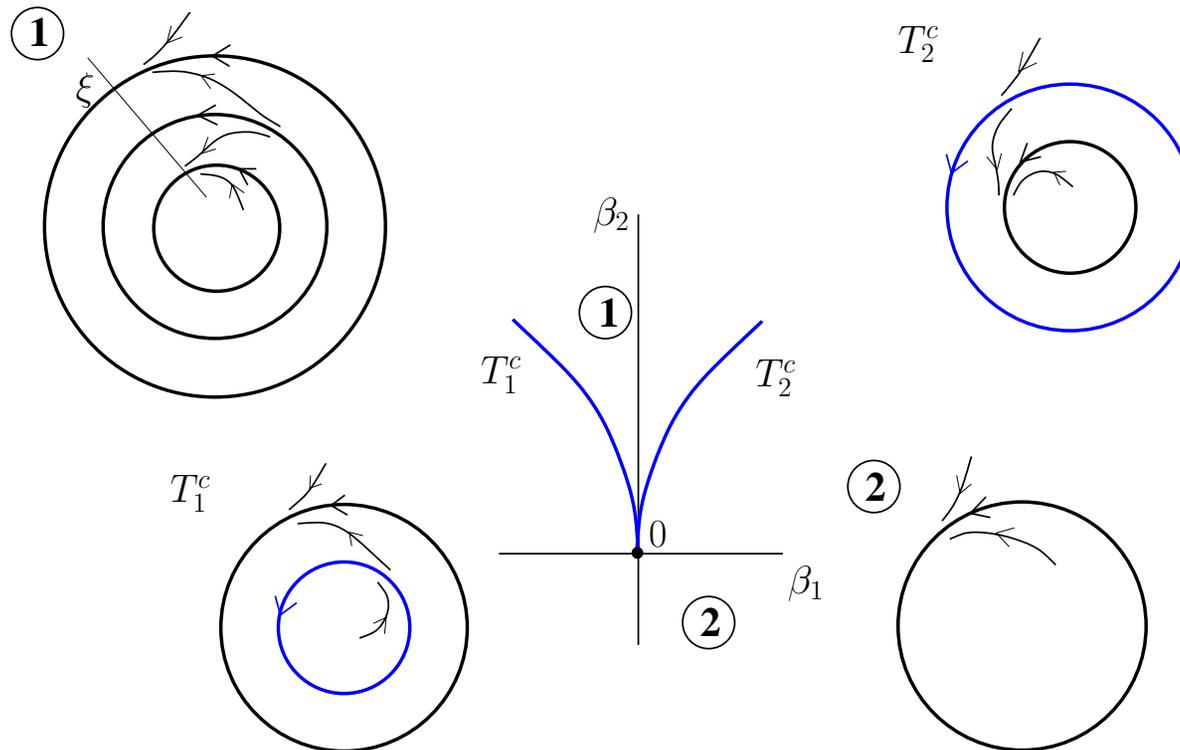
2. SOME GLOBAL CODIM 2 BIFURCATIONS IN 2D

- **Cyclic cusp** ($b = 0$): Critical Poincaré map $\xi \mapsto \xi + c\xi^3 + \dots$

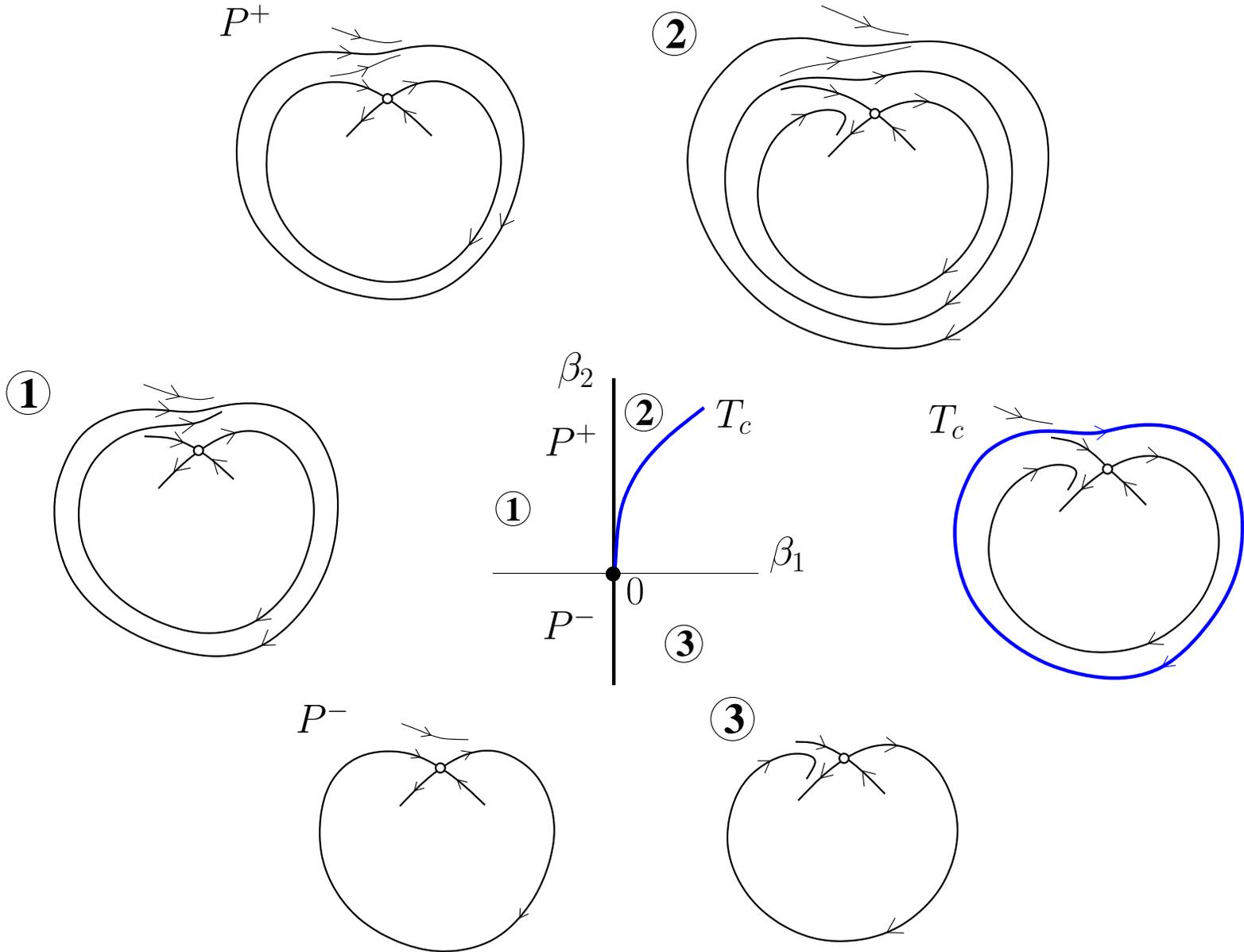
If $c \neq 0$ then the Poincaré map is locally topologically equivalent to

$$\xi \mapsto \beta_1(\alpha) + \beta_2(\alpha)\xi + \xi + s\xi^3,$$

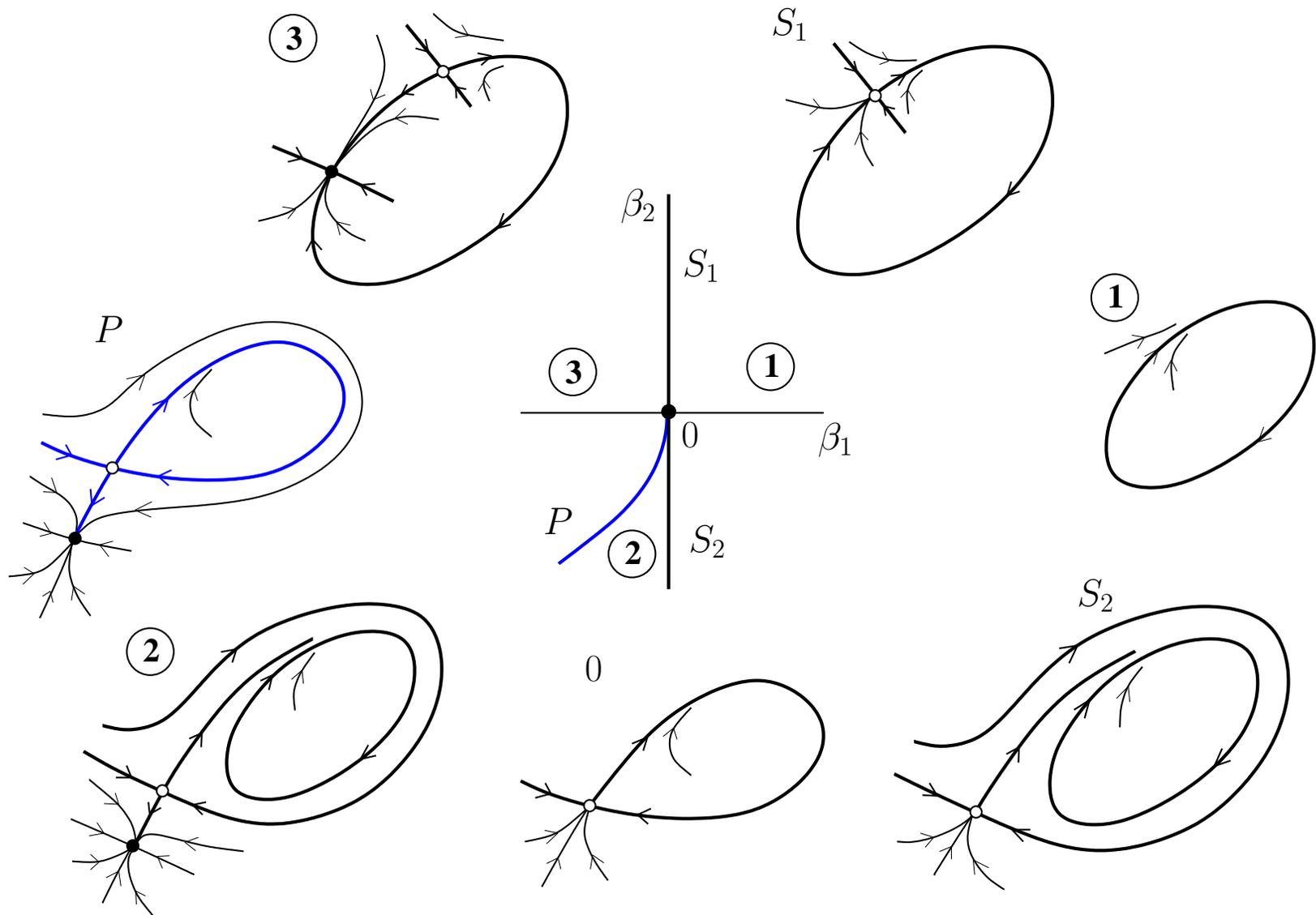
where $\beta_1(0) = \beta_2(0) = 0$ and $s = \text{sign}(c) = \pm 1$.



- Neutral saddle homoclinic orbit: $\int_{-\infty}^{\infty} (\text{div } f)(X^0(t)) dt < 0$

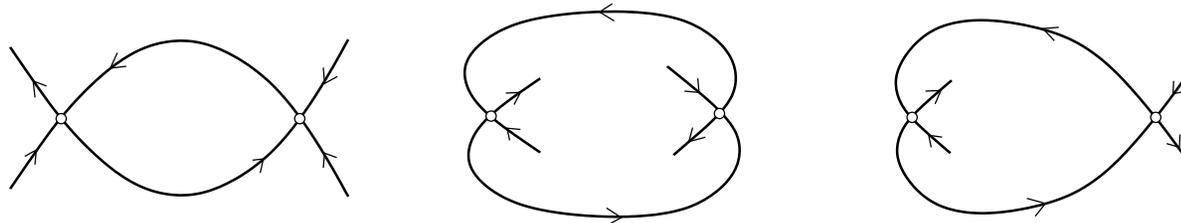


- Non-central saddle-node homoclinic orbit

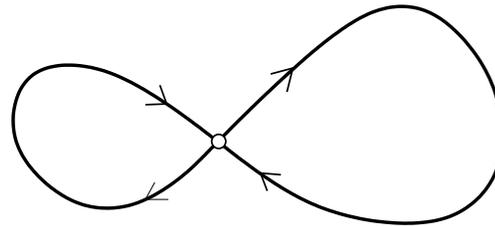


Other global codim 2 cases:

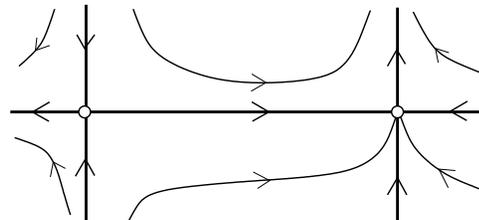
- Heteroclinic cycles



- “Figure of 8”

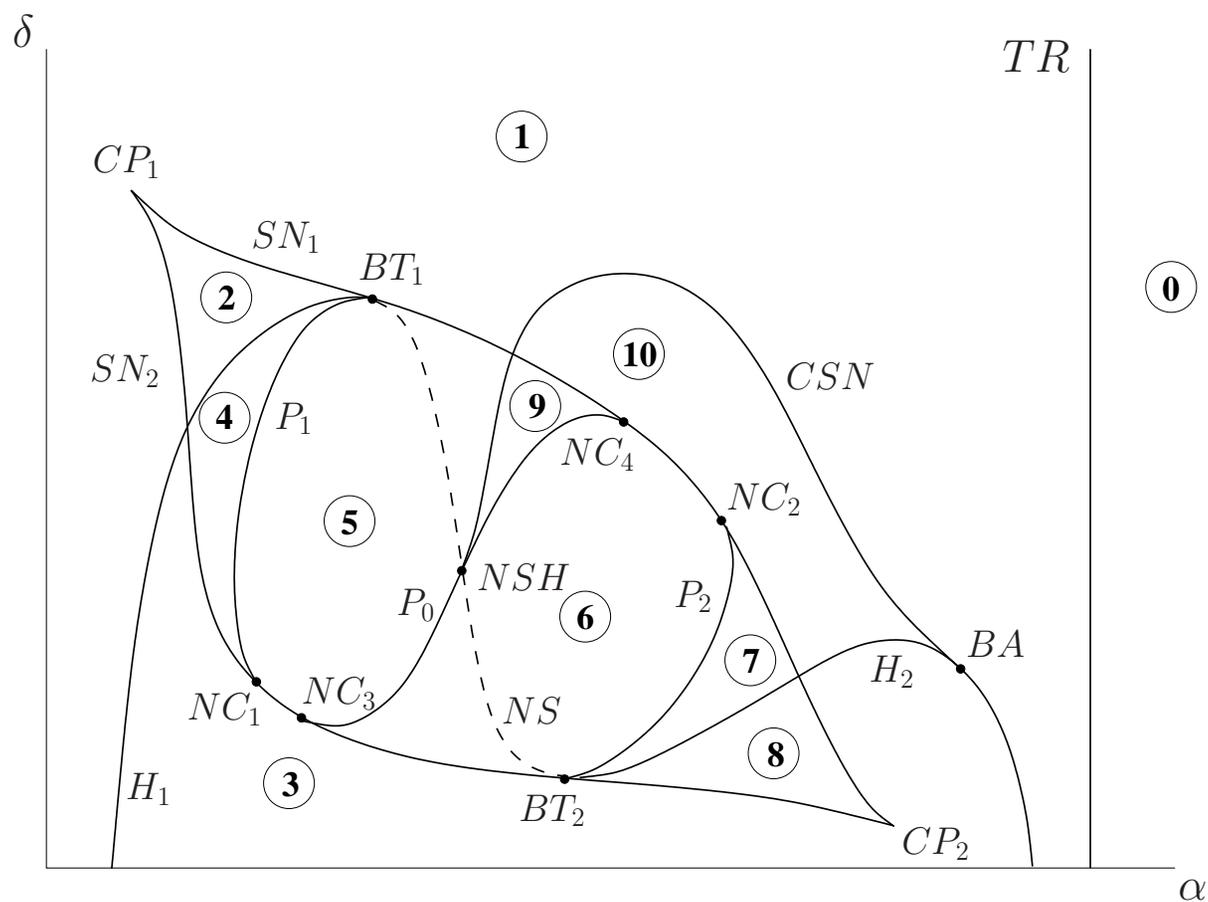


- Saddle-to-saddle-node heteroclinic orbits



Example: Bazykin's prey-predator model

$$\begin{cases} \dot{x}_1 = x_1 - \frac{x_1 x_2}{1 + \alpha x_1} - \varepsilon x_1^2, \\ \dot{x}_2 = -\gamma x_2 + \frac{x_1 x_2}{1 + \alpha x_1} - \delta x_2^2. \end{cases}$$



Generic phase portraits:

