

QCD and the QGP at the LHC

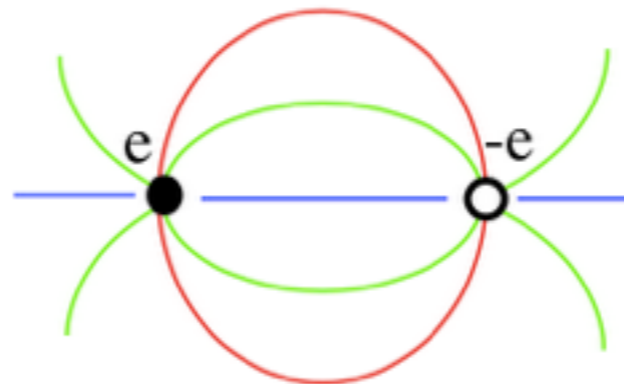
*Marco van Leeuwen,
Nikhef and Utrecht University*

**IoP colloquium UvA,
Amsterdam 19 March 2015**

The QCD potential

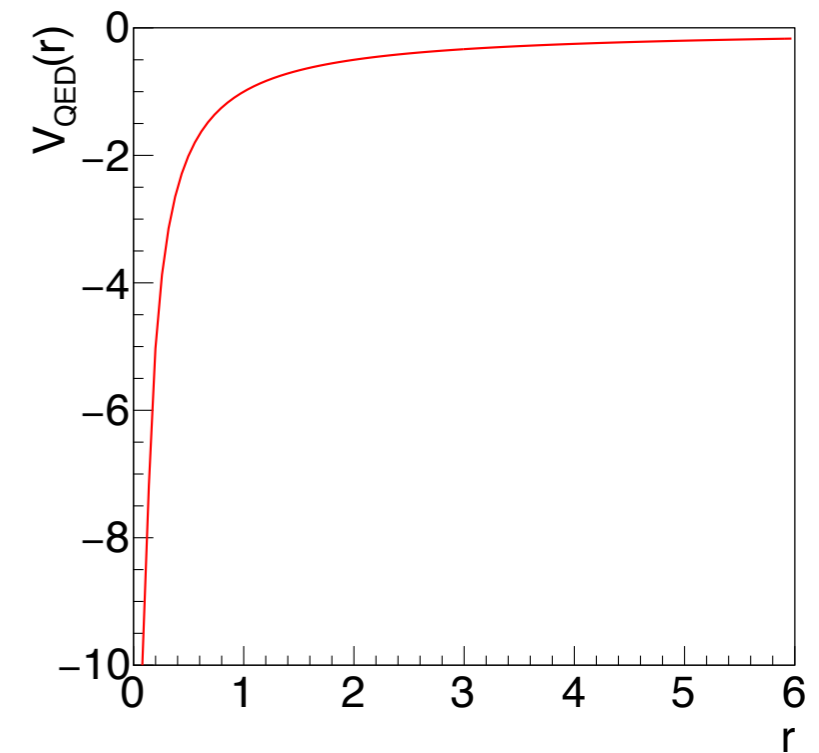
QED
electromagnetic interaction

Field lines
in dipole system

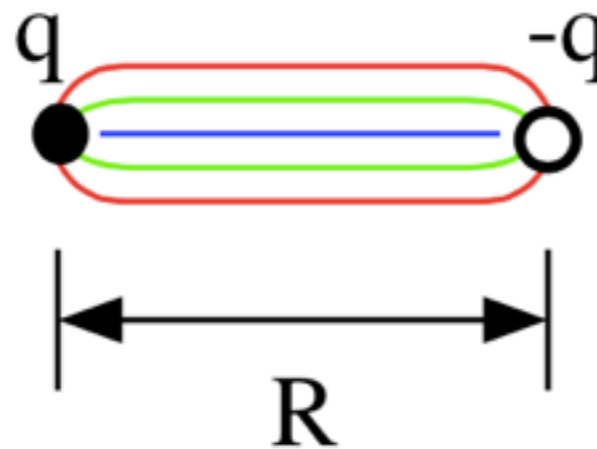


$$V(R) \sim -\alpha / R$$

Potential

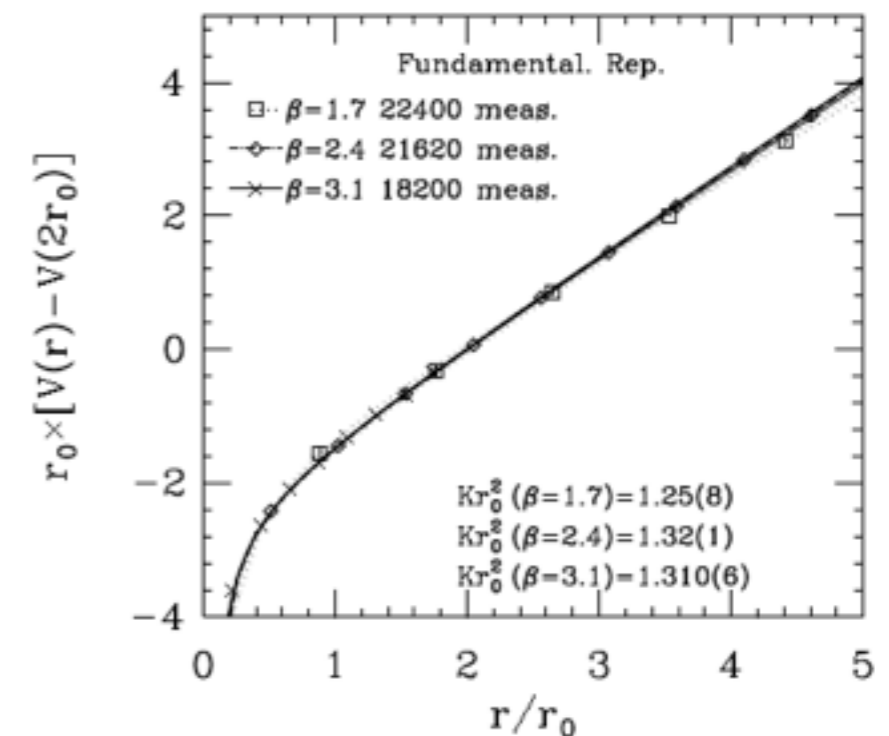


QCD
strong interaction



$$V(R) \sim \sigma R$$

S. Deldar, hep-lat/9909077

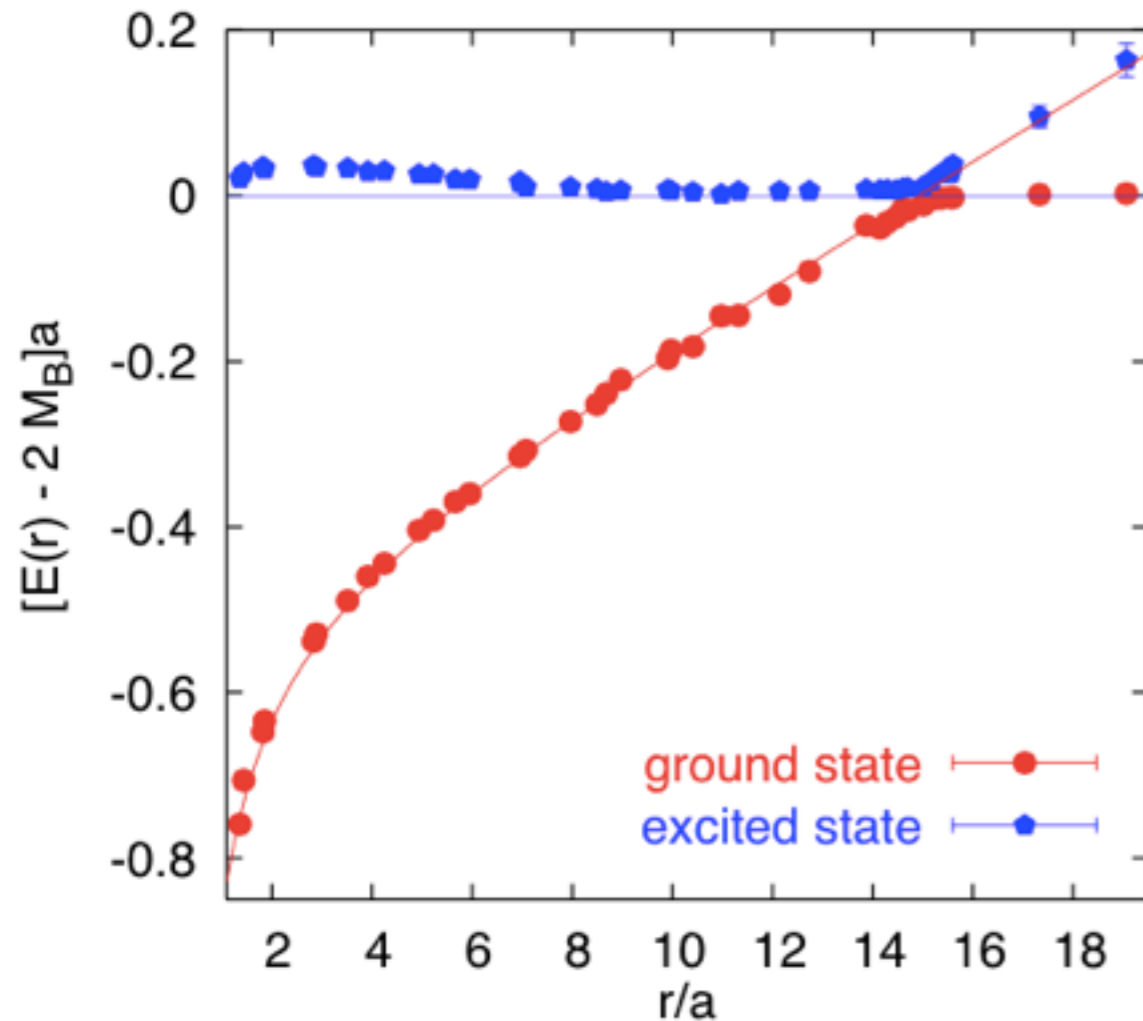


QCD is very different from EM, gravity; common intuition may fail

QCD strings

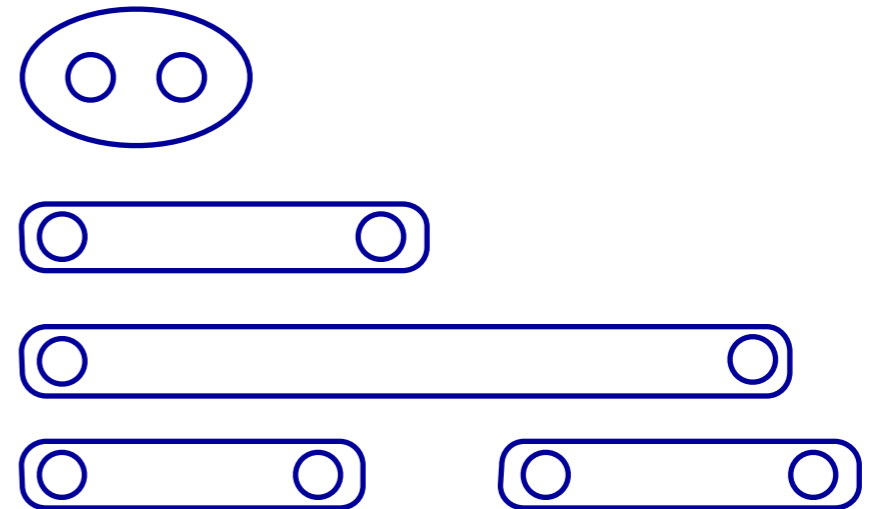
A simple picture of the strong interaction

G.S. Bali, hep-ph/0411206



QCD potential for 2-quark system rises indefinitely

Thought experiment: separating charges



For larger separation: generating a $q\bar{q}$ pair is energetically favoured

Color charges (quarks and gluons) cannot be freed

Confinement important at length scale $1/\Lambda_{\text{QCD}} \sim 1 \text{ fm}$

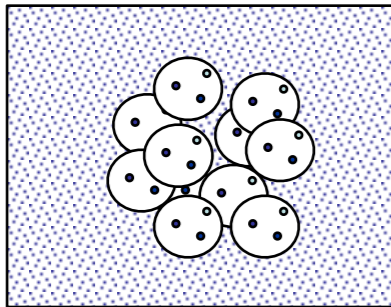
The extremes of QCD

QCD Lagrangian $\mathcal{L}_{QCD} = \bar{\psi}(i\cancel{D} - g\mathbf{A} \cdot \mathbf{t} - m)\psi + \frac{1}{4}\text{Tr}G_{\mu\nu}G^{\mu\nu}$

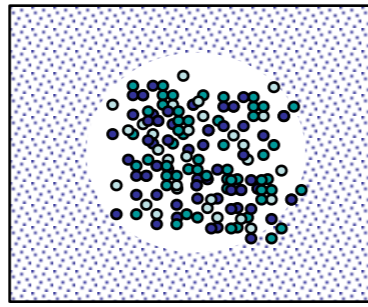
This is the basic theory, but what is the phenomenology?

Bulk QCD matter

Nuclear matter



Quark Gluon Plasma



Calculable with Lattice QCD

High density
Quarks and gluons
are quasi-free

Hard scattering



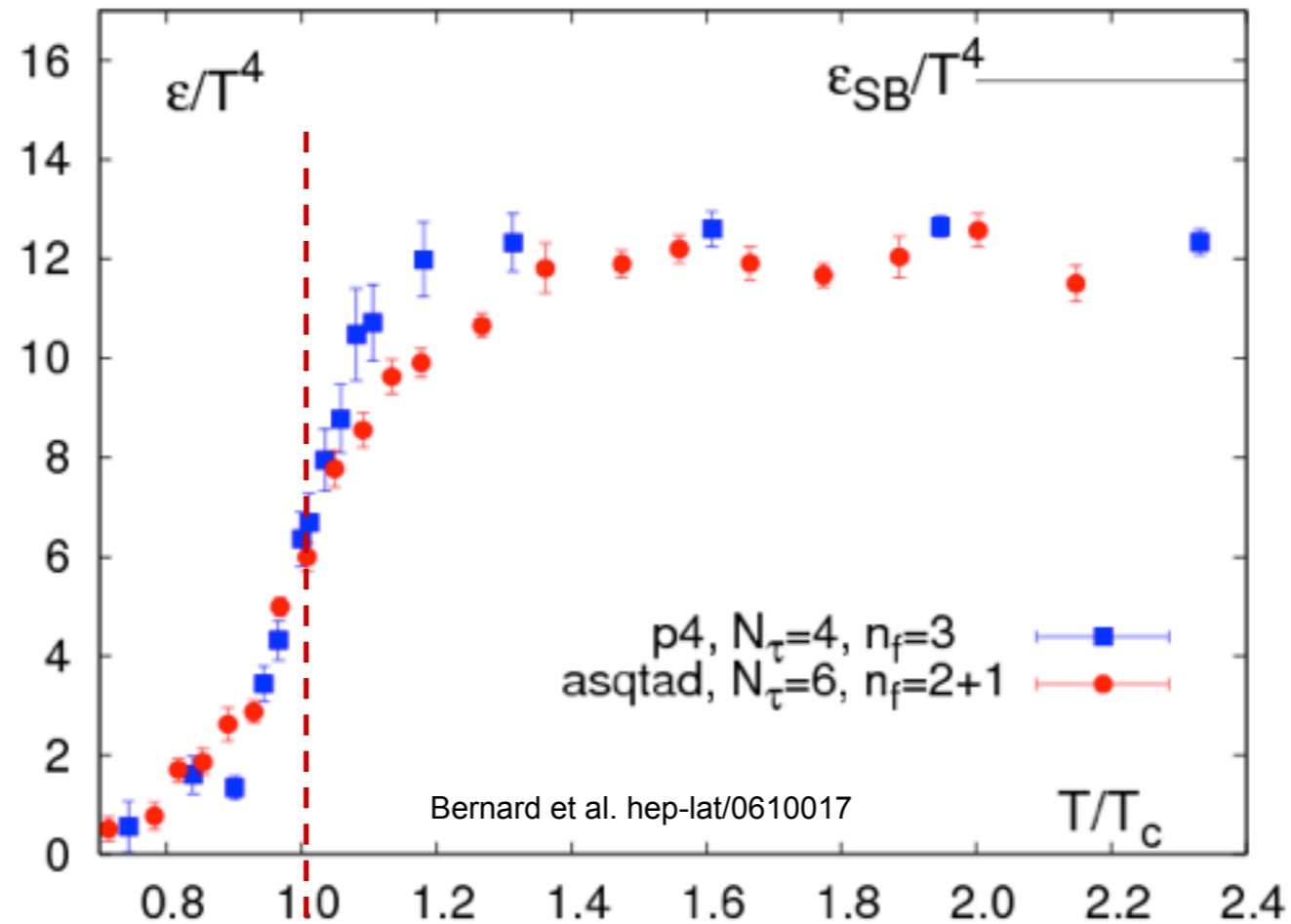
Calculable with pQCD

Small coupling
Quarks and gluons
are quasi-free

Two basic regimes in which QCD theory gives quantitative results:
Hard scattering and bulk matter

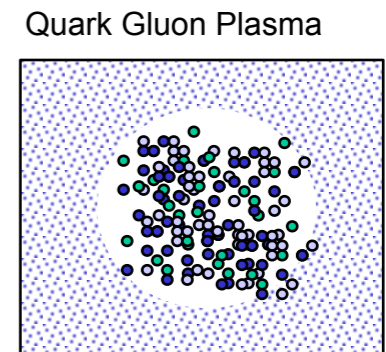
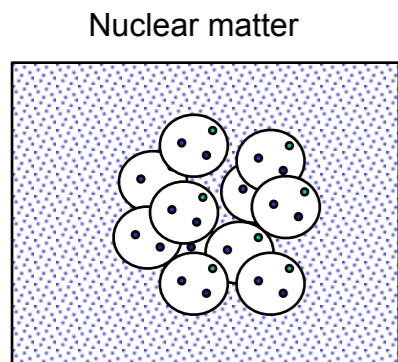
The Quark Gluon Plasma

Energy density from Lattice QCD



$$\epsilon \propto gT^4$$

g : deg of freedom



$$T_c \sim 170 - 190 \text{ MeV}$$

$$\epsilon_c \sim 1 \text{ GeV/fm}^3$$

Deconfinement transition: sharp rise of energy density at T_c
 Increase in degrees of freedom: hadrons (3 pions) \rightarrow quarks+gluons (37)

RHIC and LHC

RHIC, Brookhaven
 $\text{Au+Au } \sqrt{s_{\text{NN}}} = 200 \text{ GeV}$



First run: 2000

STAR, PHENIX,
PHOBOS, BRAHMS

LHC, Geneva
 $\text{Pb+Pb } \sqrt{s_{\text{NN}}} = 2760 \text{ GeV}$

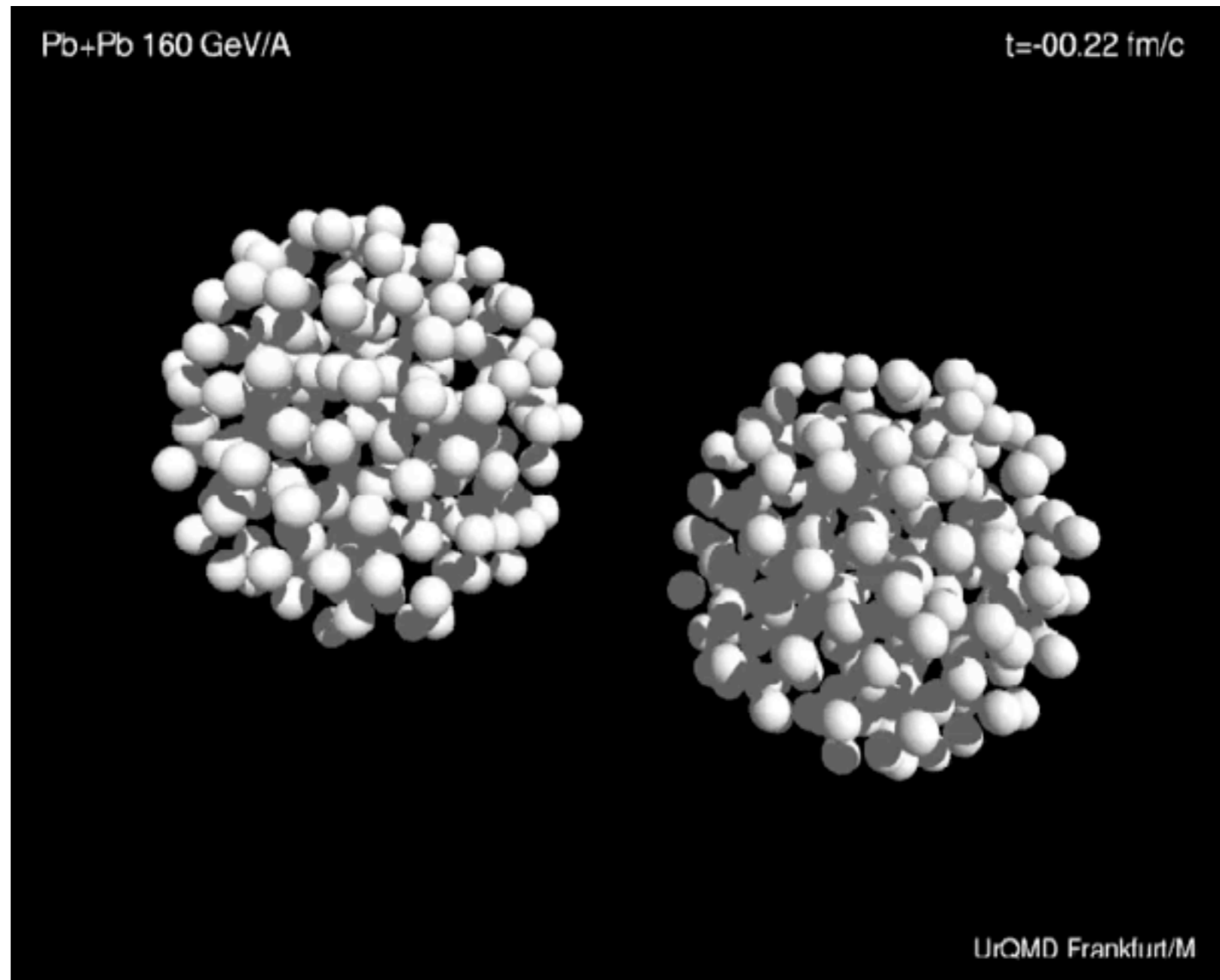


First run: 2009/2010

Currently under maintenance
Restart 2015 with higher energy:
 $\text{pp } \sqrt{s} = 13 \text{ TeV}$, $\text{PbPb } \sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$

ALICE, ATLAS,
CMS, (LHCb)

A nucleus-nucleus collision



Colored spheres: quarks

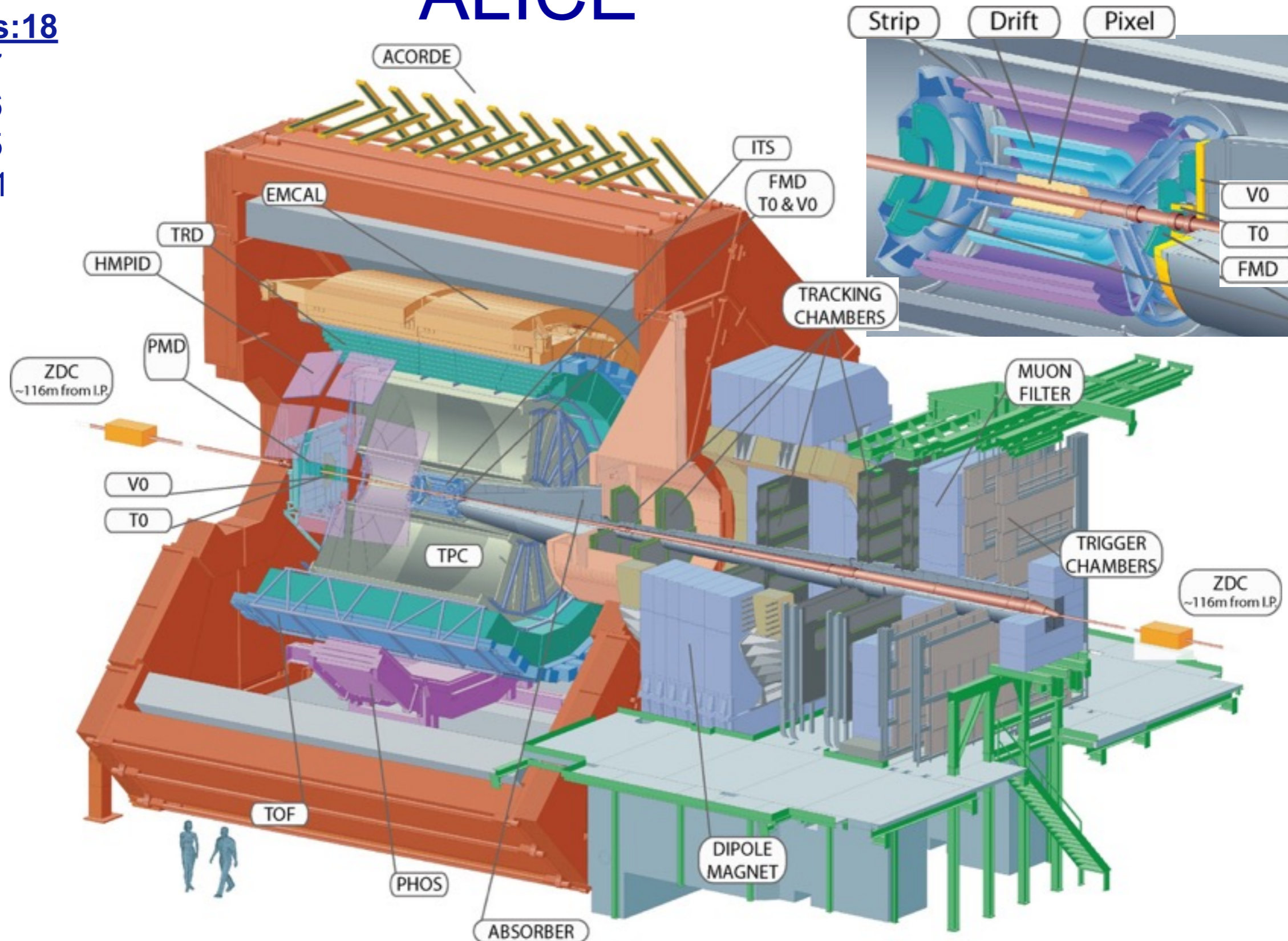
White spheres: hadrons, i.e. bound quarks

In a nuclear collision, a Quark-Gluon Plasma (liquid) is formed
⇒ Study this new state of matter

ALICE

Technologies:18

Tracking: 7
PID: 6
Calo.: 5
Trigger, N_{ch} :11



Size: 16 x 26 meters
Weight: 10,000 tons

Optimised for low momentum, high-multiplicity tracking and Particle Identification

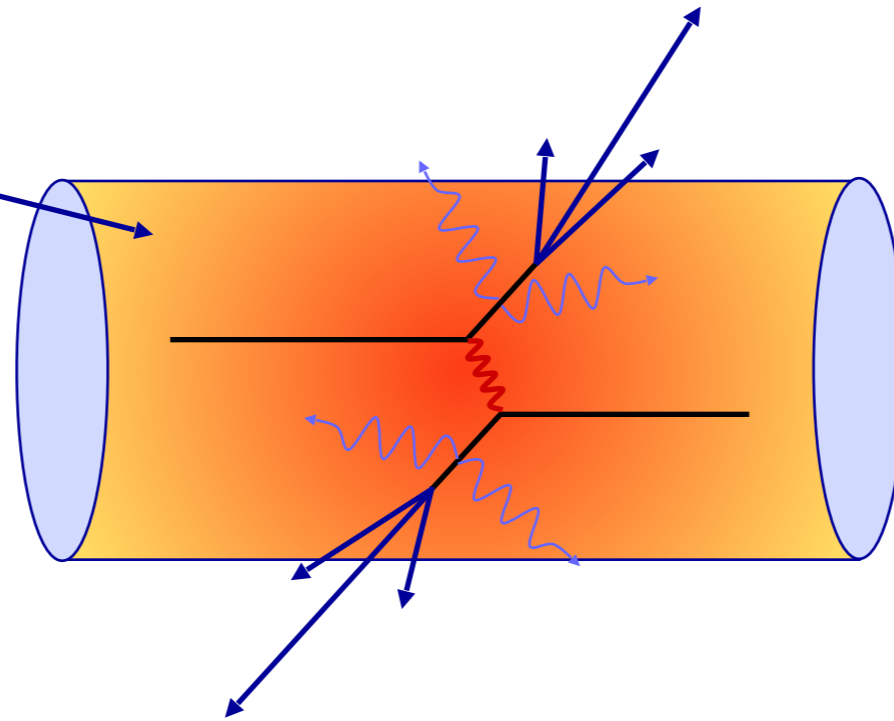
Heavy ion collisions

Heavy-ion collisions produce
'quasi-thermal' QCD matter

Dominated by soft partons
 $p \sim T \sim 100\text{-}300 \text{ MeV}$

'Bulk observables'

Study hadrons produced by the QGP
Typically $p_T < 1\text{-}2 \text{ GeV}$



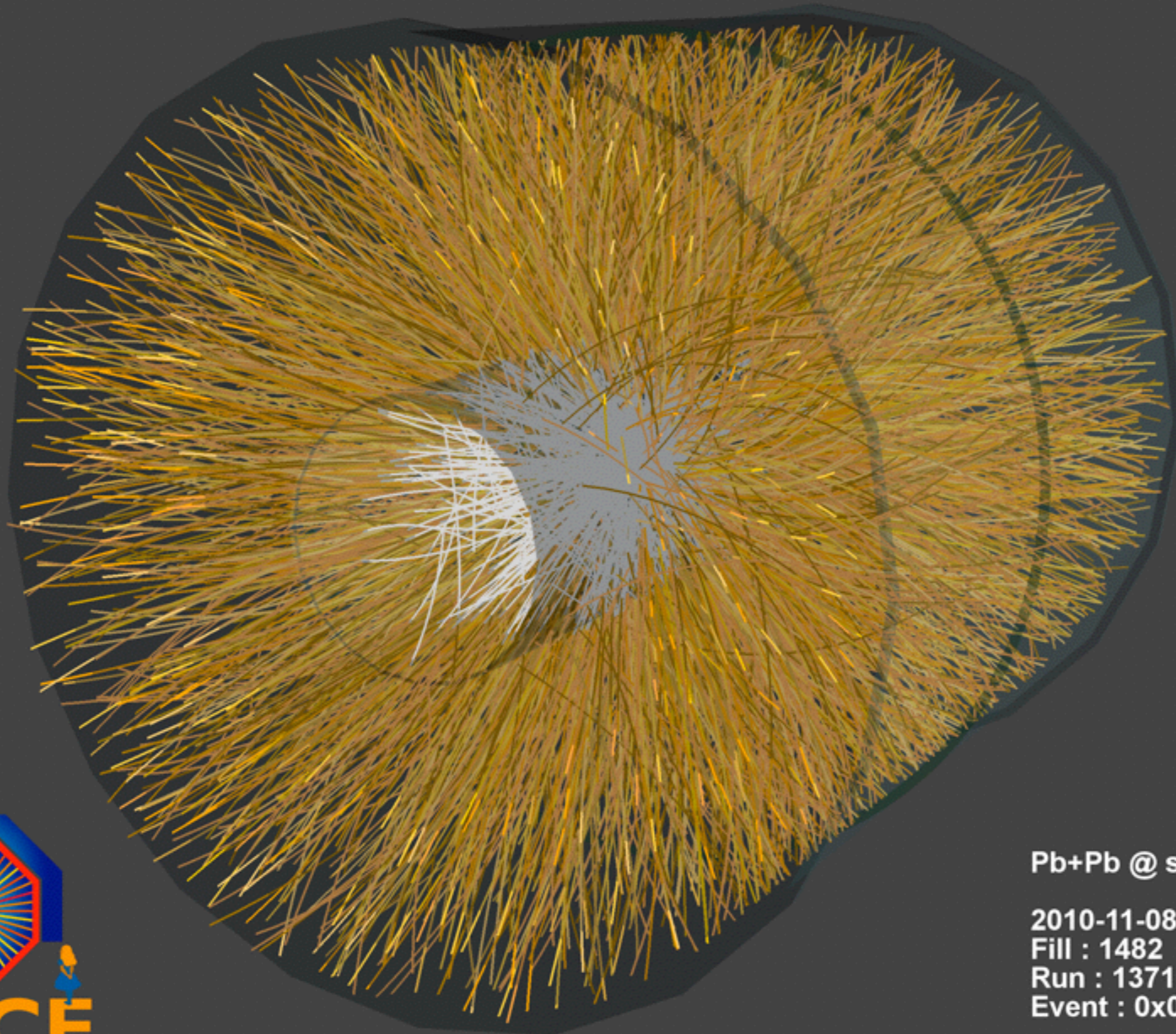
'Hard probes'

Hard-scatterings produce 'quasi-free' partons
 \Rightarrow Probe medium through energy loss
 $p_T > 5 \text{ GeV}$

Two basic approaches to learn about the QGP

- 1) Bulk observables
- 2) Hard probes

Part I: the bulk; QGP fragments



Pb+Pb @ $\sqrt{s} = 2.76$ ATeV

2010-11-08 11:30:46

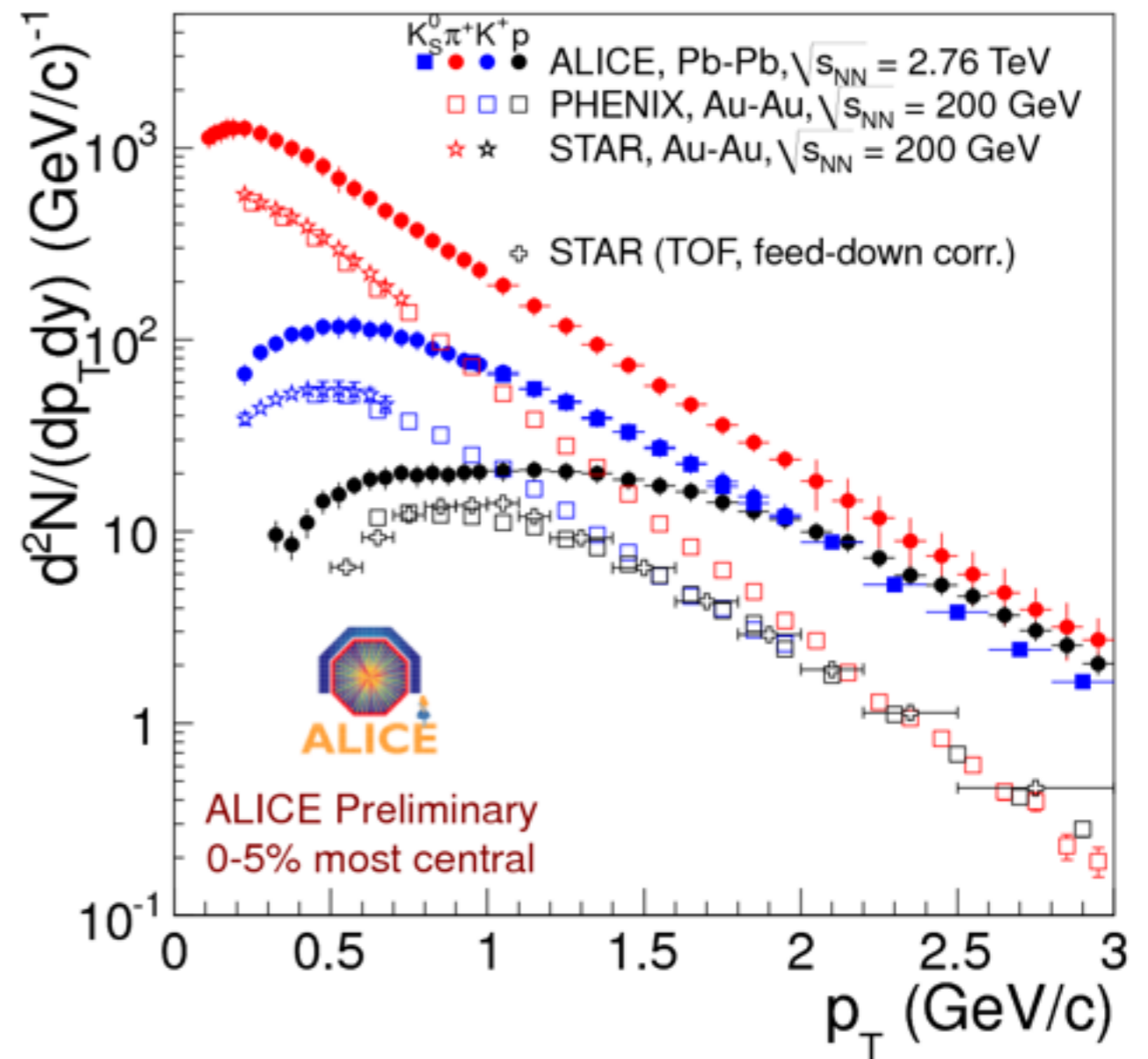
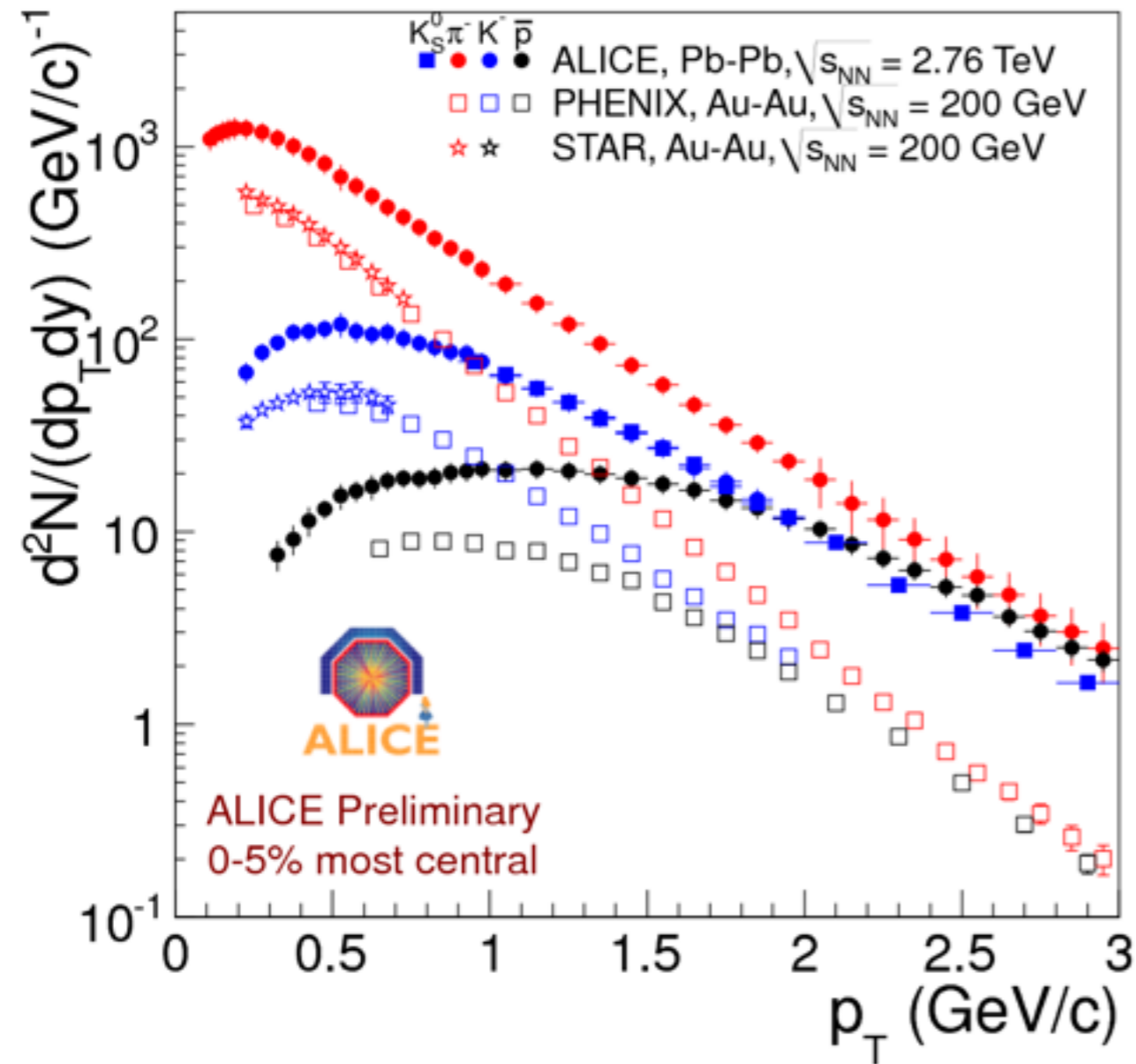
Fill : 1482

Run : 137124

Event : 0x00000000D3BBE693

p_T -spectra – radial flow

PRL109, 252301, PRC, arXiv:1303.0737



Large increase in mean p_T from RHIC ($\sqrt{s_{NN}}=200$ GeV) to LHC

First indication of collective behaviour; pressure

Mass dependence: same Lorentz boost (β) gives larger momentum for heavier particles
 $(m_p > m_K > m_\pi)$

Hydrodynamics

PLB, arXiv:1401.1250

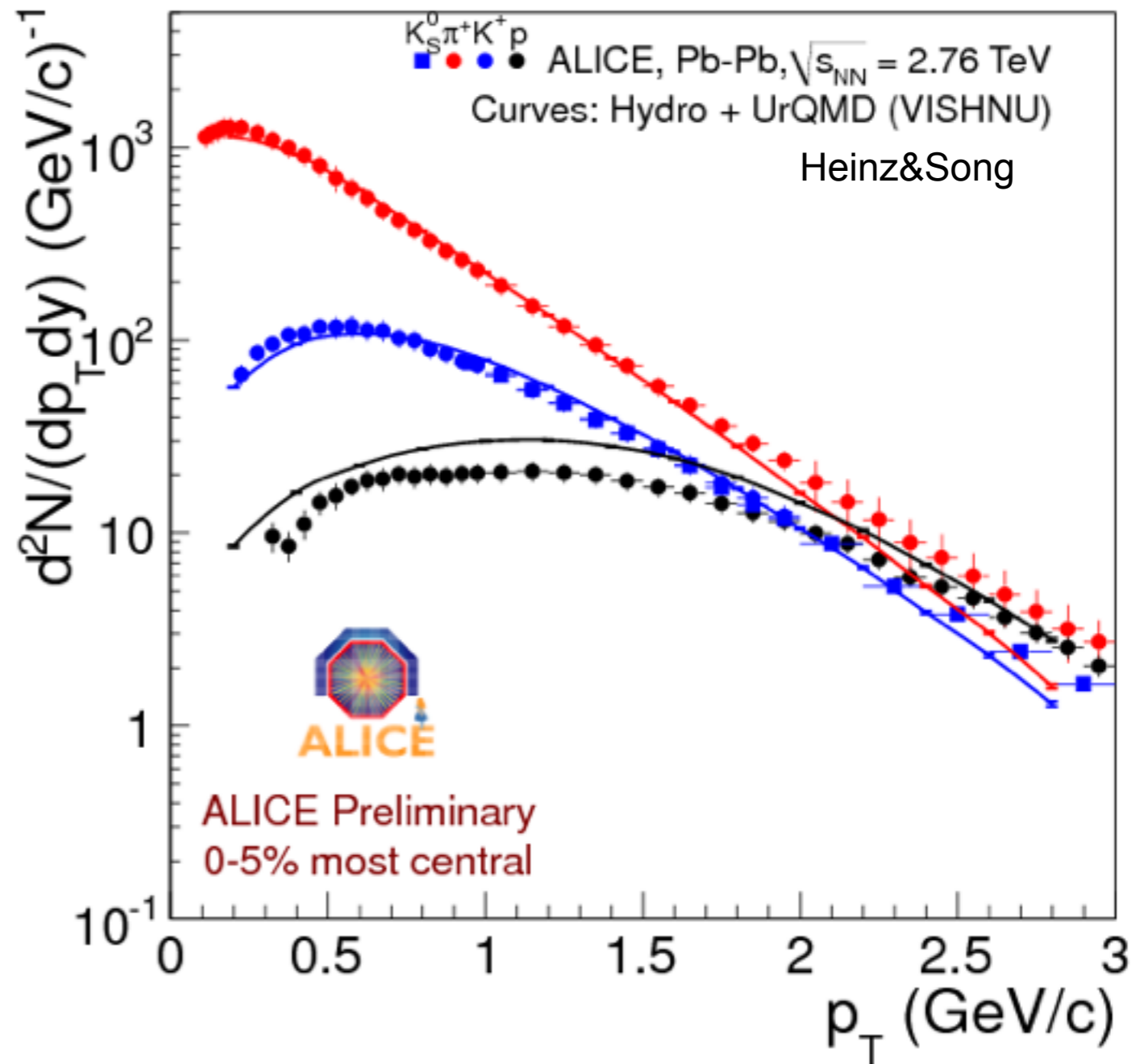
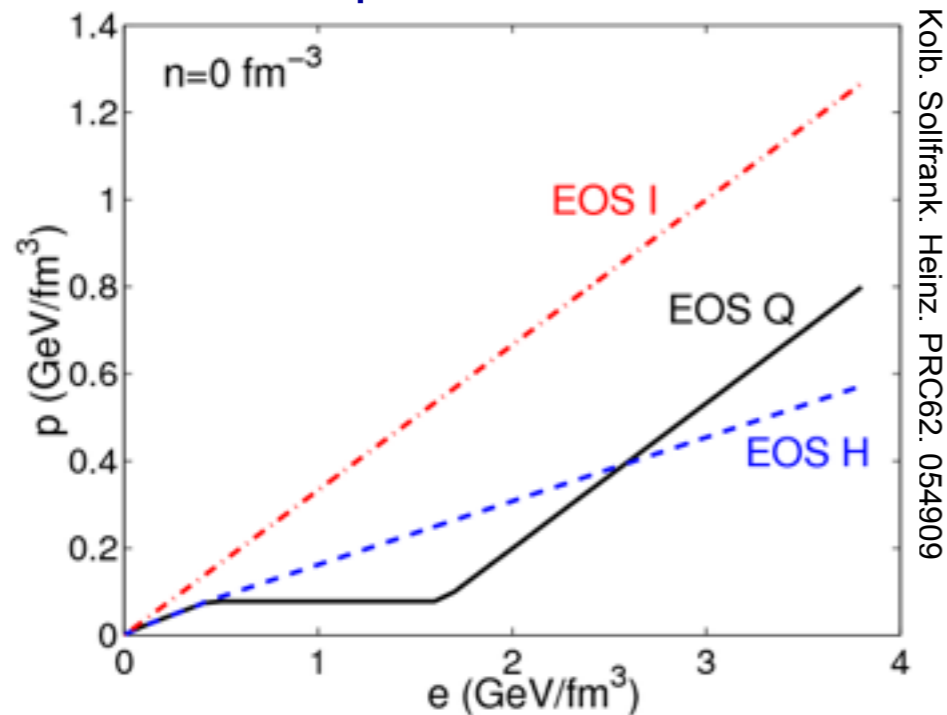
Energy-momentum conservation

$$\partial_{\mu} T^{\mu\nu} = 0$$

Continuity equations for conserved currents

$$\partial_{\mu} J_i^{\mu} = 0$$

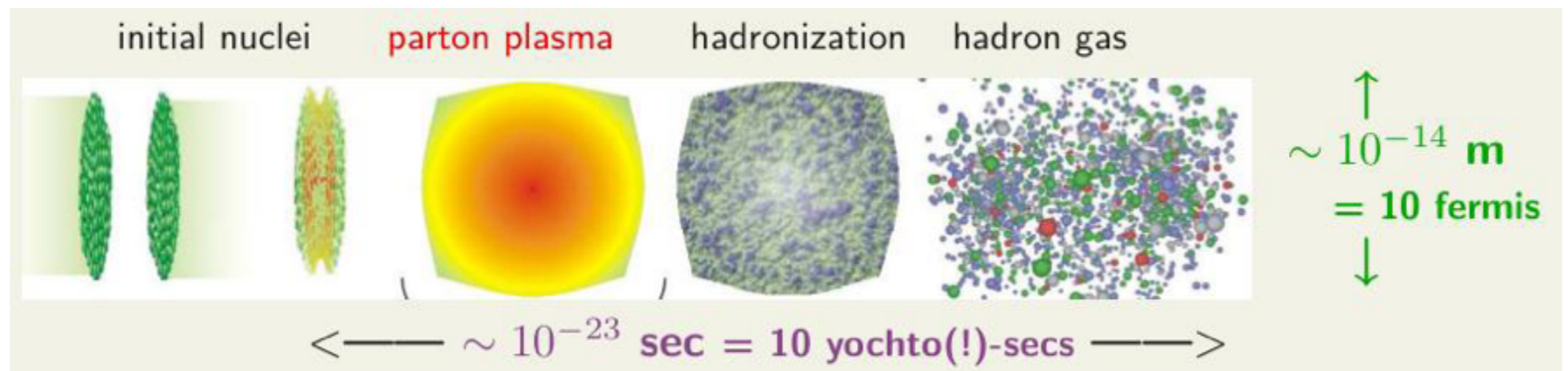
Equation of state



Measured spectra shapes agree with hydrodynamical calculation
 \Rightarrow It really looks like a fluid

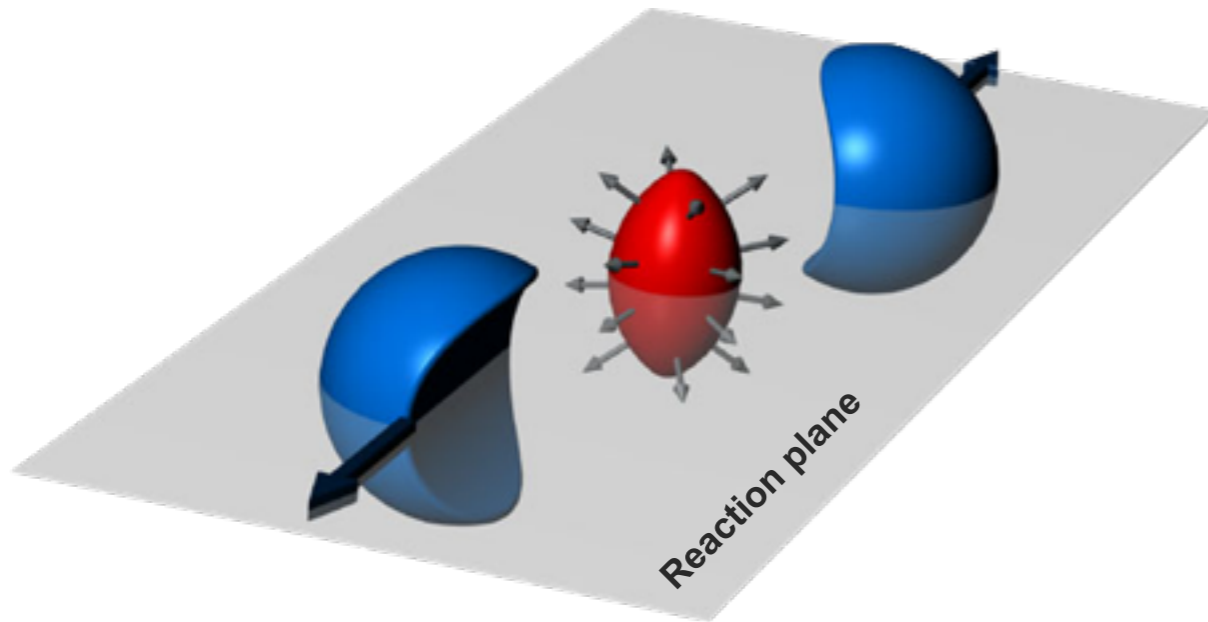
Time evolution

All observables integrate over evolution

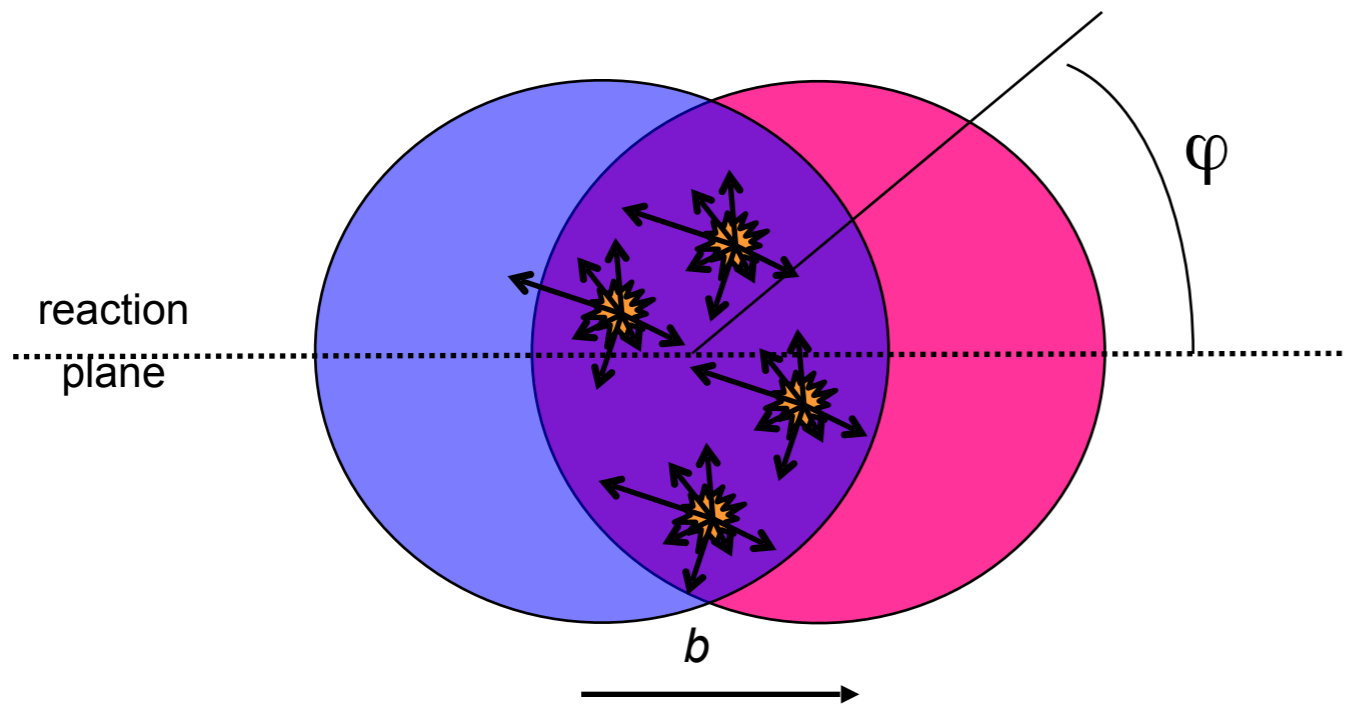


Radial flow integrates over entire 'push'

Elliptic flow

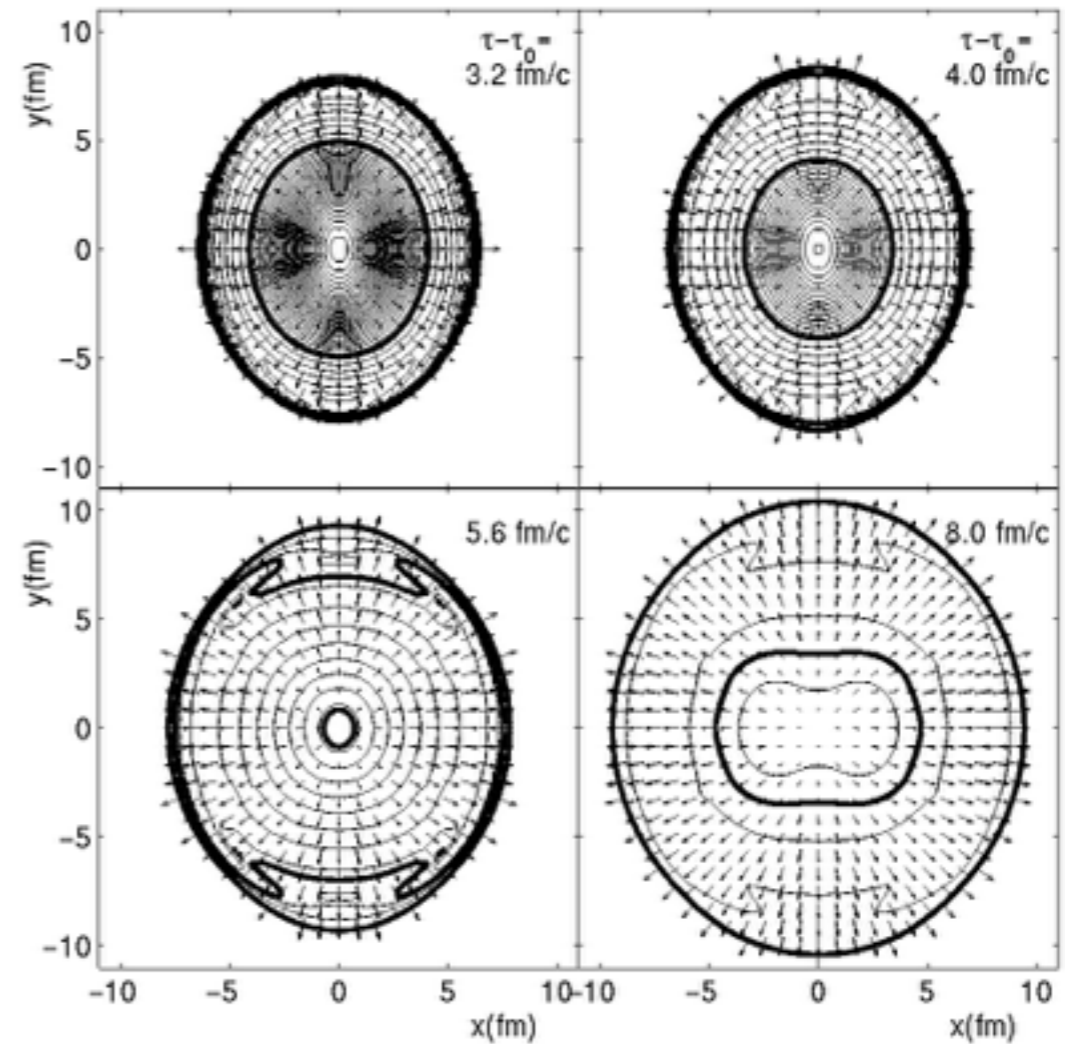


Elliptic flow:
Yield modulation in-out reaction plane



$$\frac{dN}{d\varphi} = N(1 + 2v_2 \cos 2\varphi)$$

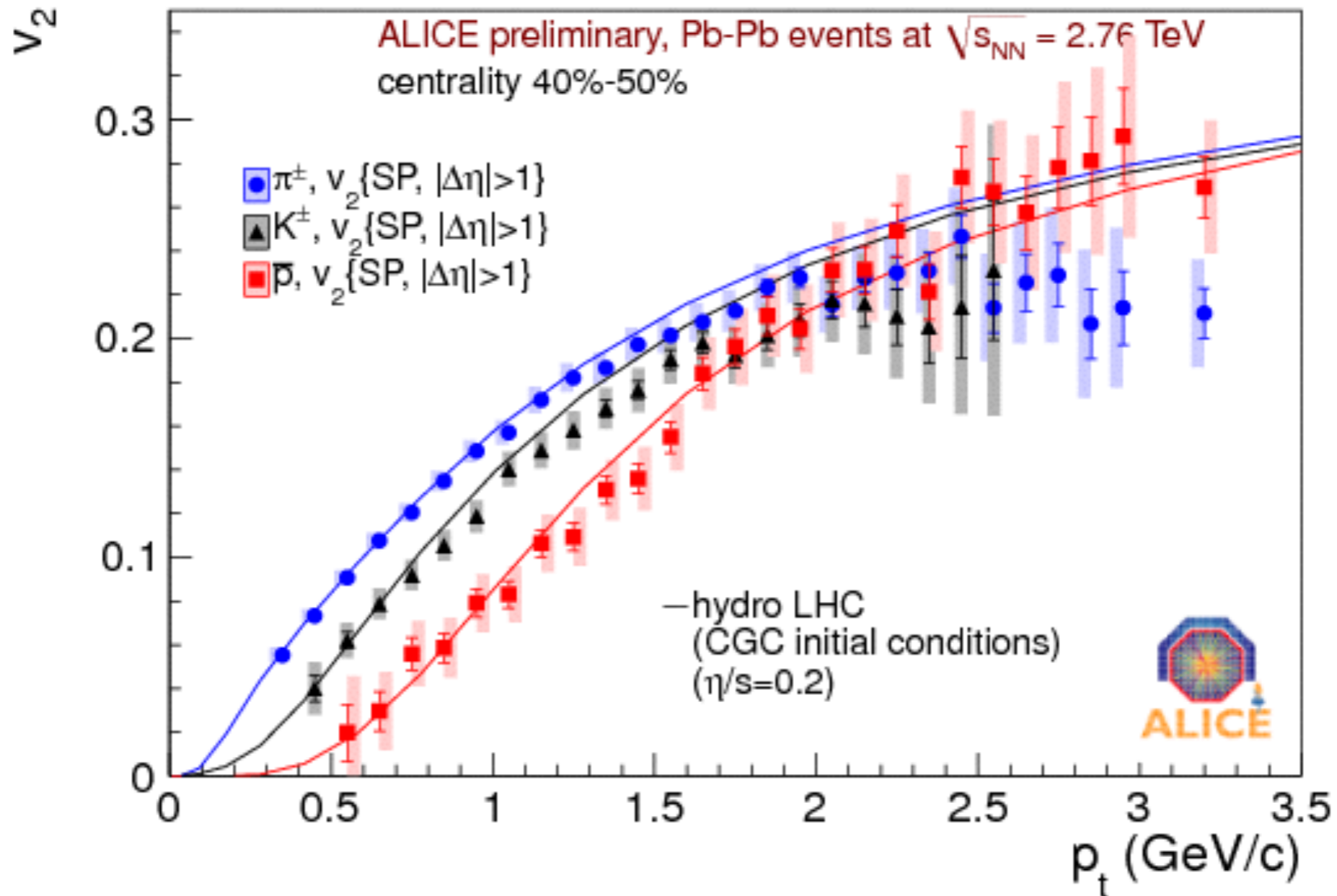
Hydrodynamical calculation



Anisotropy reduces during evolution
 v_2 more sensitive to early times

Elliptic flow

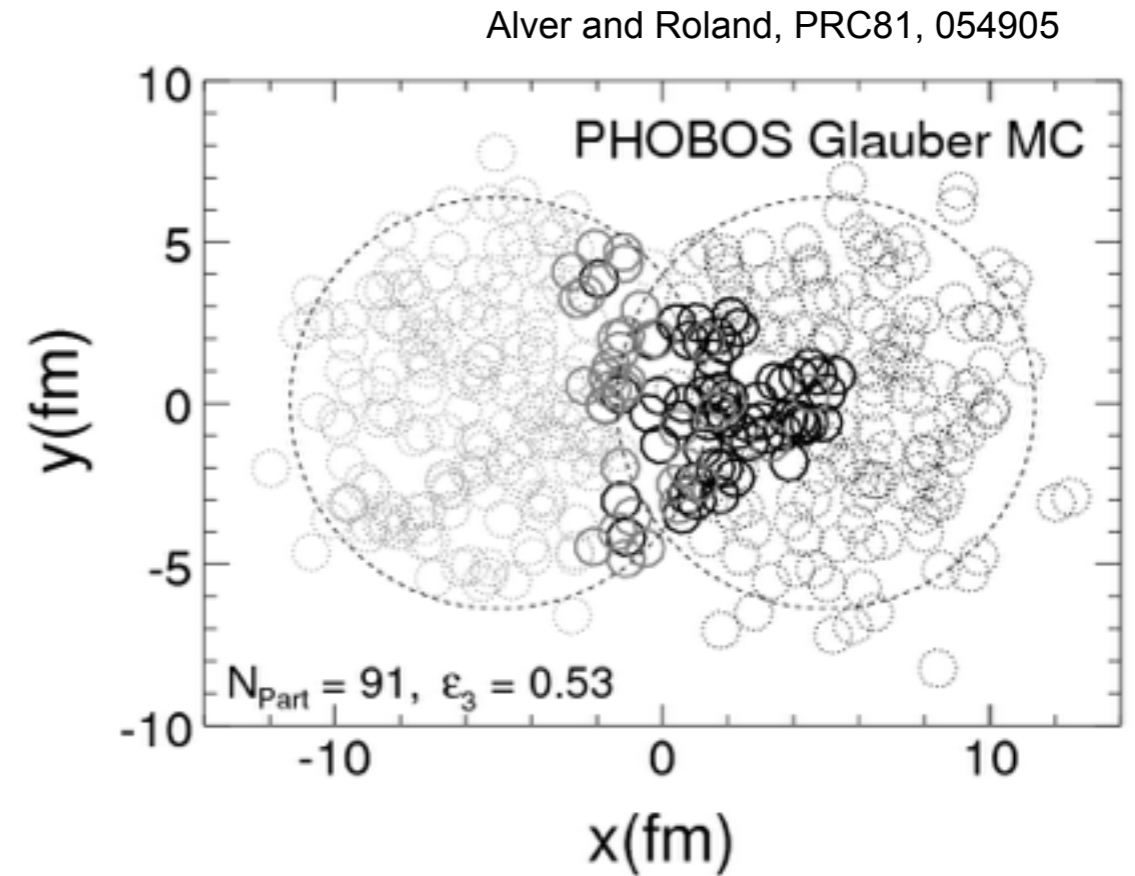
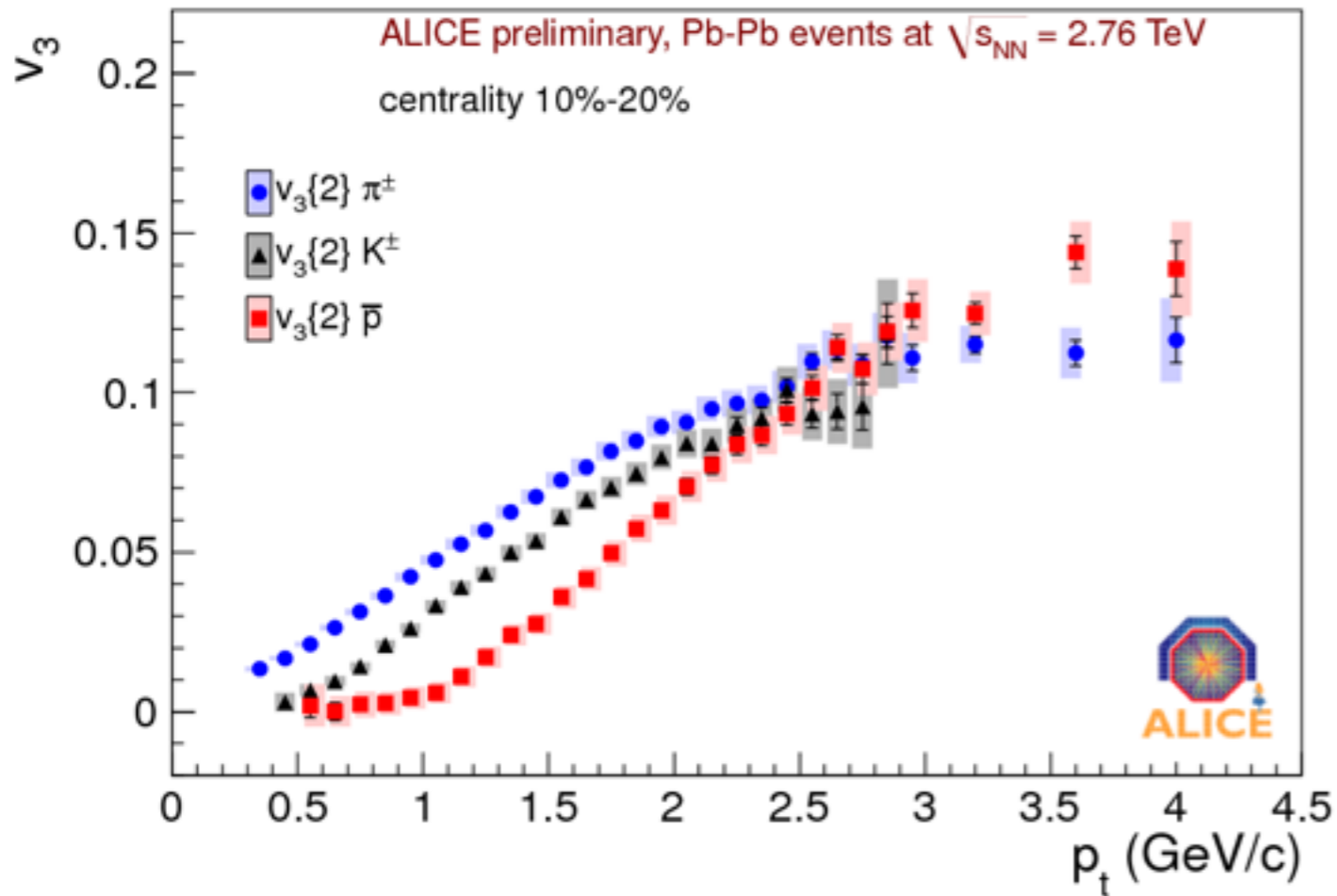
$$\frac{dN}{d\varphi} = N(1 + 2v_2 \cos 2\varphi)$$



Mass-dependence of v_2 measures flow velocity

Good agreement between data and hydro

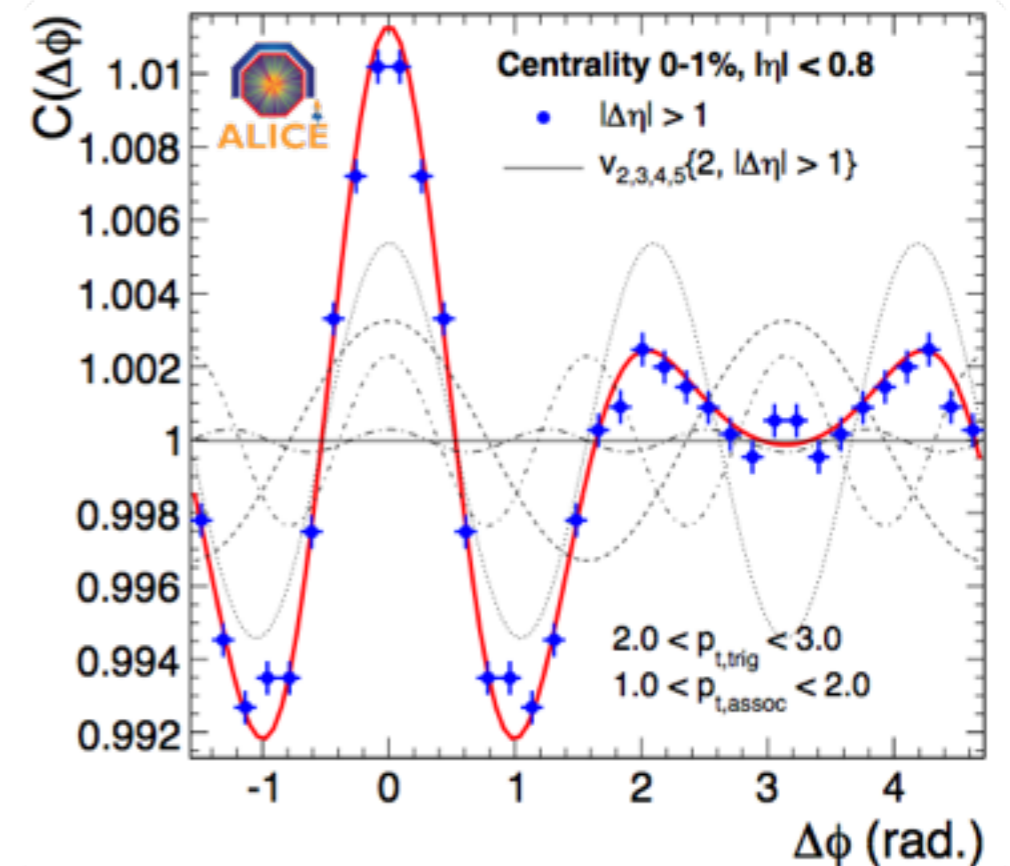
Higher harmonics



3rd harmonic 'triangularity' v_3 is large
(in central events)

Mass ordering also seen for v_3
indicates collective flow

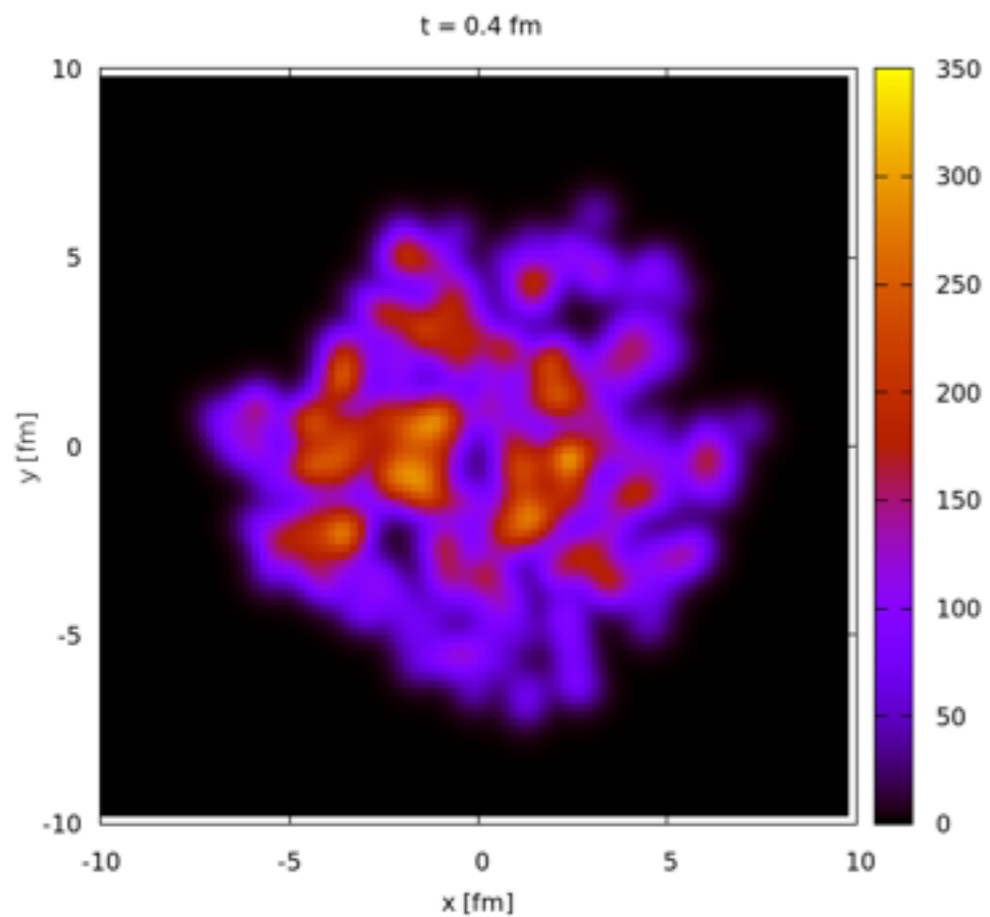
Dominant effect in azimuthal correlations
at $p_T = 1-3$ GeV



Higher harmonics

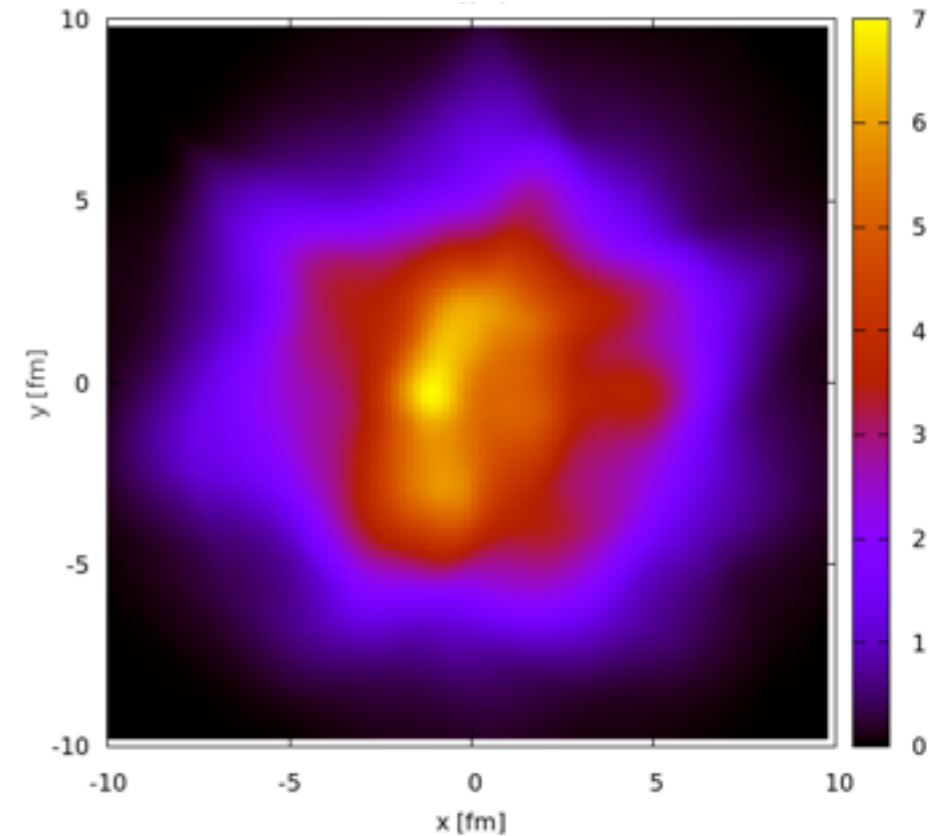
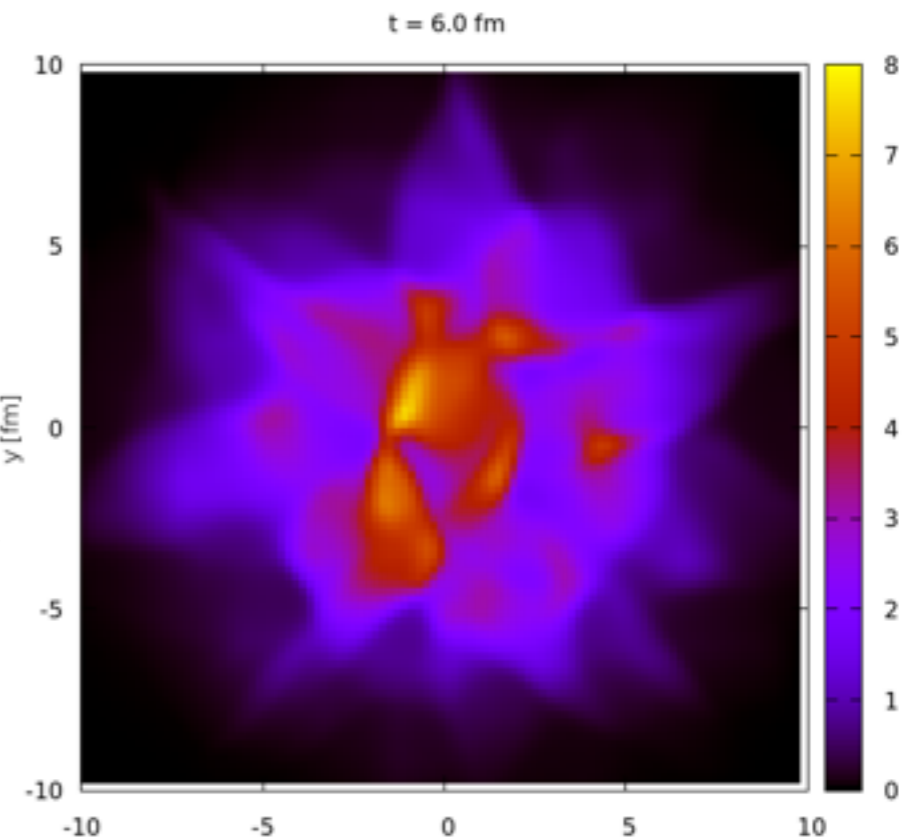
Schenke and Jeon, Phys.Rev.Lett.106:042301

In general: initial state may be 'lumpy'
(not a smooth ellipse)



$$\eta/s = 0$$

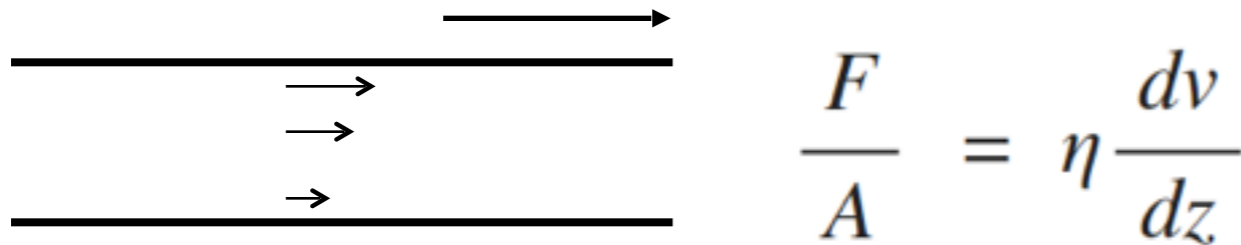
$$\eta/s = 0.16$$



How much of this is visible in the final state,
depends on shear viscosity η

Viscosity

Viscous liquids dissipate energy



For a dilute gas:

$$\eta = \frac{1}{3} np\lambda \quad \lambda = \frac{1}{n\sigma}$$

$$\eta \propto T^{\frac{1}{2}}$$

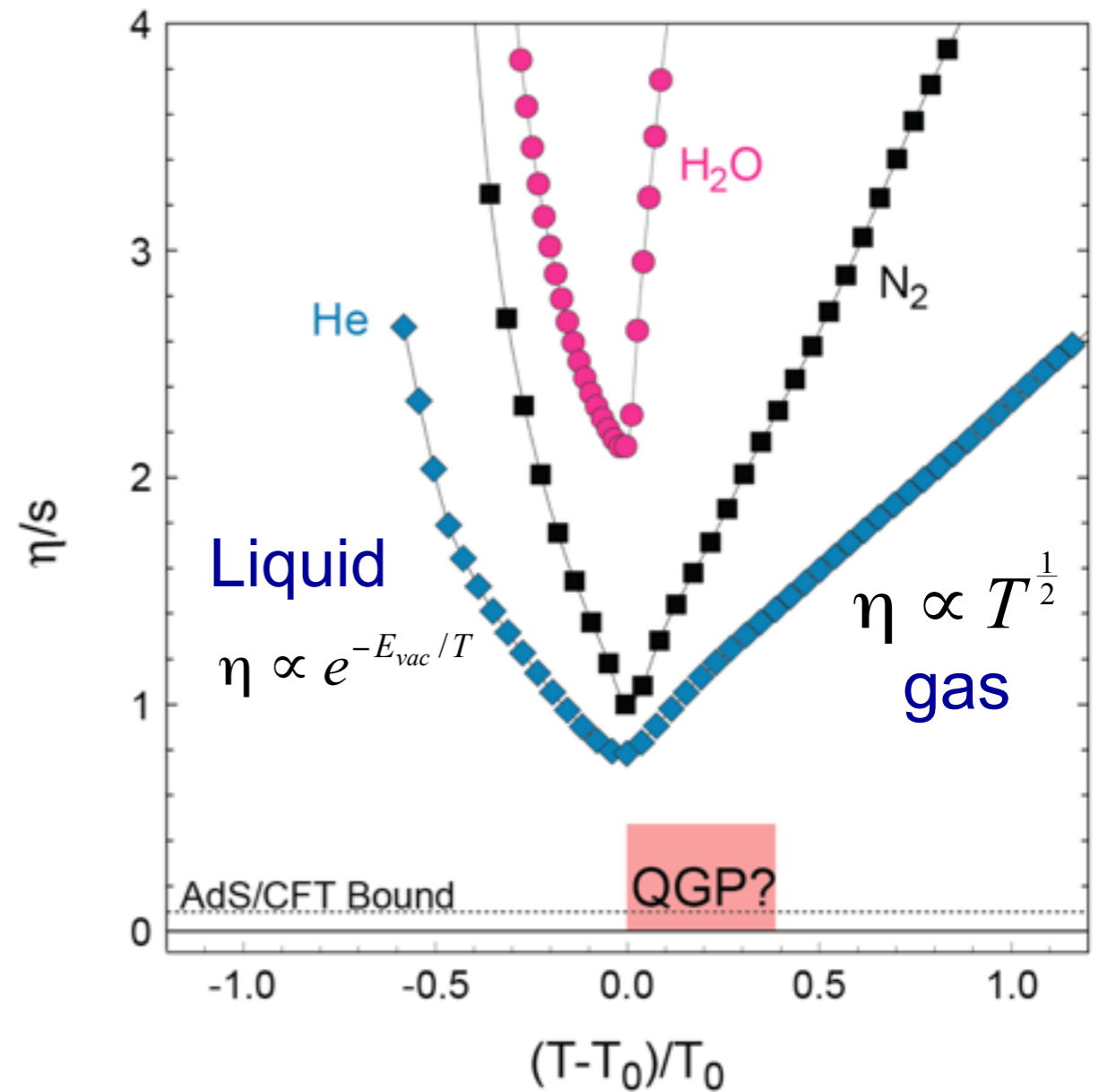
η increases with T

Liquid, densely packed, so:

$$\eta \propto e^{-E_{vac}/T}$$

E_{vac} : activation energy for jumps of vacancies

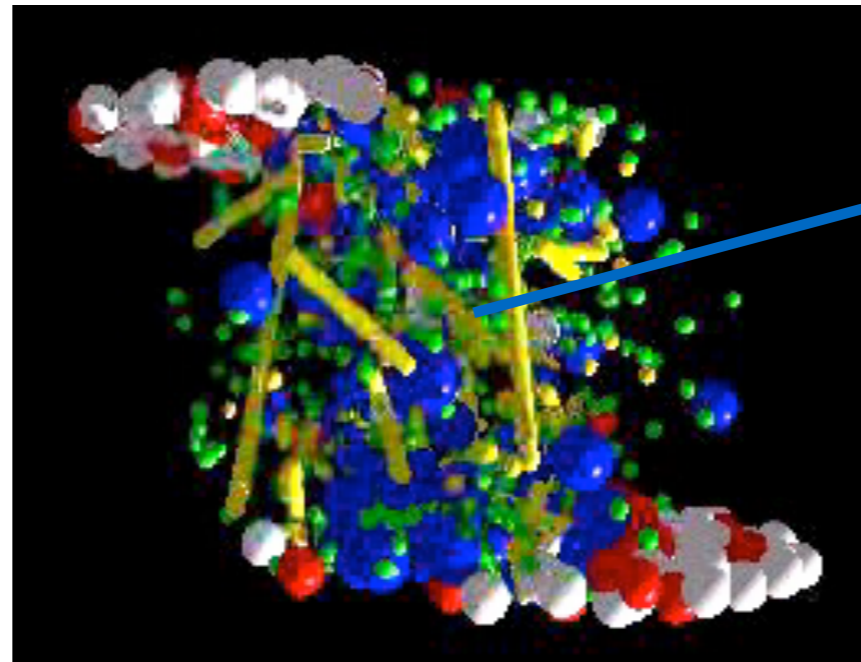
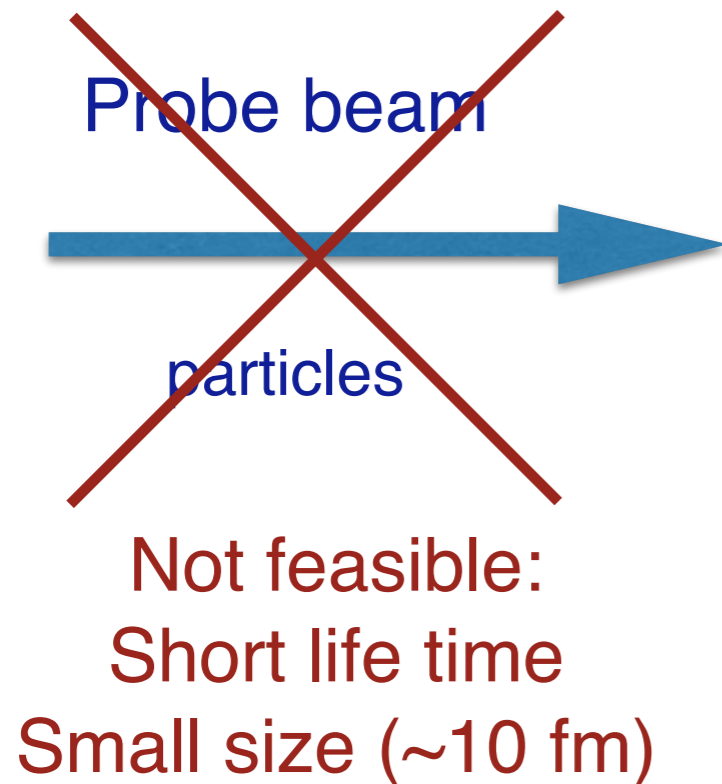
η decreases with T



Viscosity minimal at liquid-gas transition

QGP viscosity lower than any atomic matter

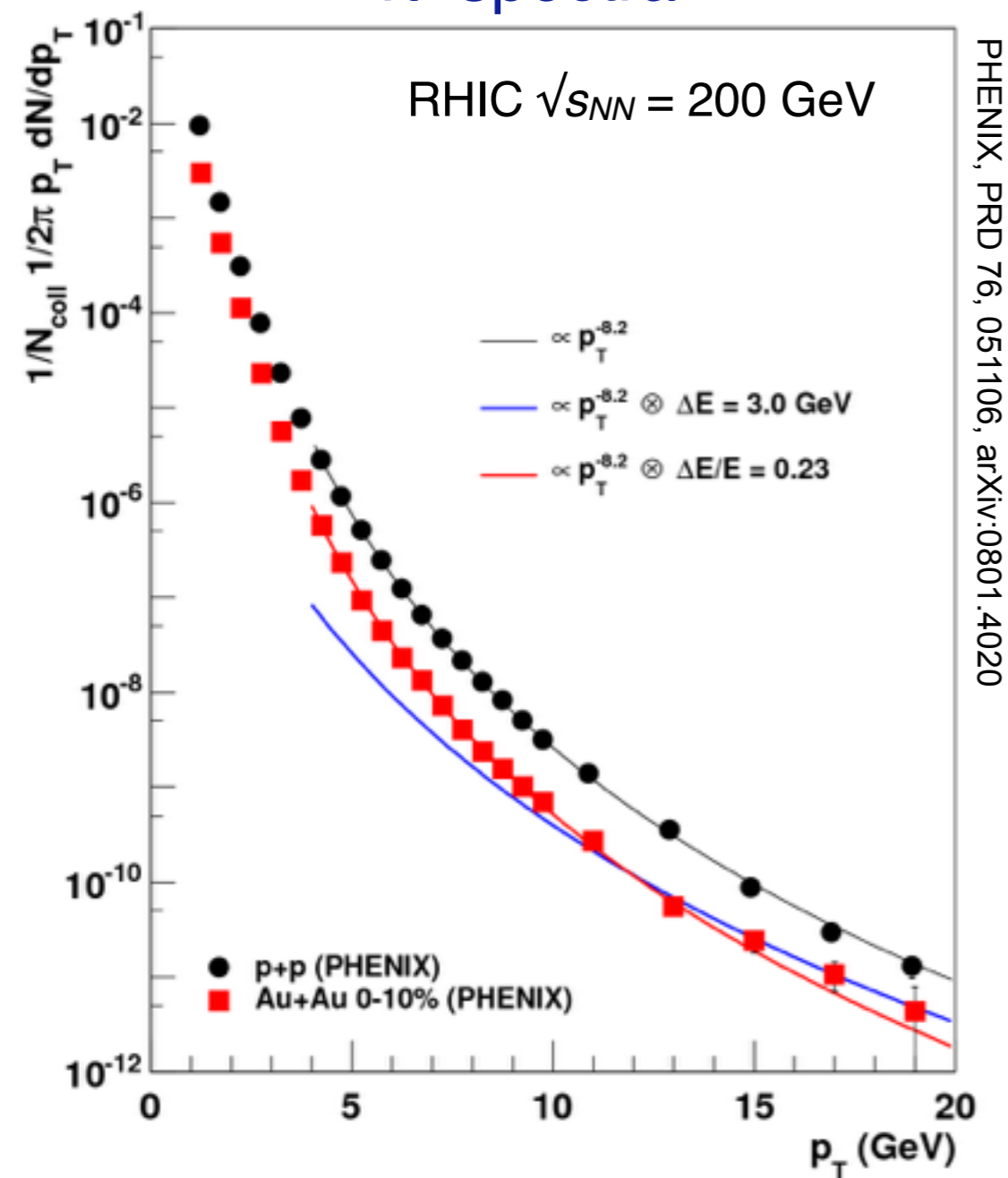
Probing the Quark-Gluon Plasma



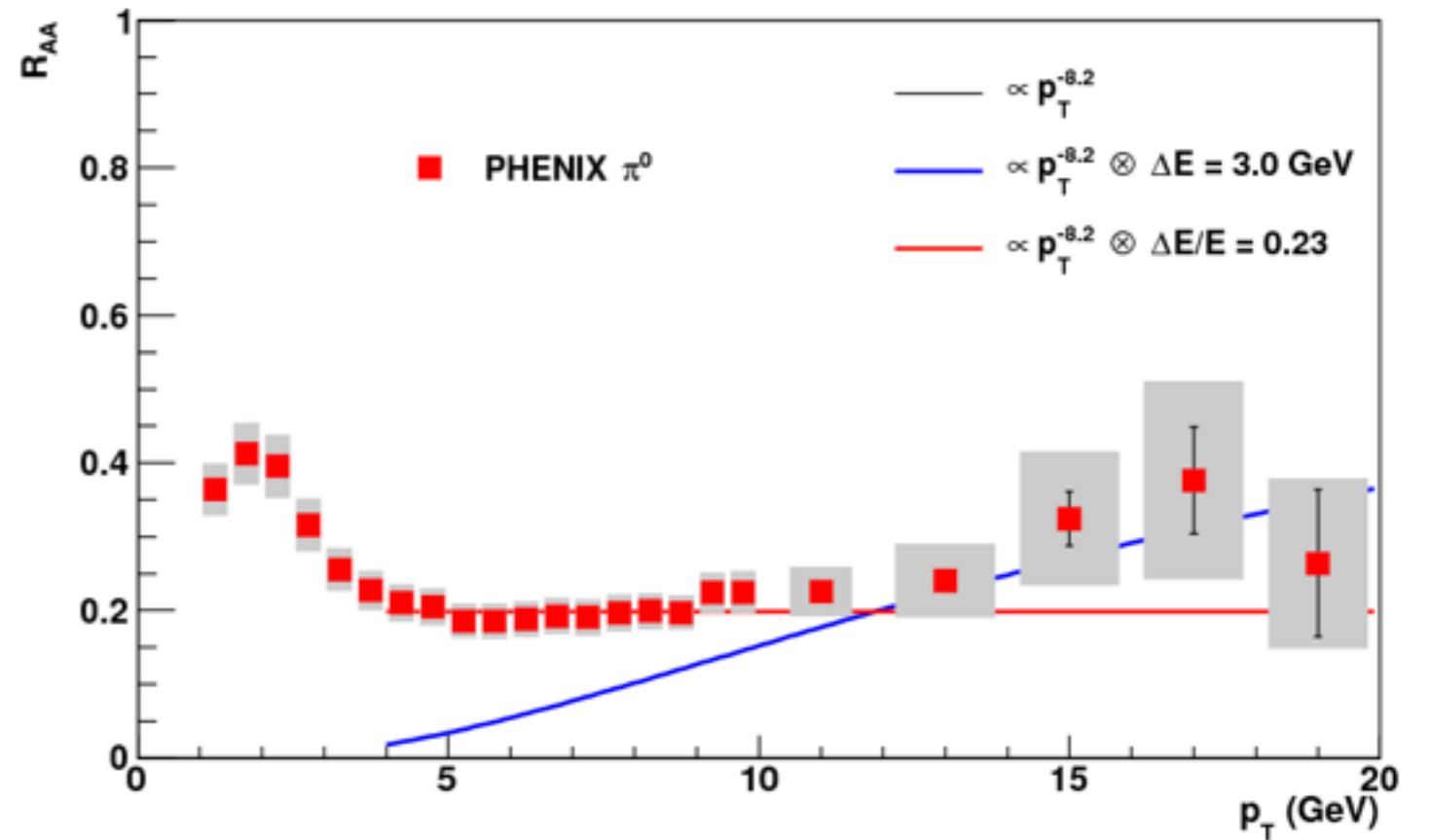
Use self-generated probe:
quarks, gluons from hard scattering
large transverse momentum

Nuclear modification factor at RHIC

π^0 spectra



Nuclear modification factor



Oversimplified calculation:
 -Fit pp with power law
 -Apply energy shift or relative E loss
Not even a model !

Ball-park numbers: $\Delta E/E \approx 0.2$, or $\Delta E \approx 3$ GeV
 for central collisions at RHIC

From RHIC to LHC

RHIC: 200 GeV
LHC: 2.76 TeV

per nucleon pair

LHC: spectrum less steep,
larger p_T reach

$$\frac{1}{2\pi p_T} \frac{dN}{dp_T} \propto p_T^{-n}$$

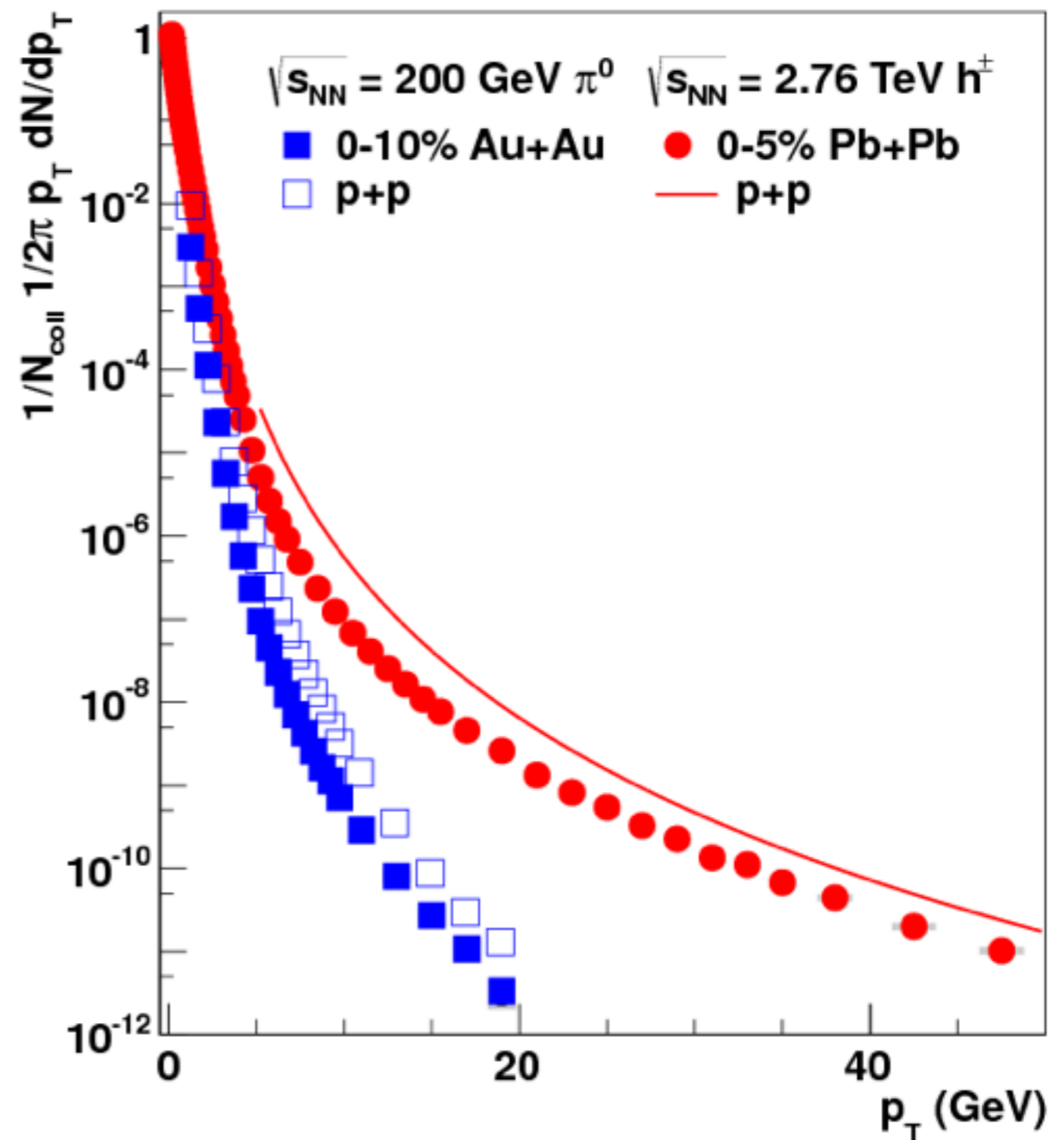
RHIC: $n \sim 8.2$

LHC: $n \sim 6.4$

Fractional energy loss:

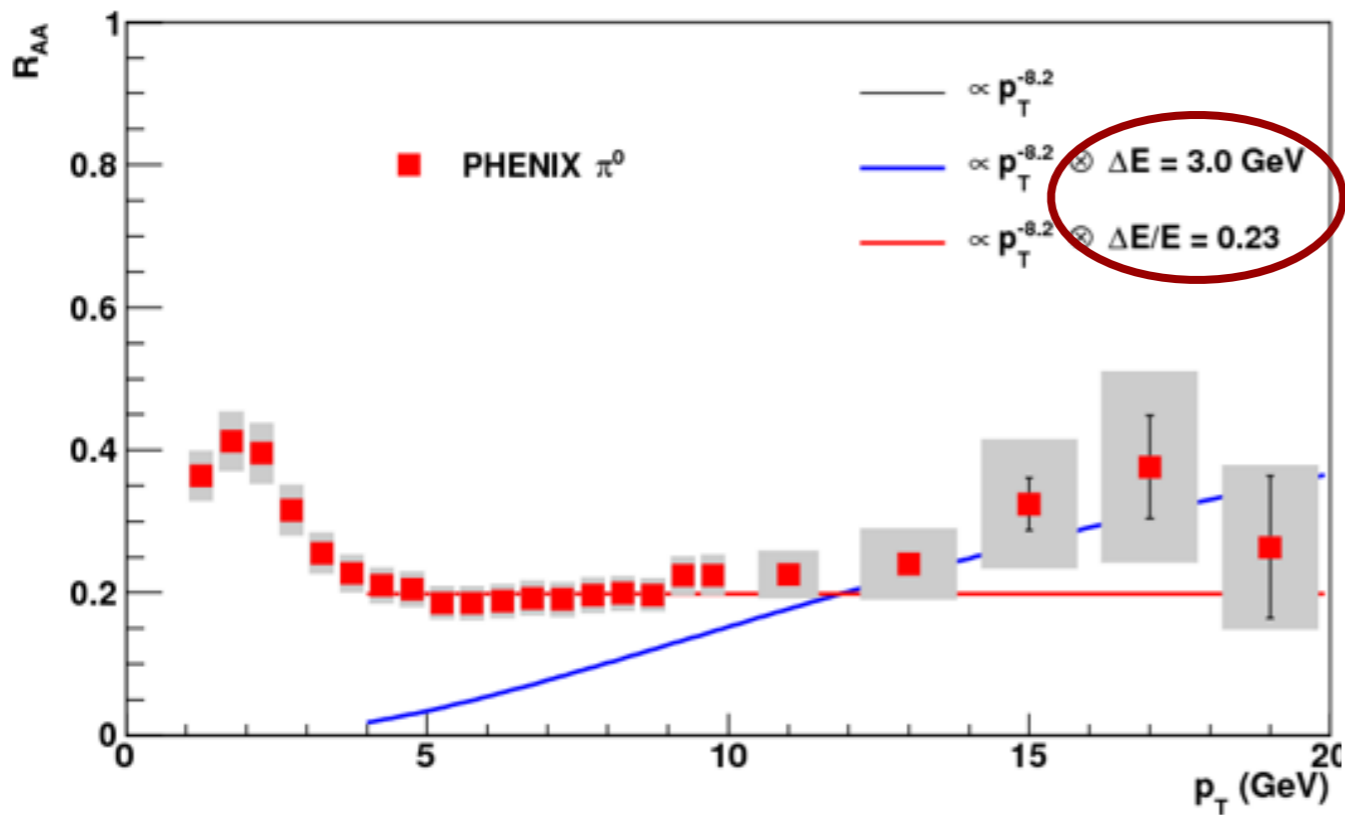
$$R_{AA} = \left(1 - \frac{\Delta E}{E}\right)^{n-2}$$

R_{AA} depends on n , steeper spectra, smaller R_{AA}



From RHIC to LHC

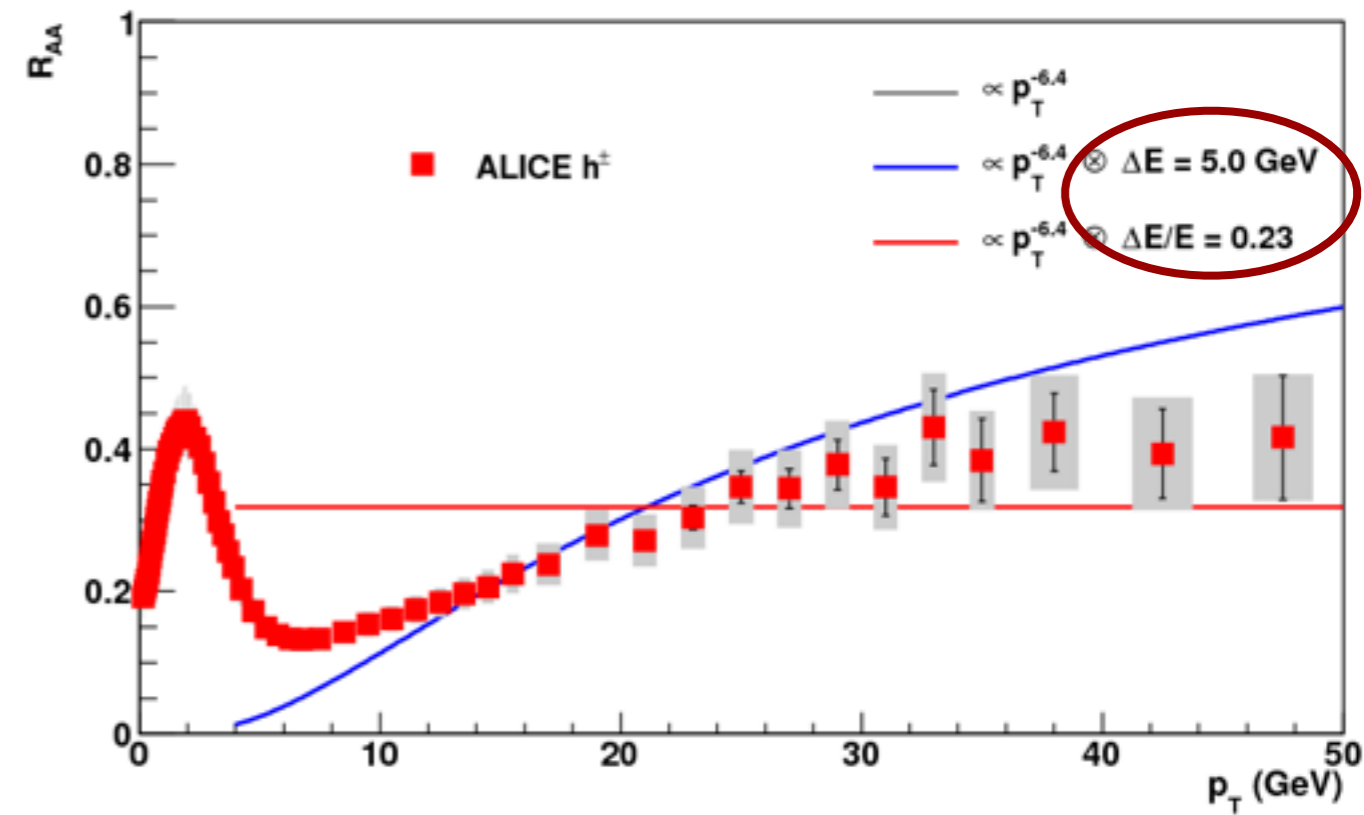
RHIC



RHIC: $n \sim 8.2$

$$(1 - 0.23)^{6.2} = 0.20$$

LHC



LHC: $n \sim 6.4$

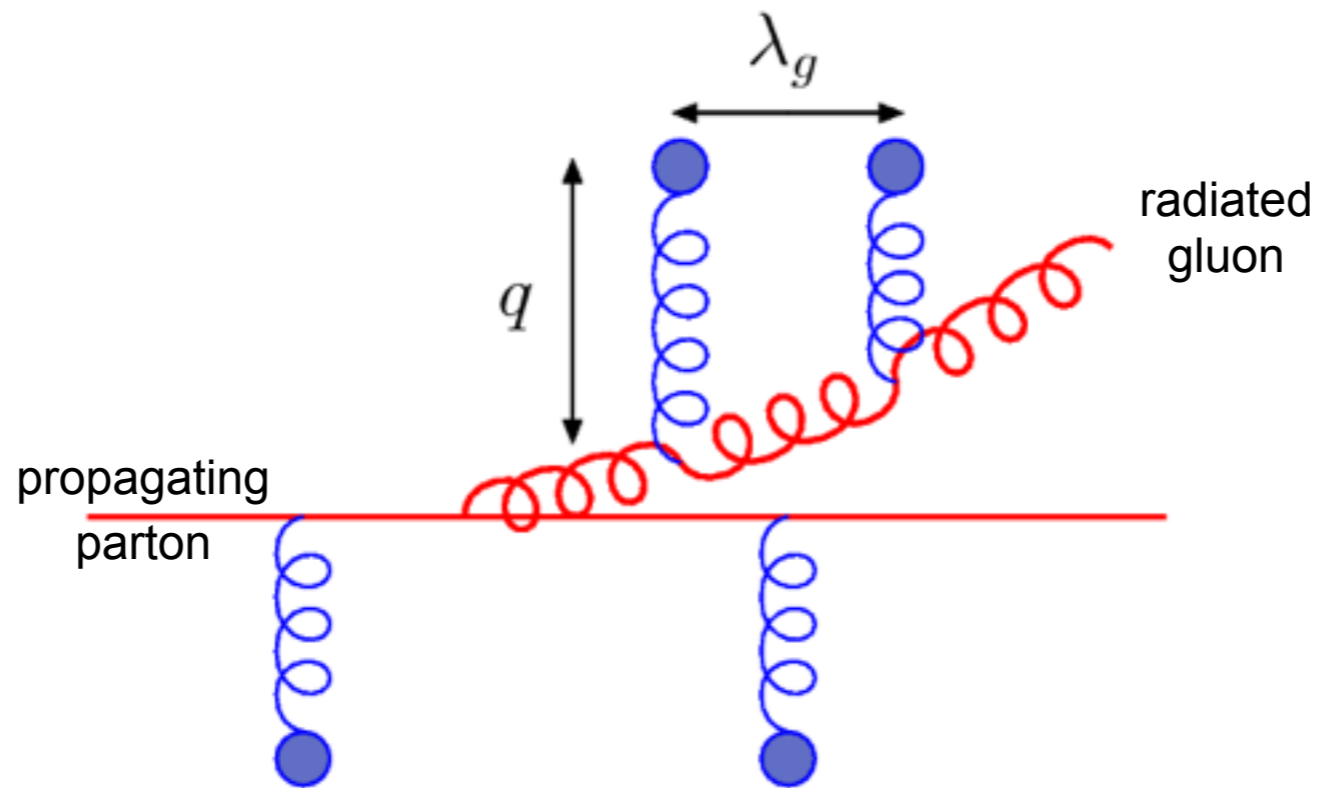
$$(1 - 0.23)^{4.4} = 0.32$$

Energy loss at LHC is larger than at RHIC
 R_{AA} is similar due to flatter p_T dependence

Towards a more complete picture

- Geometry: couple energy loss model to model of evolution of the density (hydrodynamics)
- Energy loss not single-valued, but a distribution
- Energy loss is partonic, not hadronic
 - Full modeling: medium modified shower
 - Simple ansatz for leading hadrons: energy loss followed by fragmentation
 - Quark/gluon differences

Medium-induced radiation



Key parameter:

Transport coefficient

$$\hat{q} \equiv \frac{\langle q_{\perp}^2 \rangle}{\lambda}$$

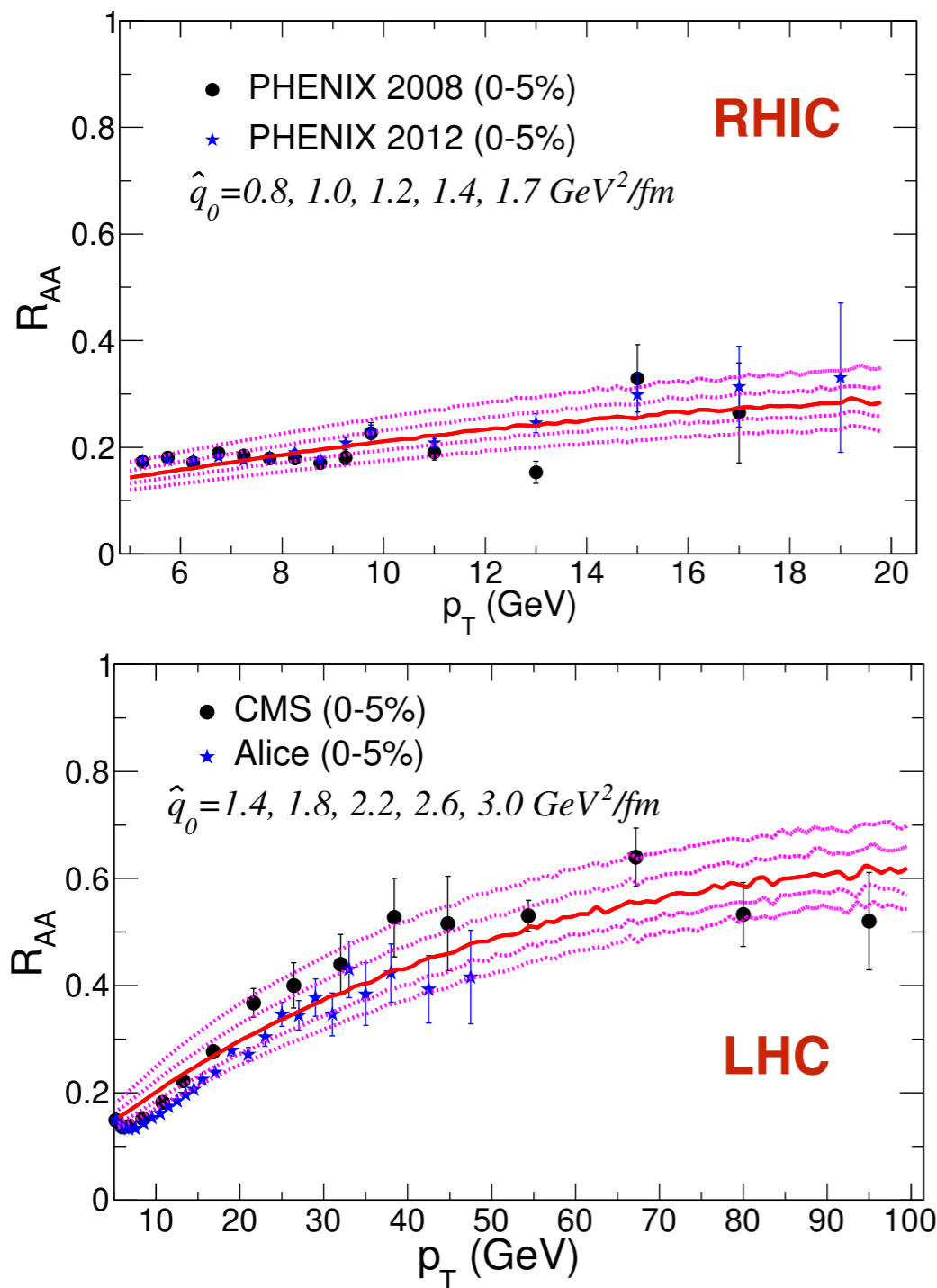
Mean transverse kick per unit path length

$$\Delta E_{med} \sim \alpha_s \hat{q} L^2$$

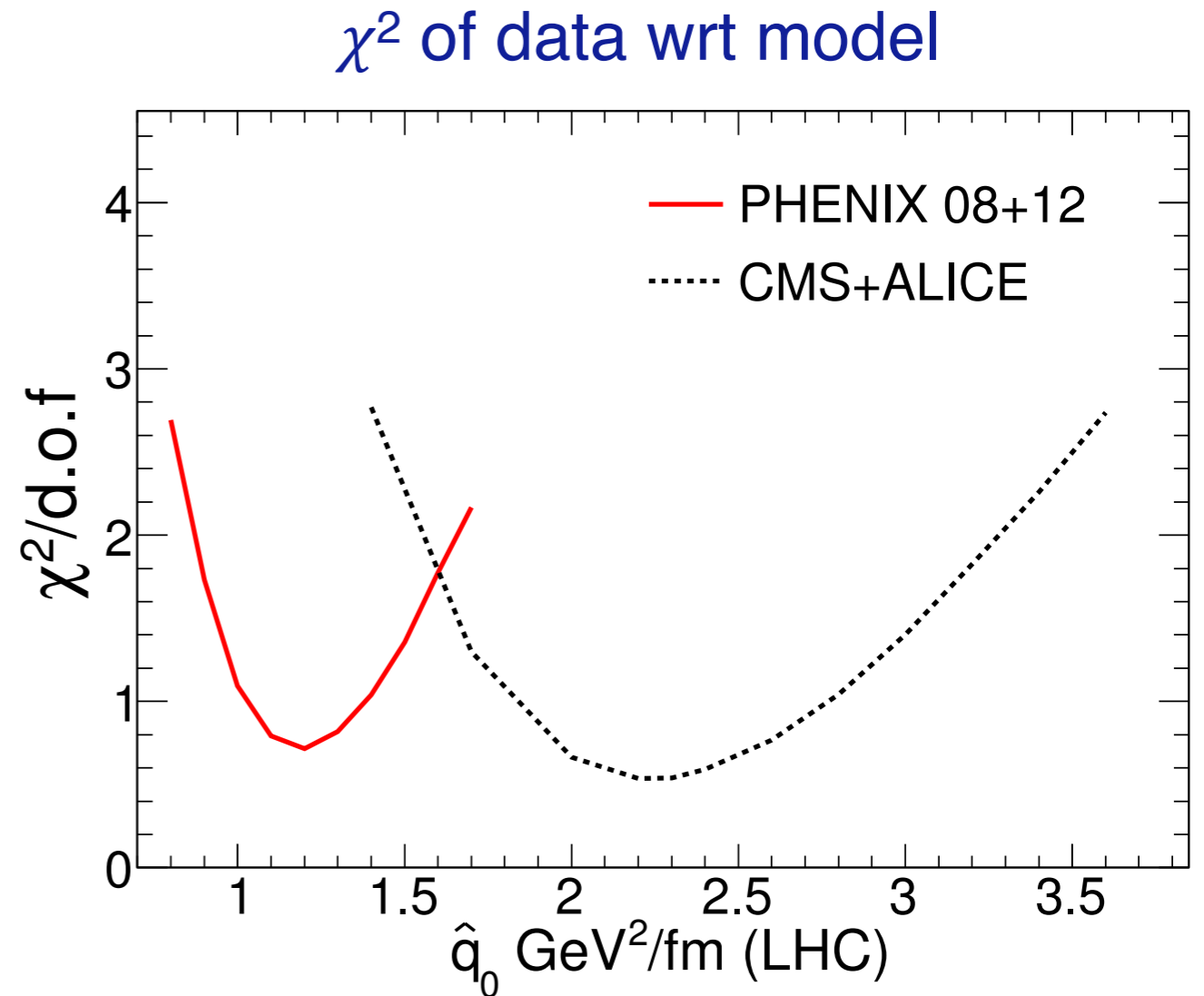
Depends on density ρ through mean free path λ

$$\lambda \propto \frac{1}{\rho}$$

Fitting the model to the data



Burke et al, JET Collaboration, arXiv:1312.5003



Clear minimum: found best value
for transport coefficient

Factor ~ 2 larger at LHC than RHIC

Comparing several models

RHIC:

$$\hat{q} = 1.2 \pm 0.3 \text{ GeV}^2/\text{fm}$$

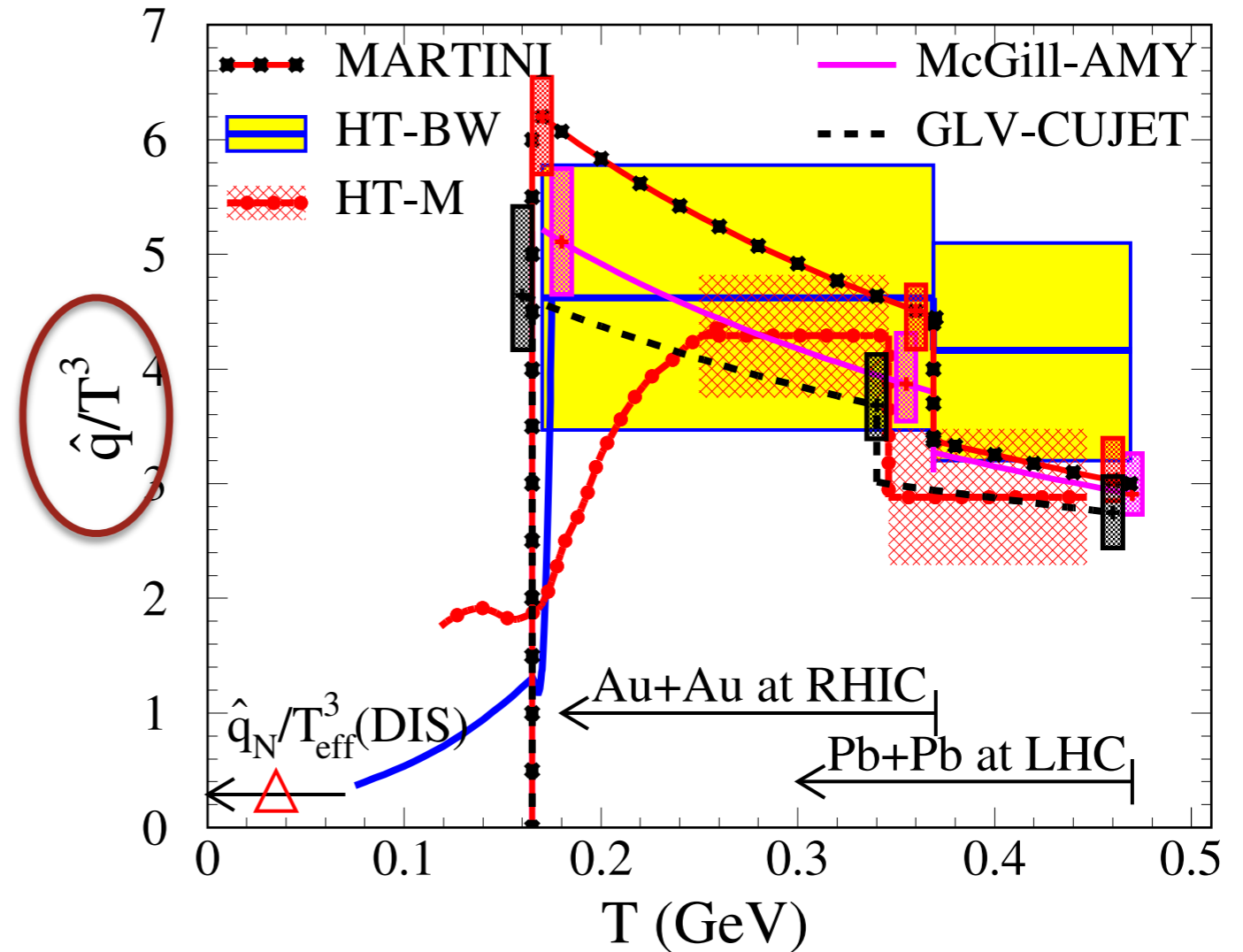
(T=370 MeV)

LHC:

$$\hat{q} = 1.9 \pm 0.7 \text{ GeV}^2/\text{fm}$$

(T=470 MeV)

Expect factor 2.2 from
multiplicity + nuclear size



Burke et al, JET Collaboration, arXiv:1312.5003

\hat{q} values from different models agree

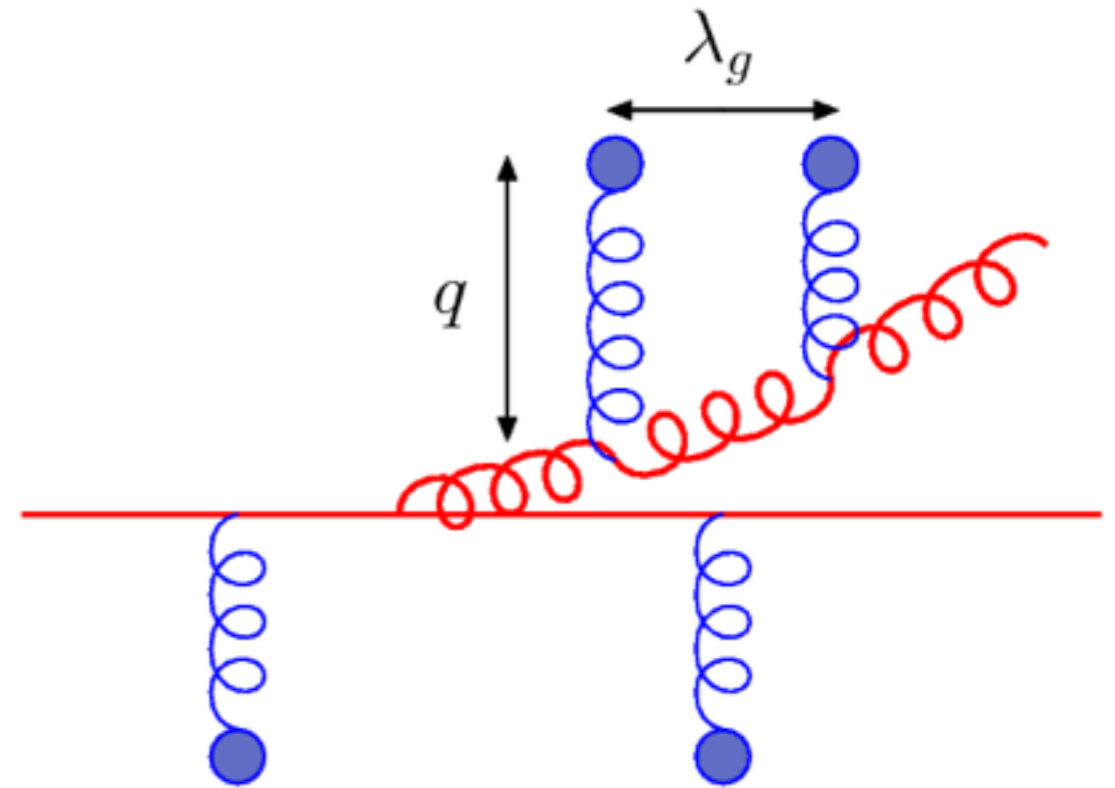
\hat{q}/T^3 larger at RHIC than LHC

Transport coefficient and viscosity

Transport coefficient:
momentum transfer per unit path length

$$\hat{q} = \frac{\langle q_{\perp}^2 \rangle}{\lambda} = \rho \int dq_{\perp}^2 q_{\perp}^2 \frac{d^2 \sigma}{dq_{\perp}^2}$$

$$\rho \propto T^3$$



Viscosity: $\eta \propto \rho \langle p \rangle \lambda$

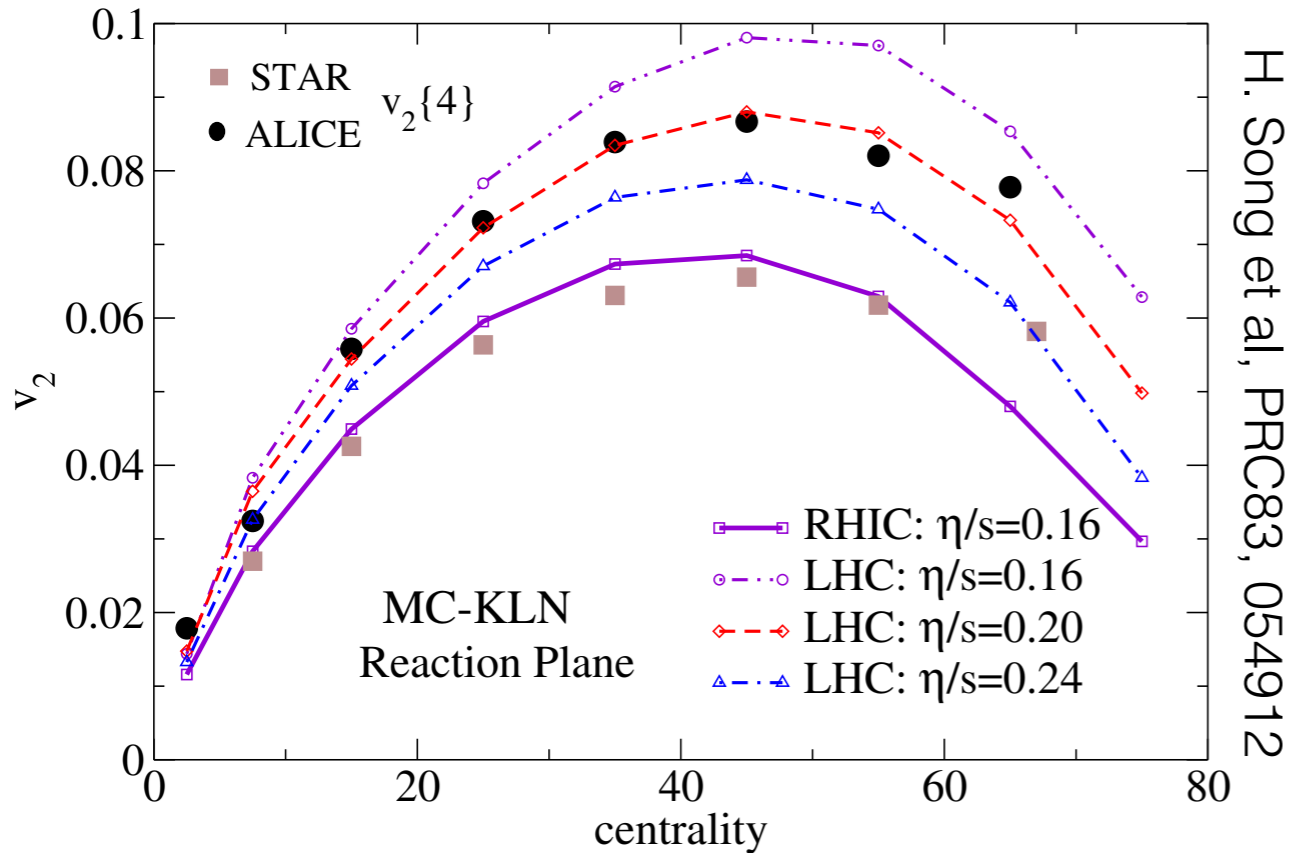
General relation: $\frac{\hat{q}}{T^3} \propto \left(\frac{\eta}{s} \right)^{-1}$

Expect $\frac{\eta}{s} \approx 1.25 \frac{T^3}{\hat{q}}$ for a QCD medium

Majumder, Muller and Wang, PRL99, 192301

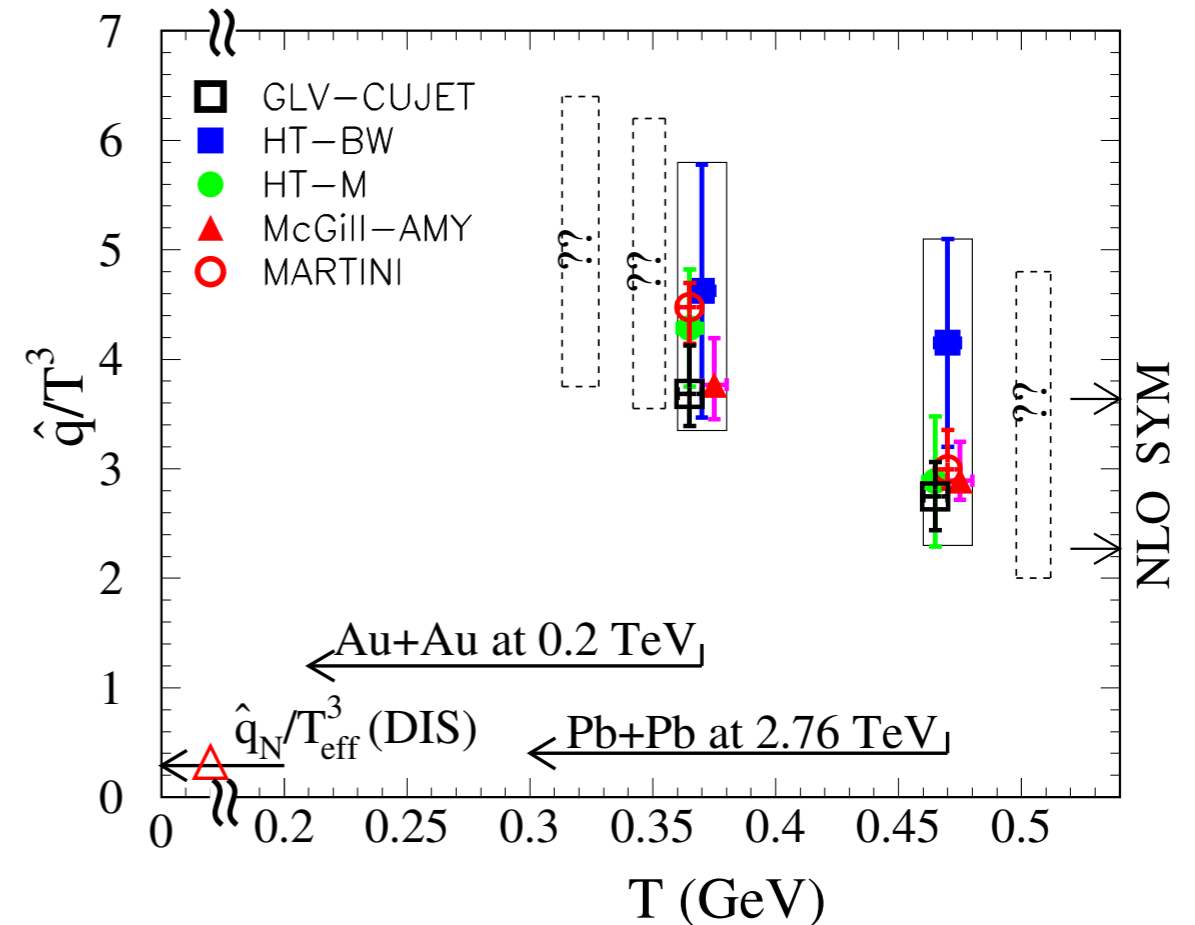
Relation transport coefficient and viscosity

Elliptic flow



(Scaled) viscosity slightly larger at LHC

Transport coefficient from R_{AA}



Scaled transport coefficient slightly smaller at LHC

Increase of η/s and decrease of q/T^3 with collision energy are probably due to a common origin, e.g. running α_s

Results agree reasonably well with expectation:
$$\frac{\eta}{s} \approx 1.25 \frac{T^3}{\hat{q}}$$

Summary

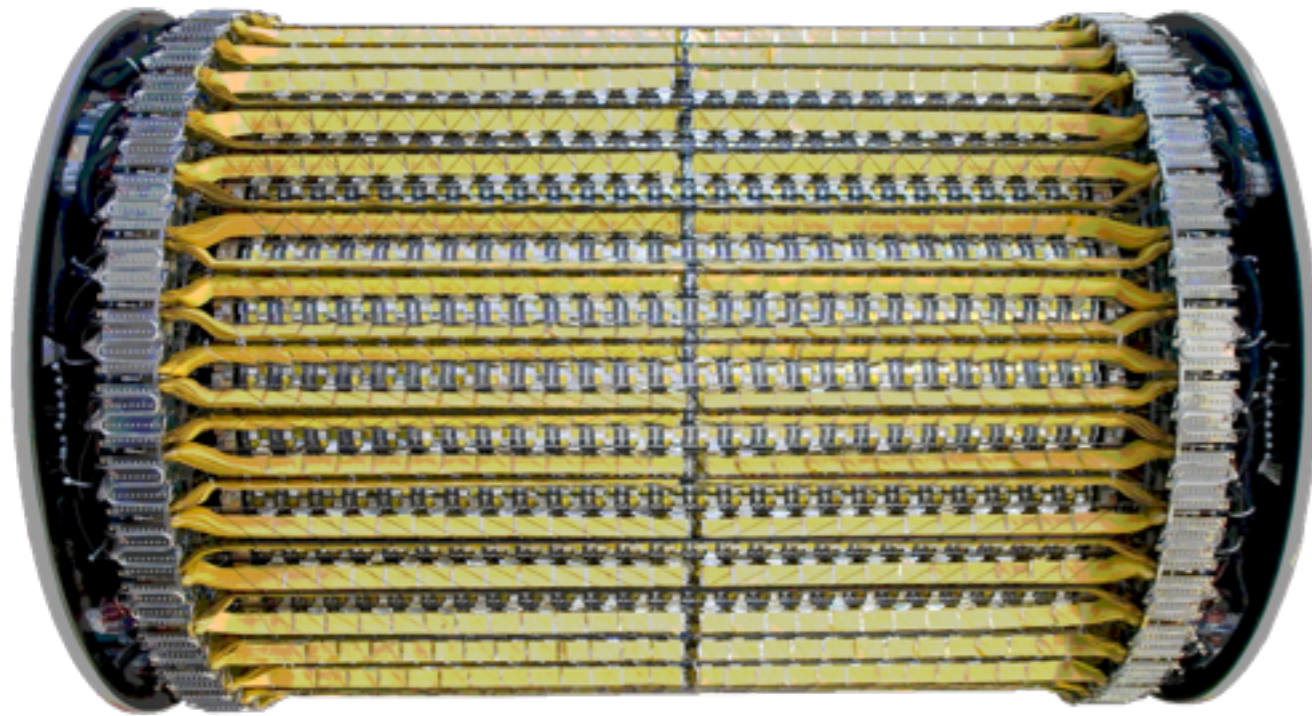
- Heavy ion collisions study hot and dense nuclear matter
- Two main experimental approaches:
 - Study QGP fragments, e.g. elliptic/triangular flow
Indicates **very low value of η/s** $\frac{\eta}{s} \approx 0.1-0.3$
 - Probe QGP with self-generated probes, e.g. high- p_T particle production
Large density, energy loss $\hat{q} \approx 1-2 \text{ GeV}^2/\text{fm}$
- Both observations consistent with very dense system, small mean free path

Run 2 of the LHC: larger data samples to explore flow, parton energy loss mechanisms

Extra slides

The Inner Tracking System

Dutch contribution to ALICE



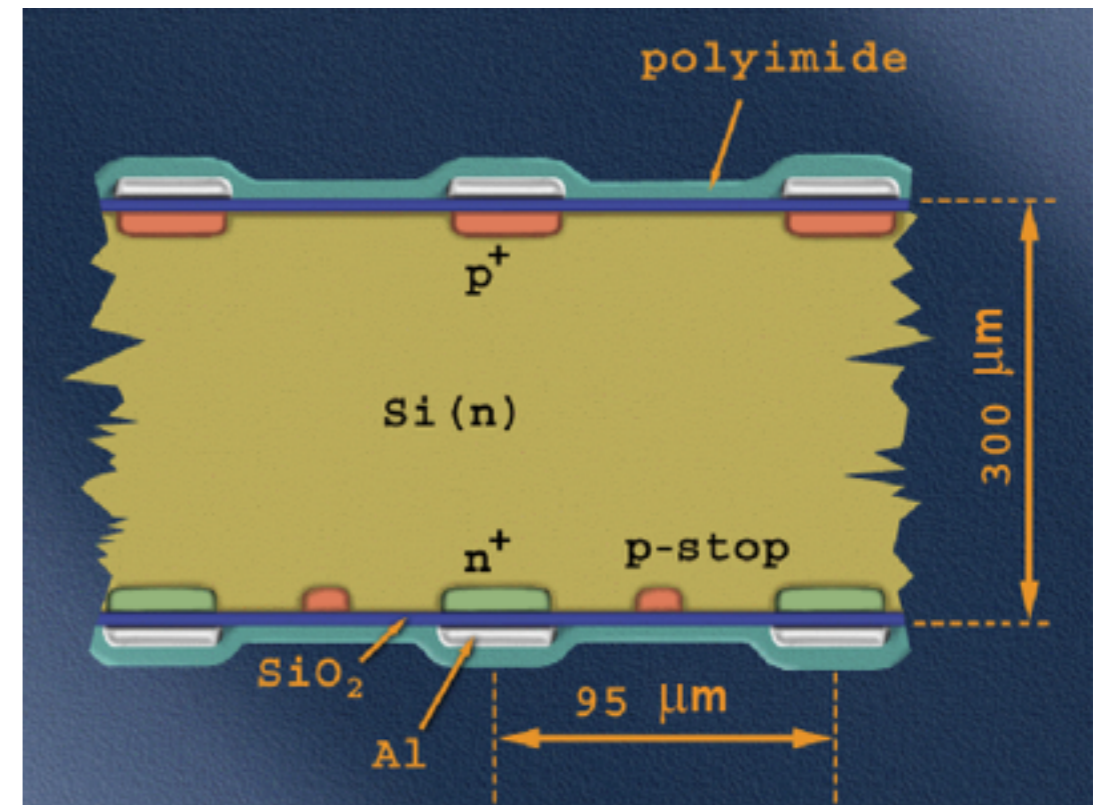
1698 double sided strip sensors

73 * 40 mm² 300 um thick

768 strips on each side

- + 2 layers of Silicon Drift detectors
- + 2 layers of Silicon Pixel detectors

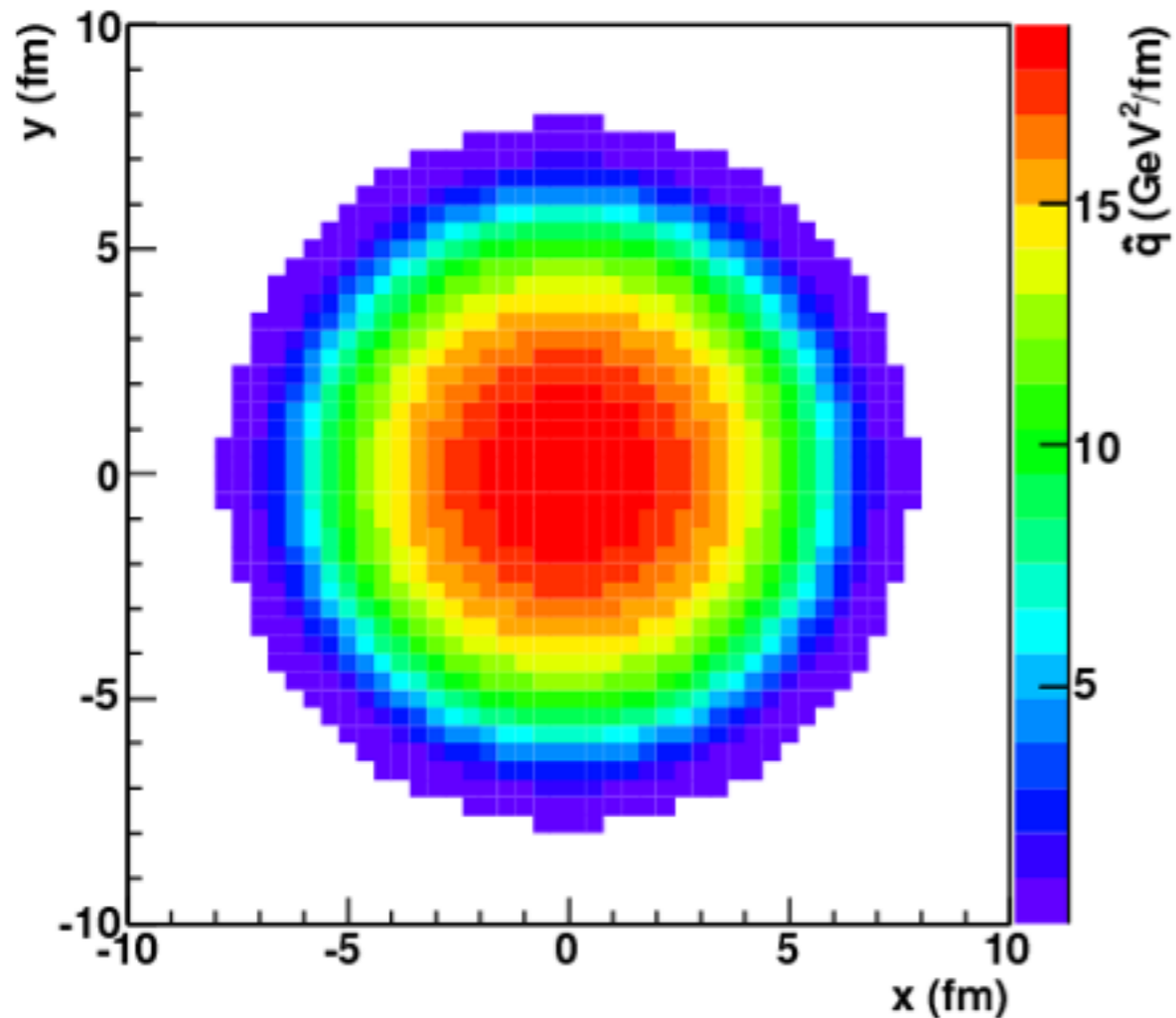
Strip detector, cross section



- Charged particle generates free electrons+holes
- Drift (E-field) to p, n-doped strips
- Detect image charge in Al strips

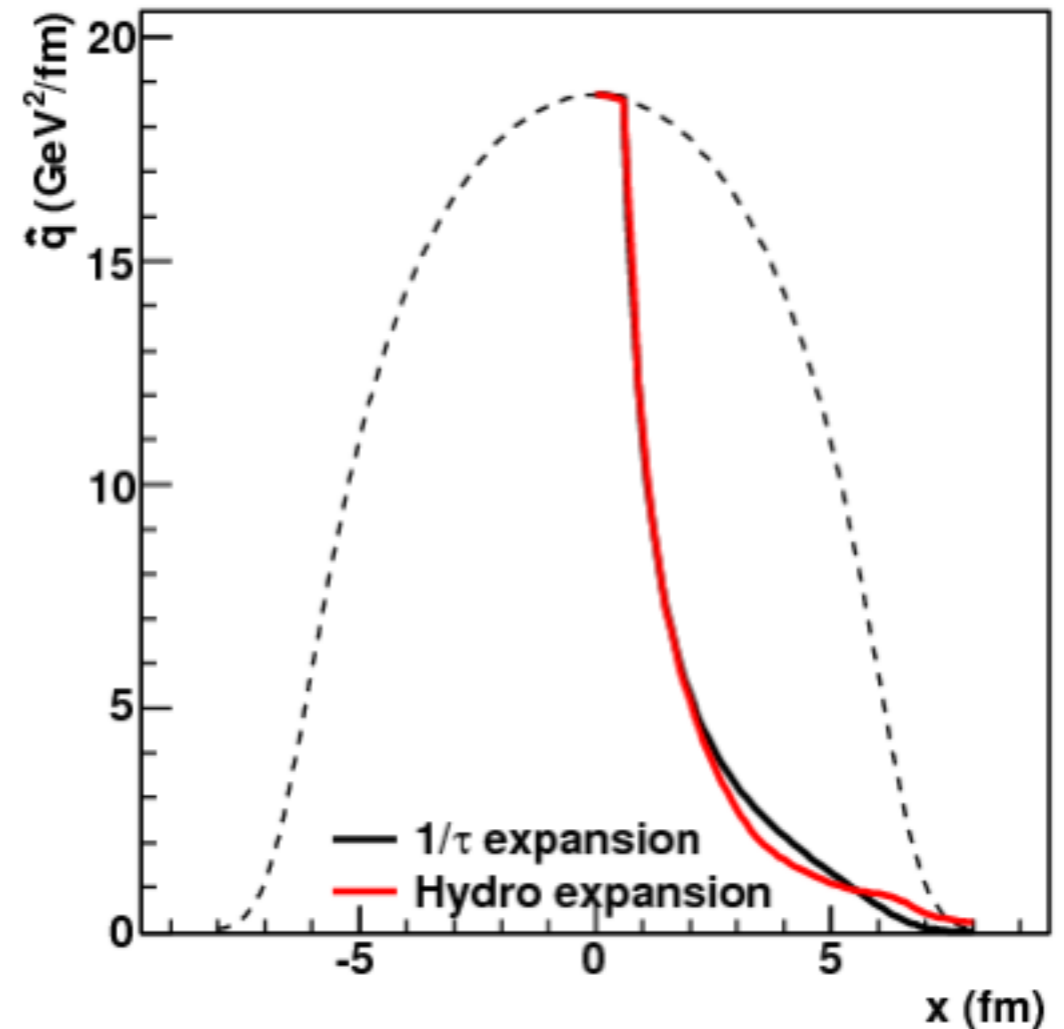
Geometry

Density profile



Profile at $\tau \sim \tau_{\text{form}}$ known

Density along parton path



Longitudinal expansion
dilutes medium
 \Rightarrow Important effect

Space-time evolution is taken into account in modeling

A simplified approach

$$\left. \frac{dN}{dp_T} \right|_{hadr} = \left[\left. \frac{dN}{dE} \right|_{jets} \right] \otimes P(\Delta E) \otimes \left[D(p_{T,hadr} / E_{jet}) \right]$$

Parton spectrum
Energy loss distribution
Fragmentation (function)

$\left. \frac{dN}{dE} \right|_{jets}$
 known
 pQCDxPDF

$P(\Delta E)$
 extract

$D(p_{T,hadr} / E_{jet})$
 'known' from e^+e^-

This is where the information about the medium is

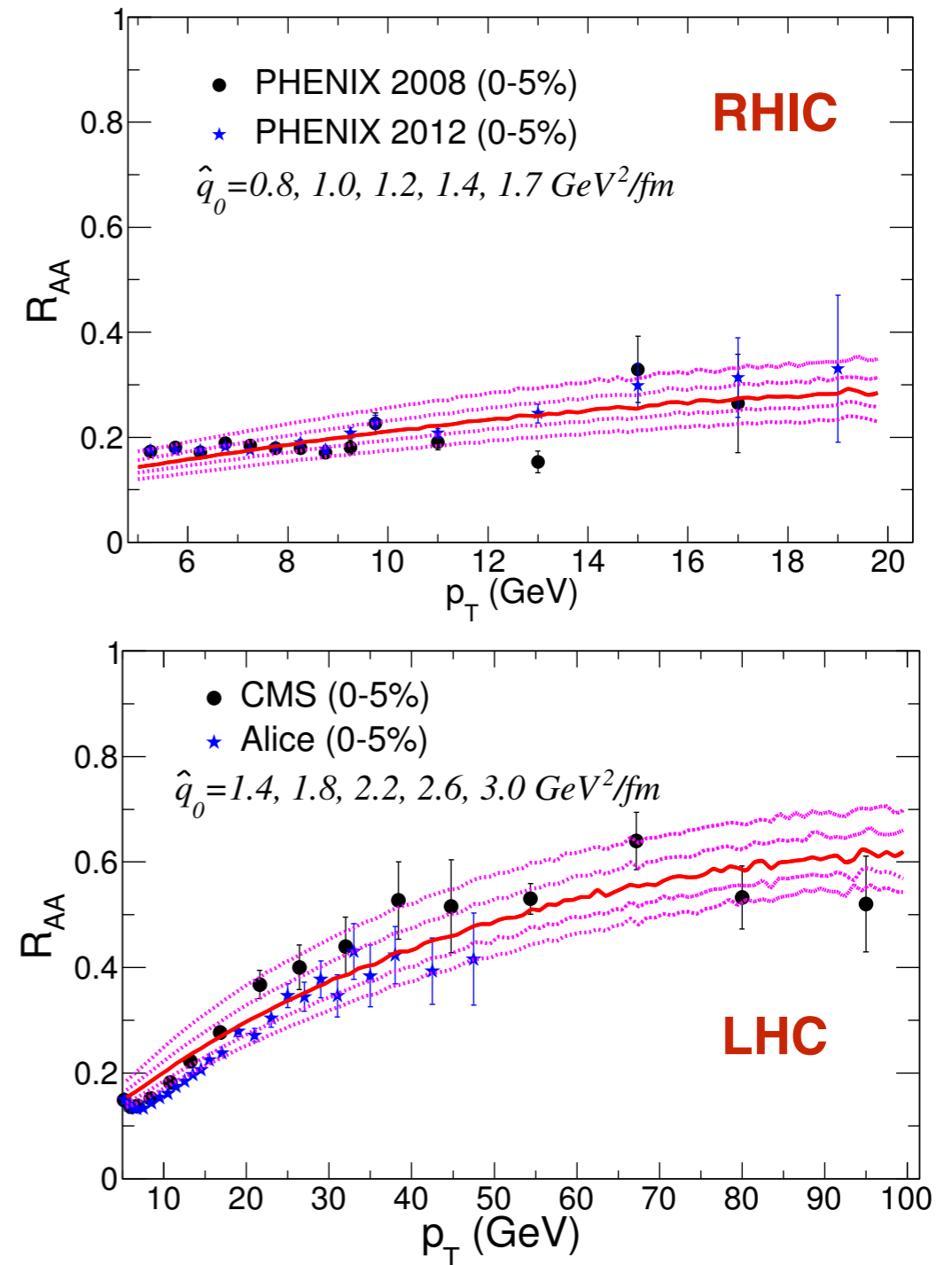
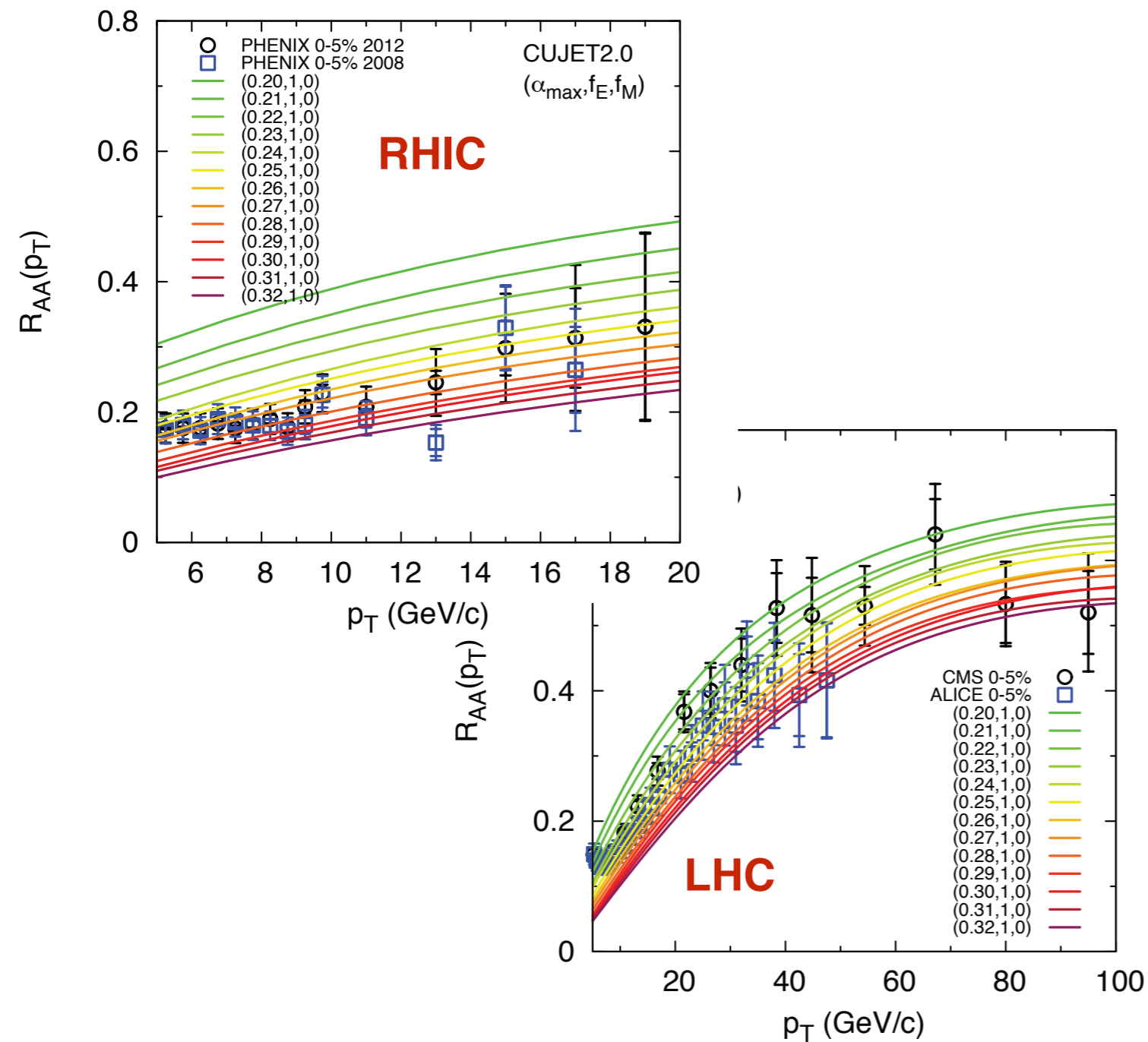
$P(\Delta E)$ combines geometry
 with the intrinsic process

– Unavoidable for many observables

Notes:

- This is the simplest ansatz – most calculation to date use it (except some MCs)
- Jet, γ -jet measurements 'fix' E , removing one of the convolutions

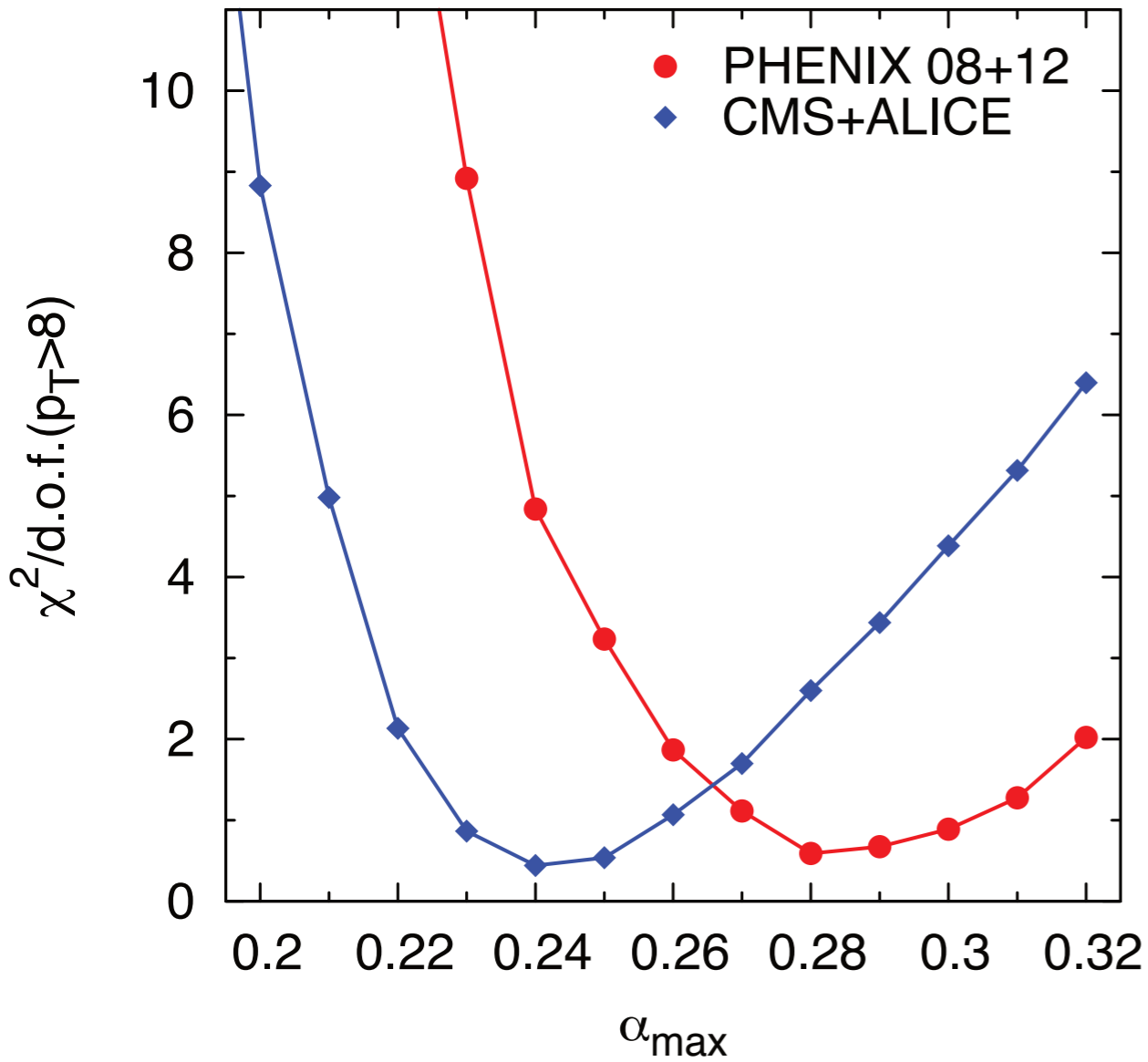
RHIC and LHC



Systematic comparison of energy loss models with data
 Medium modeled by Hydro (2+1D, 3+1D)
 p_T dependence matches reasonably well

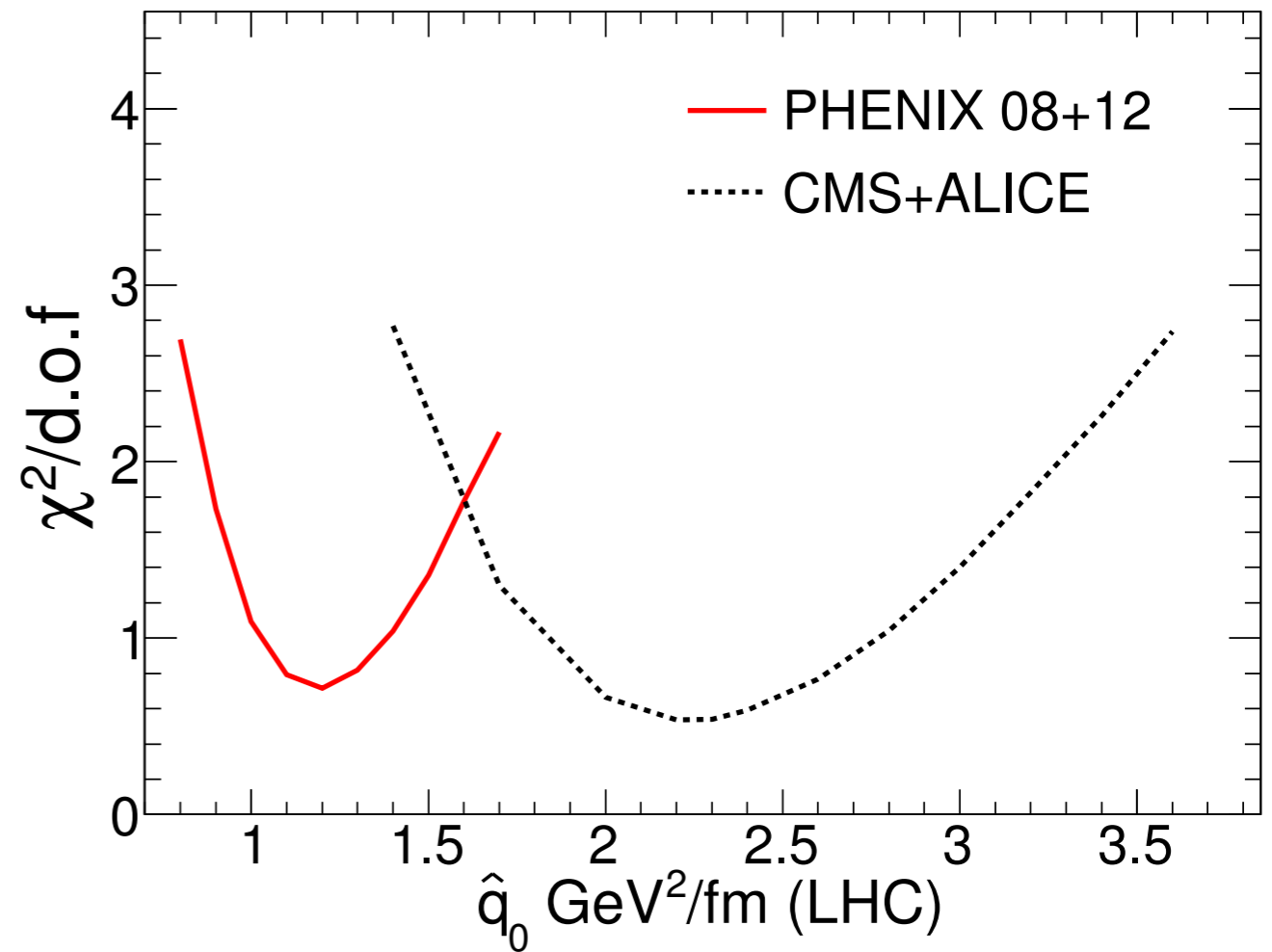
RHIC and LHC

CUJET 2.0



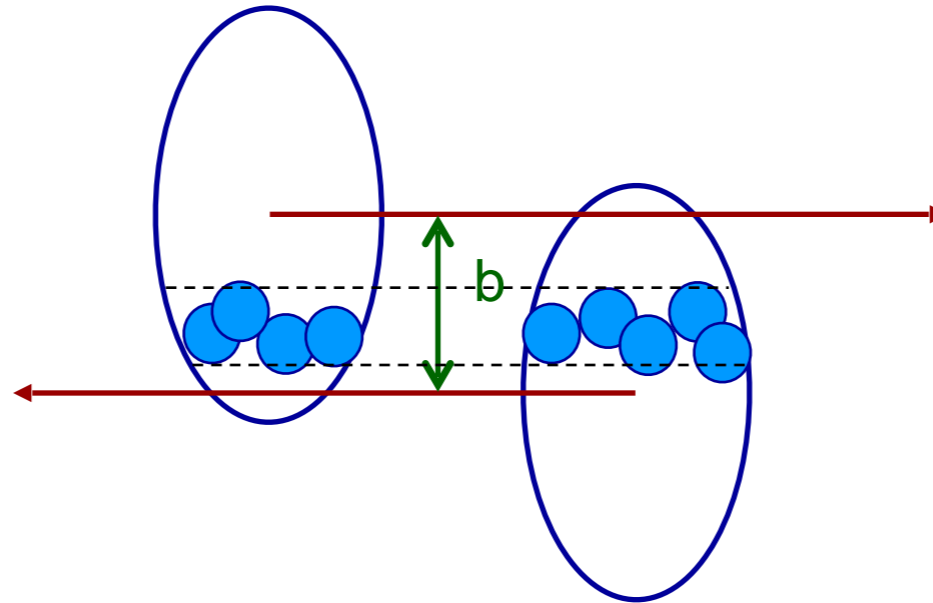
CUJET: α_s is medium parameter
Lower at LHC

HT-BW



HT: transport coeff is parameter
Higher at LHC

Nuclear geometry: N_{part} , N_{coll}



Two limiting possibilities:

- Each nucleon only **interacts once**, 'wounded nucleons'

$$N_{\text{part}} = n_A + n_B \quad (\text{ex: } 4 + 5 = 9 + \dots)$$

Relevant for **soft production**; long timescales: $\sigma \propto N_{\text{part}}$

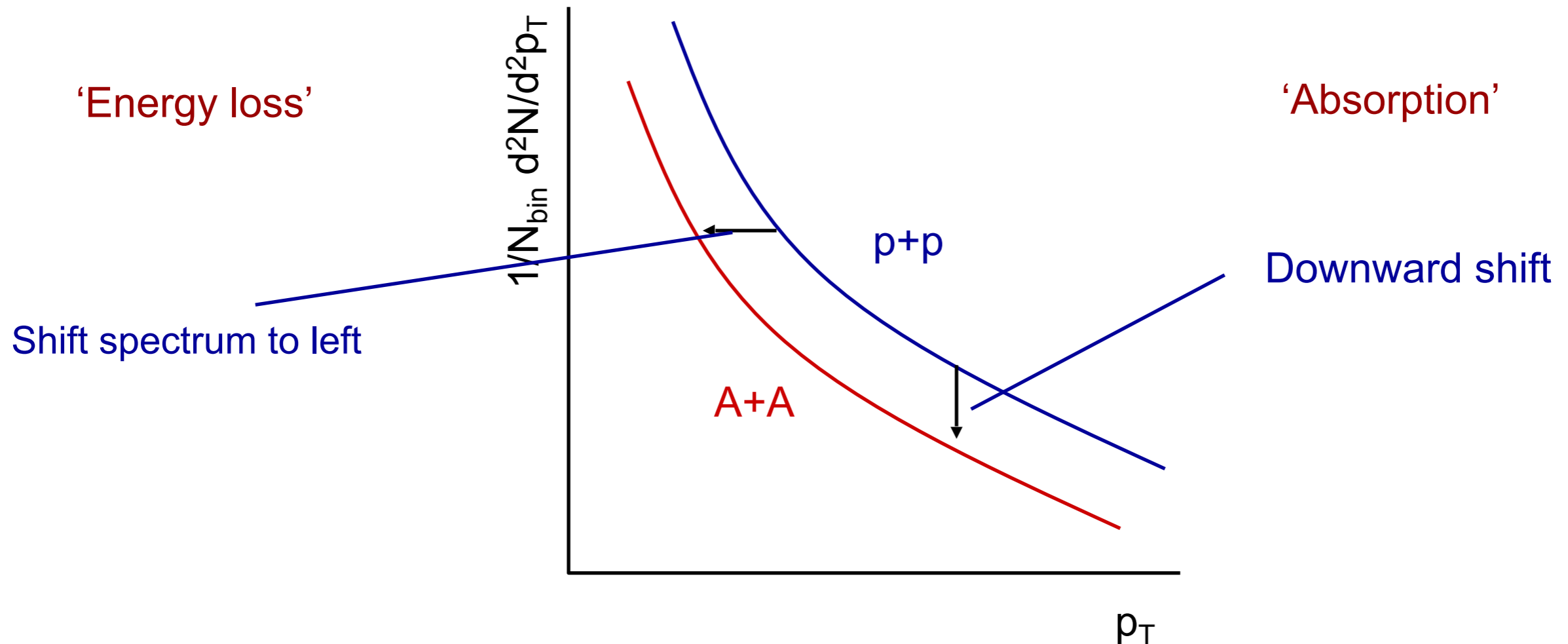
- Nucleons **interact with all** nucleons they encounter

$$N_{\text{coll}} = n_A \times n_B \quad (\text{ex: } 4 \times 5 = 20 + \dots)$$

Relevant for **hard processes**; short timescales: $\sigma \propto N_{\text{bin}}$

Nuclear modification factor R_{AA}

$$R_{AA} = \frac{dN/dp_T|_{A+A}}{N_{coll} dN/dp_T|_{p+p}}$$

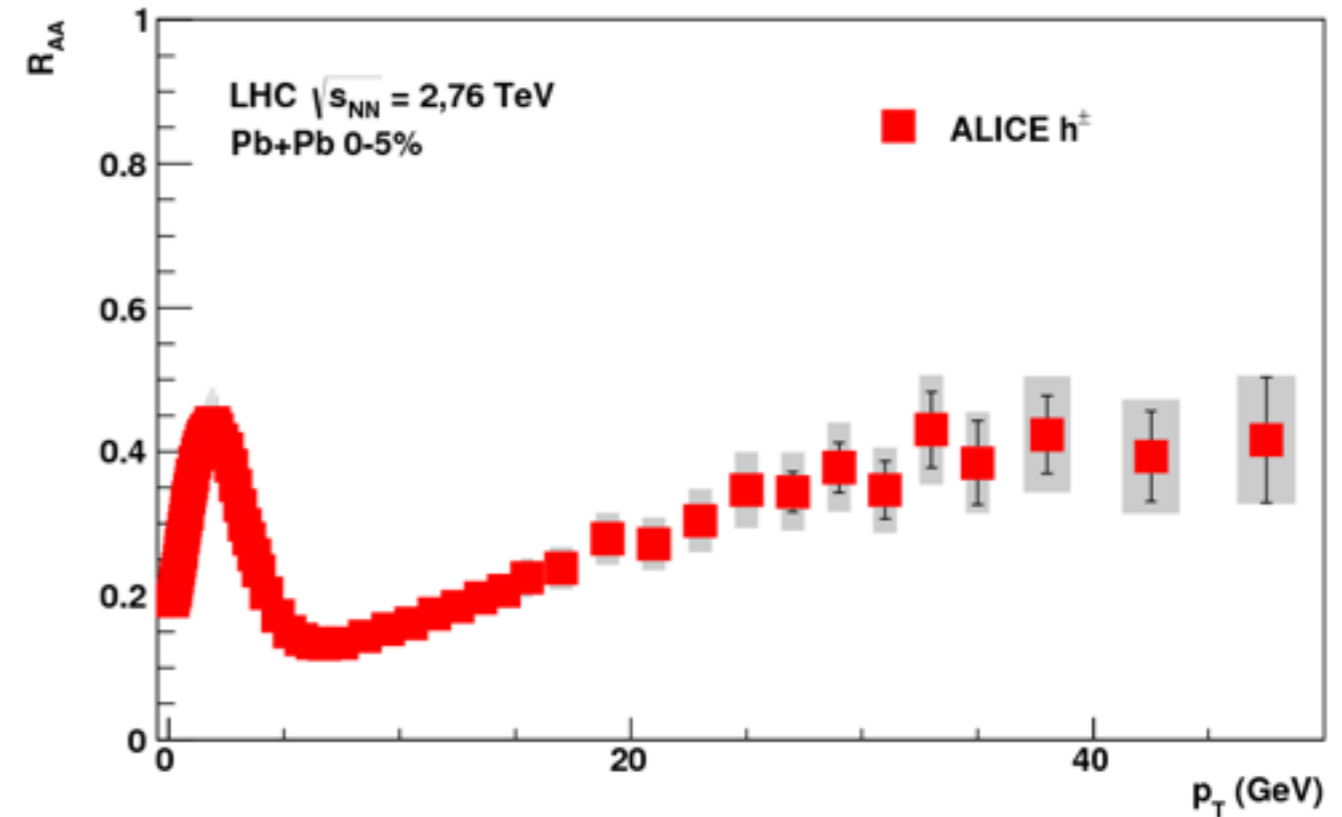
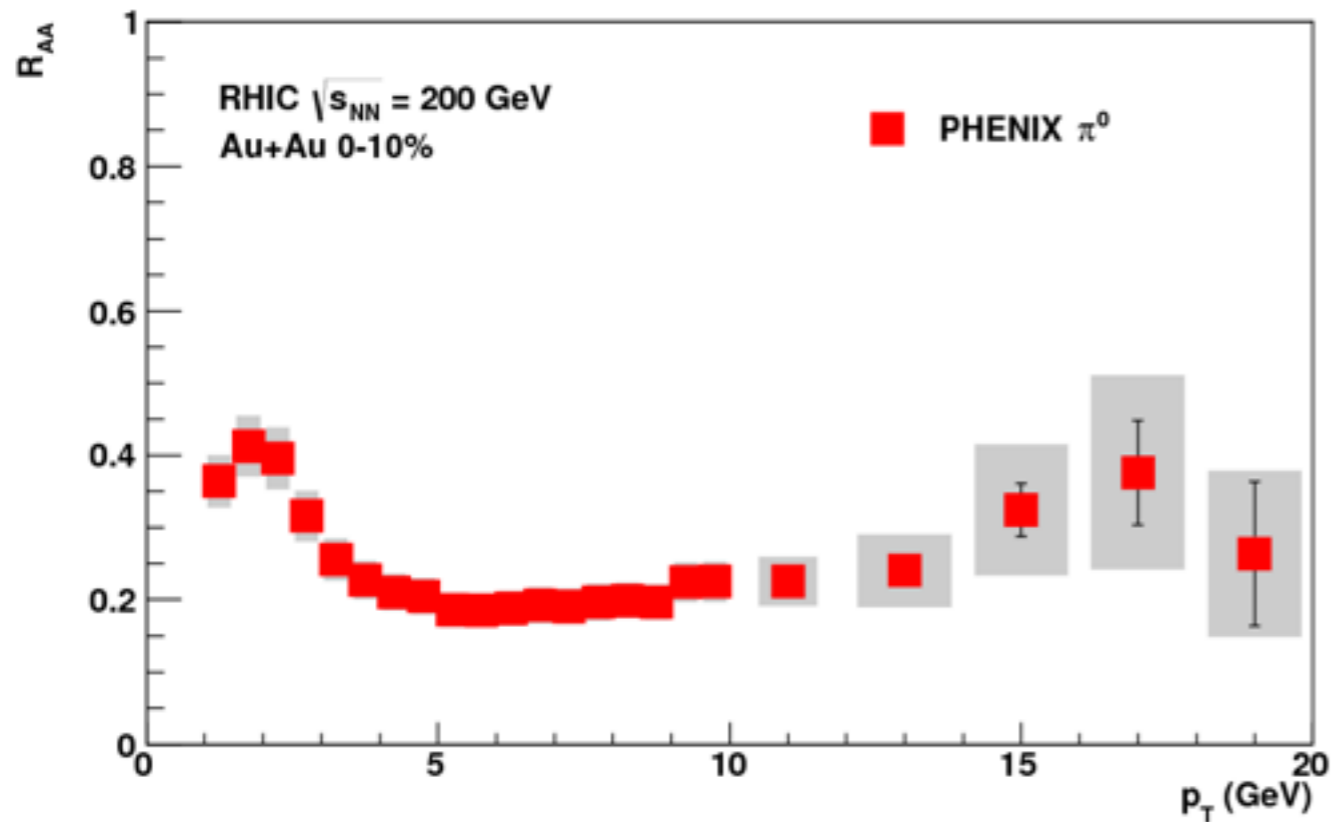


Measured R_{AA} is a ratio of yields at a given p_T
The physical mechanism is energy loss; shift of yield to lower p_T

The full range of physical pictures can be captured with an energy loss distribution $P(\Delta E)$

Nuclear modification factor

$$R_{AA} = \frac{dN / dp_T|_{Pb+Pb}}{N_{coll} dN / dp_T|_{p+p}}$$

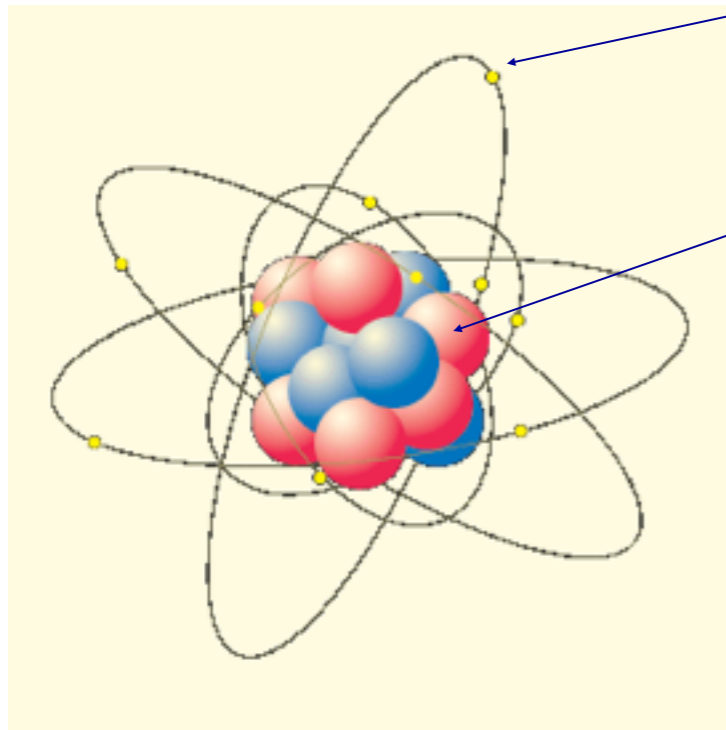


Suppression factor 2-6
Significant p_T -dependence
Similar at RHIC and LHC?

So what does it mean?

Quarks and the strong interaction

Atom



Electron
elementary, point-particle

Protons, neutrons
Composite particle
⇒ quarks

Standard Model: elementary particles

Quarks:

Electrical charge

Strong charge (color)

up

down

charm

strange

top

bottom

Leptons:

Electrical charge

electron

ν_e

Muon

ν_μ

Tau

ν_τ

Force carriers:

photon

gluon

W,Z-boson

EM force

strong force

weak force

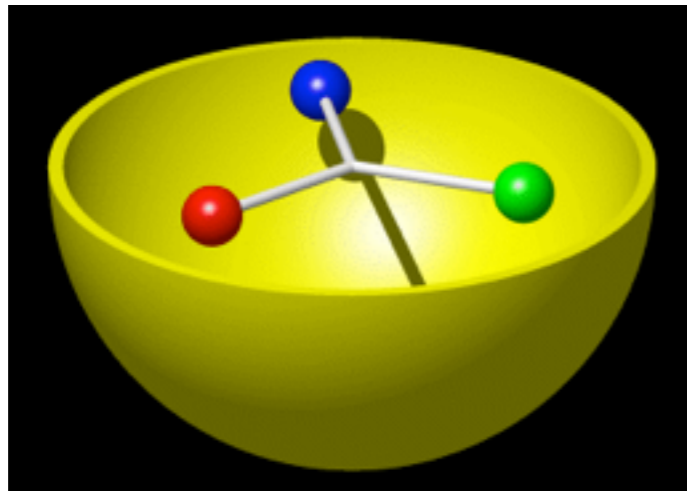
+anti-particles

EM force binds electrons
to nucleus in atom

Strong force binds nucleons
in nucleus and quarks in nucleons

QCD and quark parton model

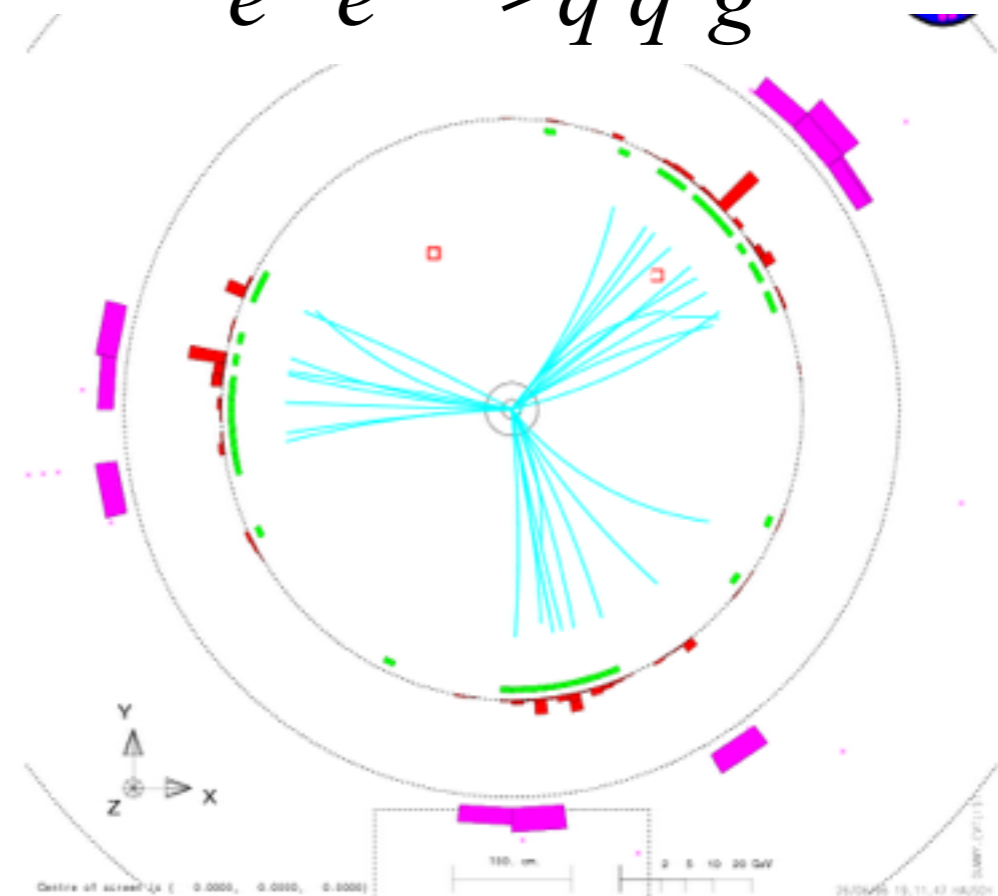
At low energies, quarks are confined in hadrons



protons, neutrons,
pions, kaons
+ many others

At high energies, quarks and gluons are manifest

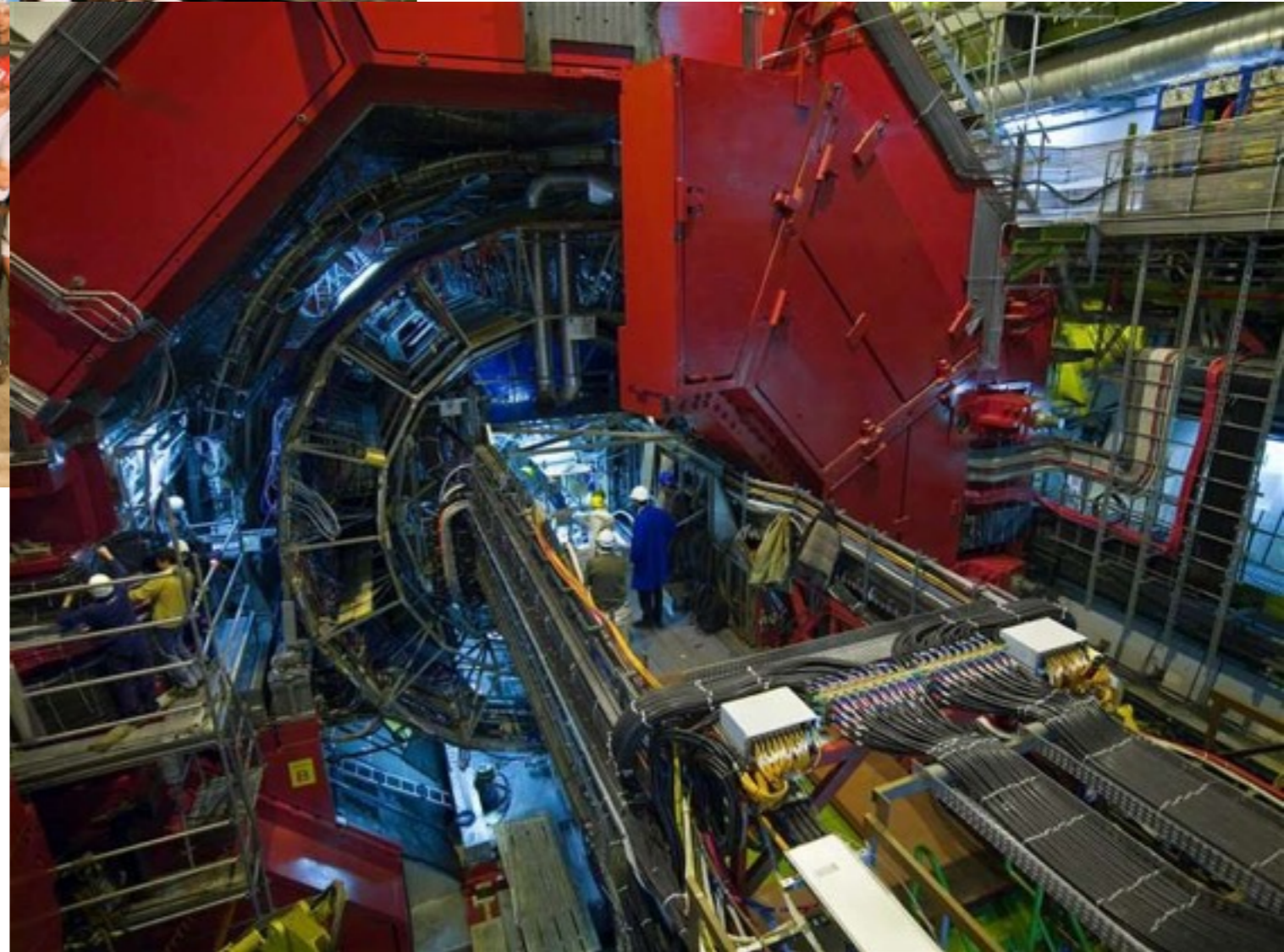
$$e^+ e^- \rightarrow q \bar{q} g$$



Experimental signature: jets of hadrons

Goal of Heavy Ion Physics:
Study dynamics of QCD and confinement
in many-body systems

ALICE in real life



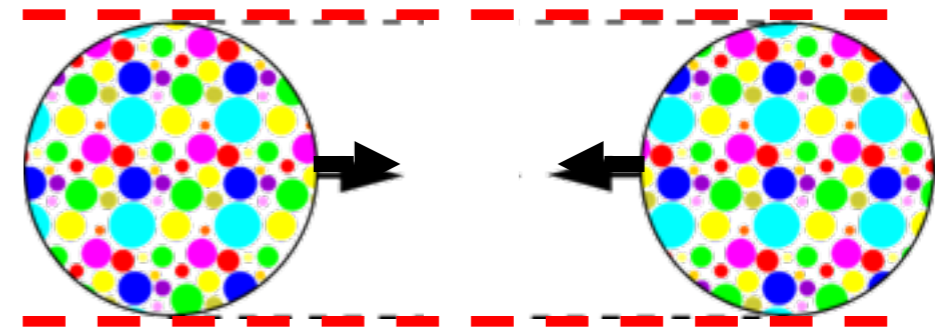
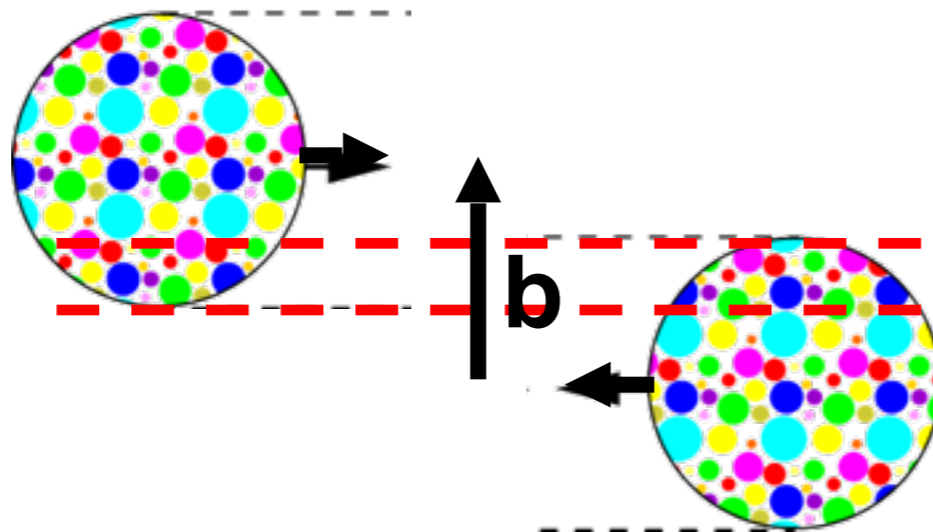
Collision centrality

Nuclei are large compared to the range of strong force

Peripheral collision

Central collision

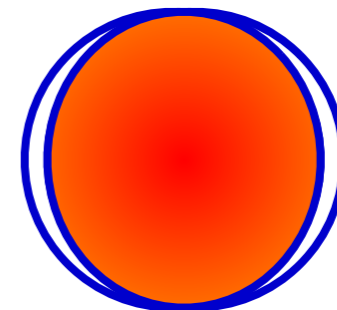
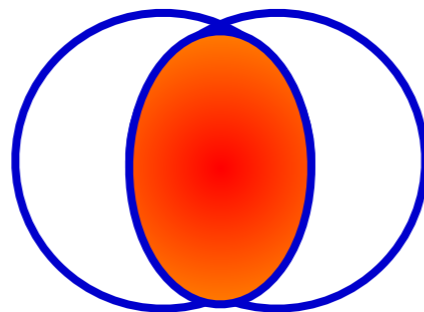
top/side view:



b finite

$b \sim 0$ fm

front view:

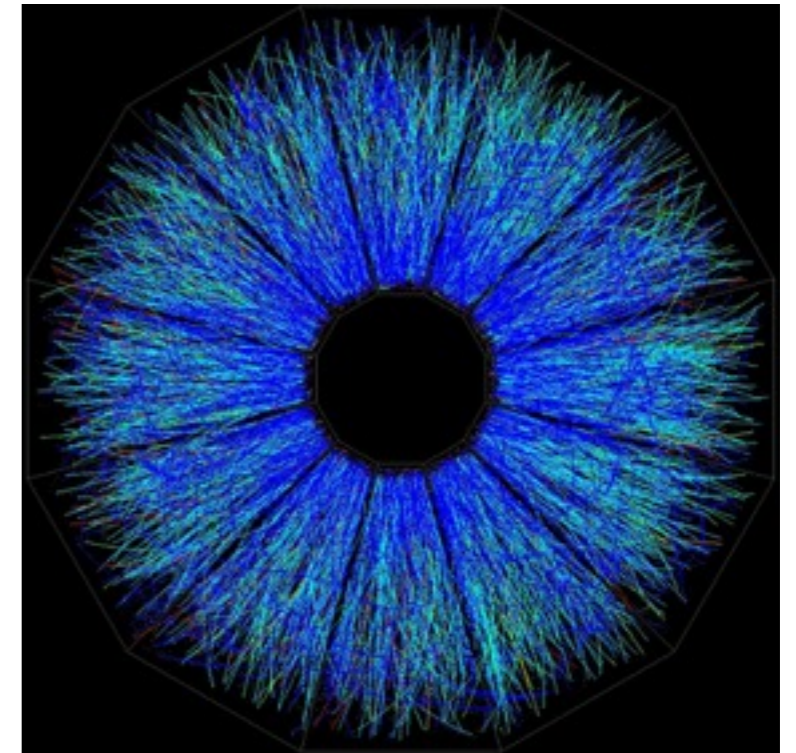
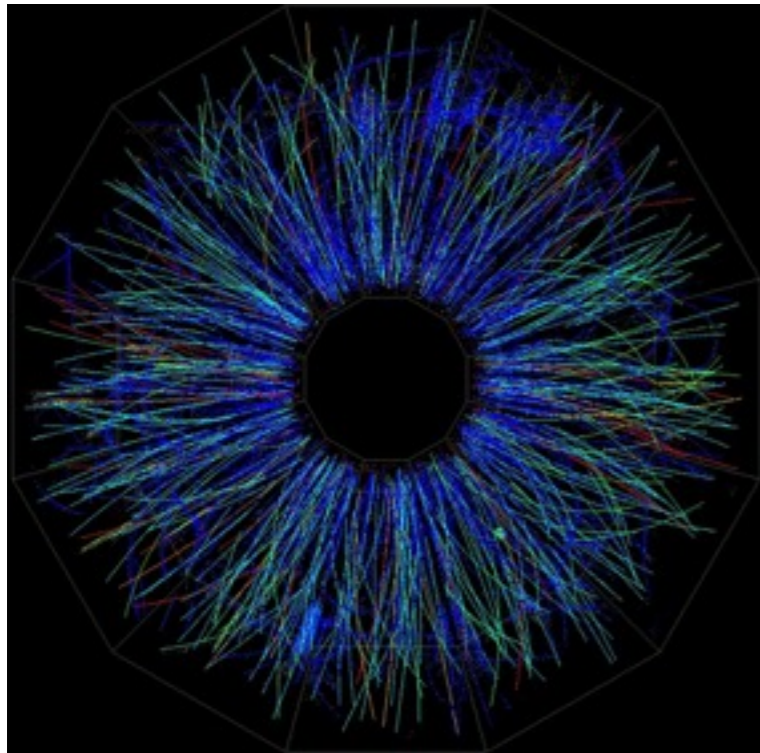


This talk: concentrate on central collisions

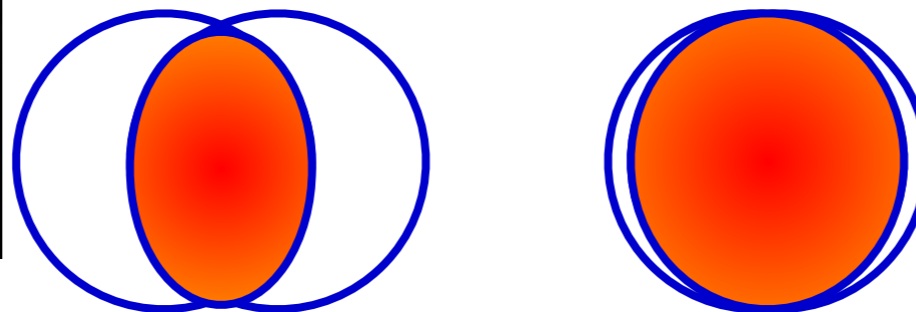
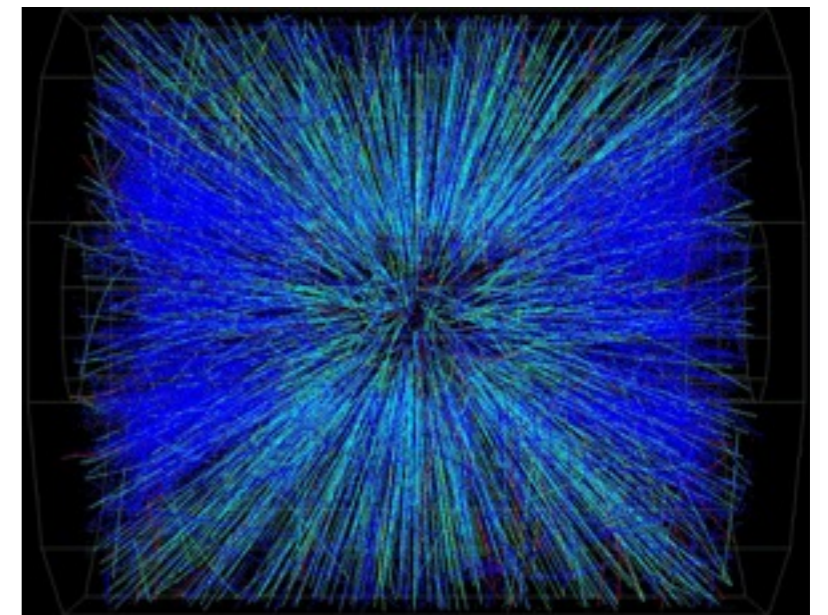
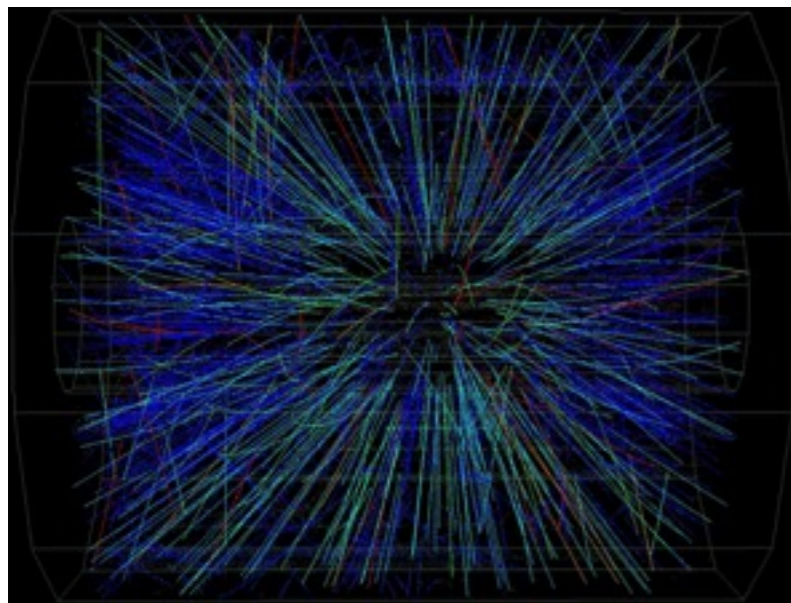
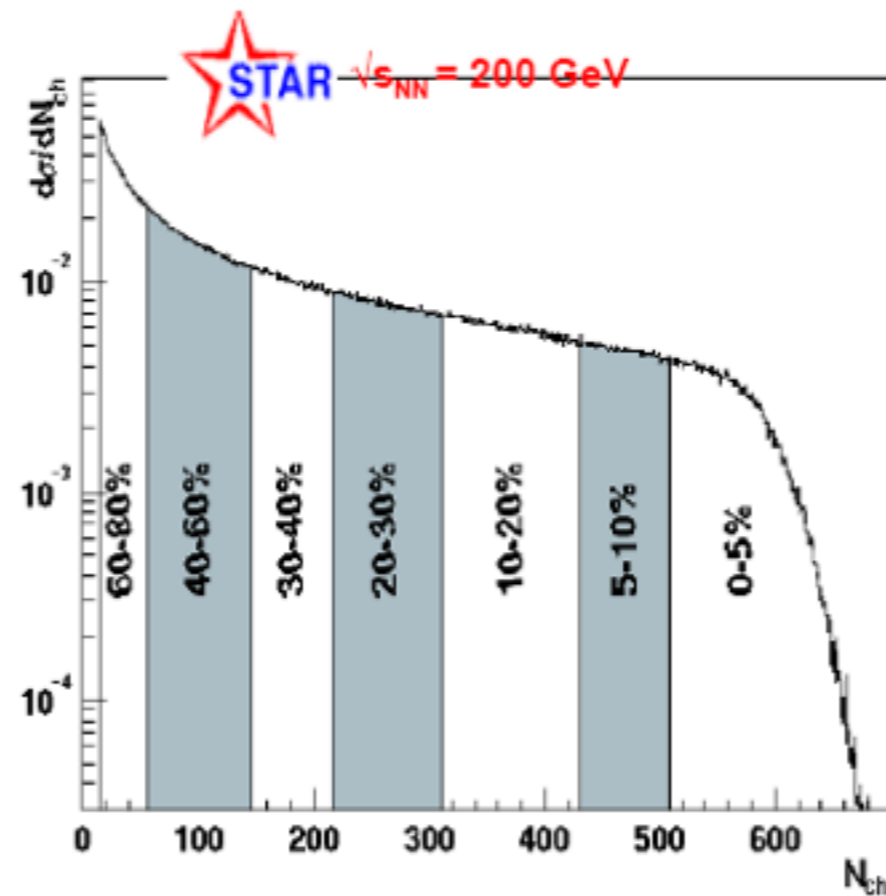
Centrality continued

peripheral

central



Multiplicity distribution



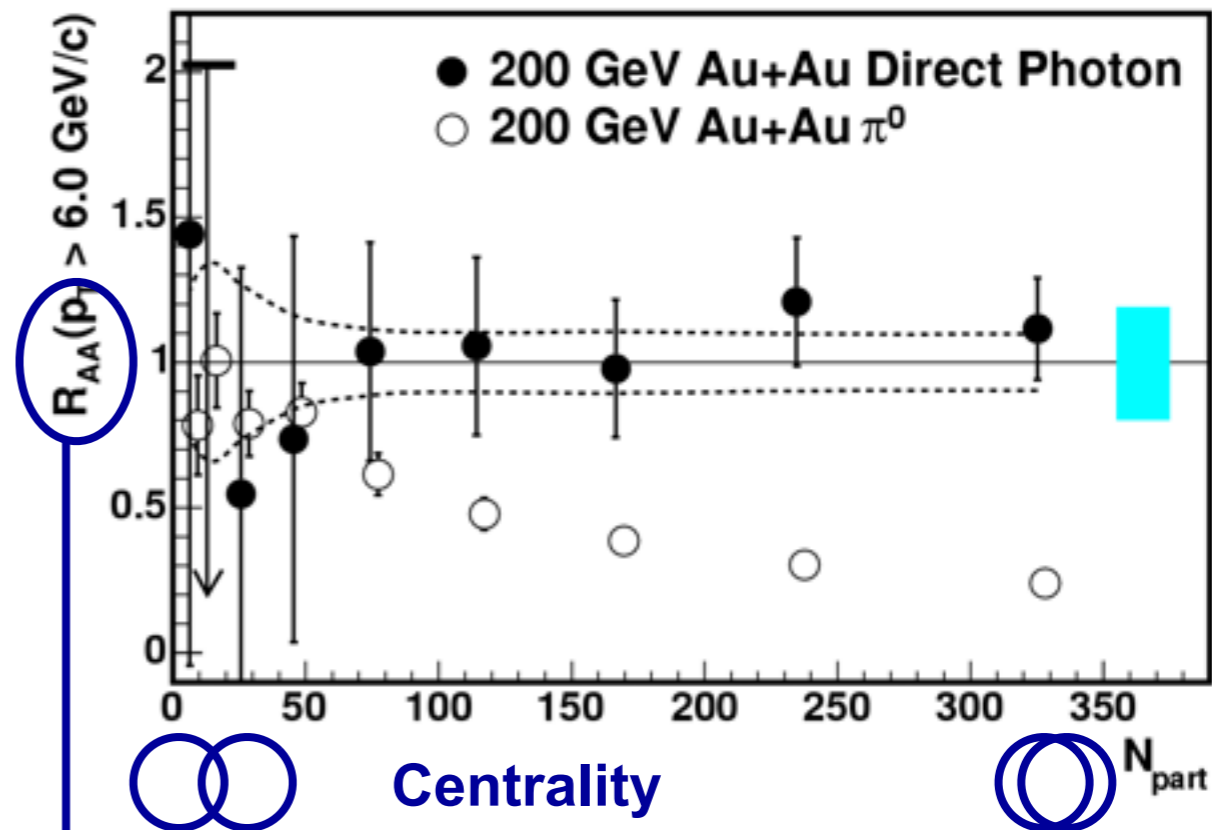
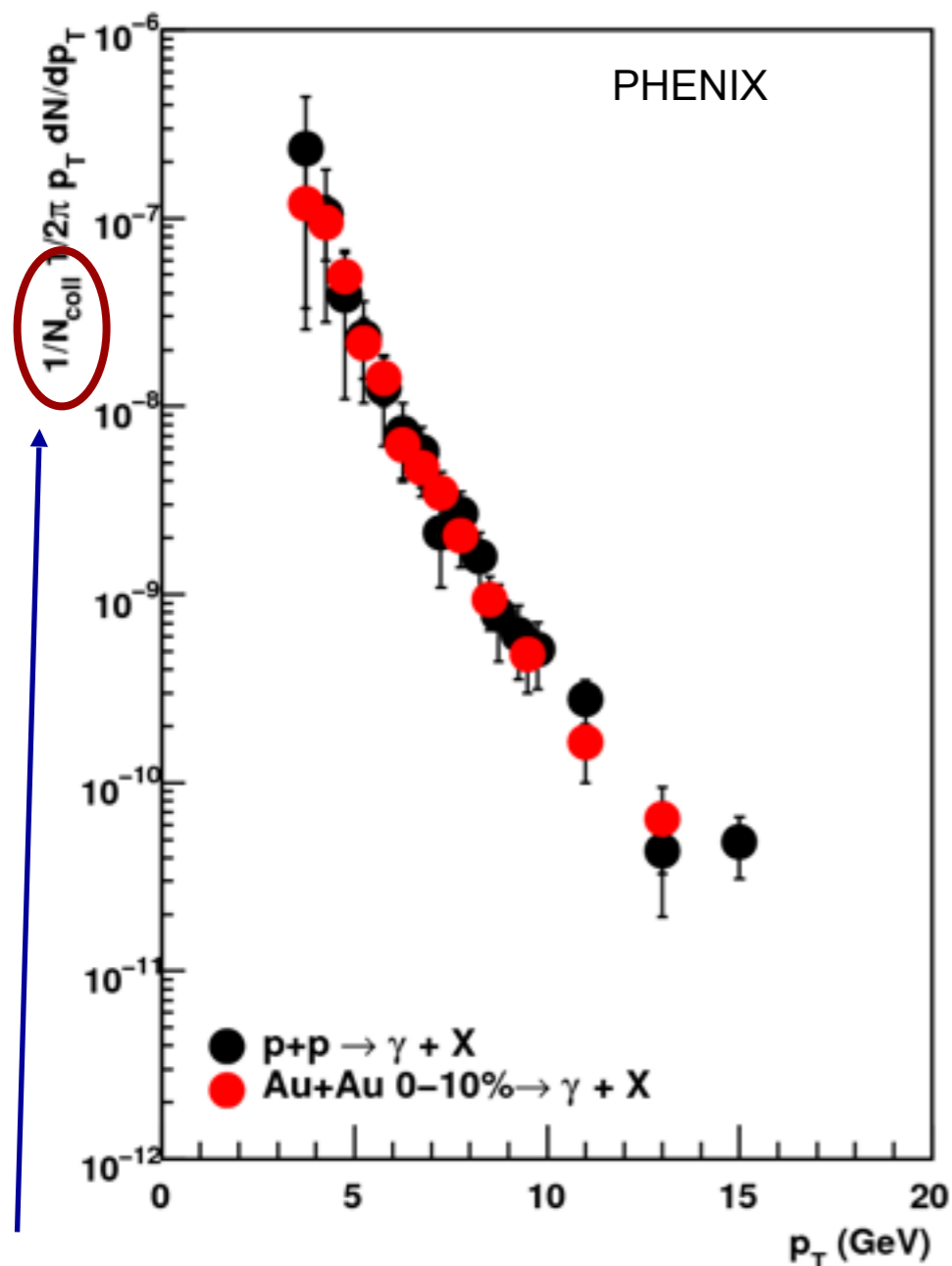
Experimental measure of centrality: multiplicity

Need to take into account volume of collision zone for production rates

Testing volume (N_{coll}) scaling in Au+Au

Direct γ spectra

PHENIX, PRL 94, 232301



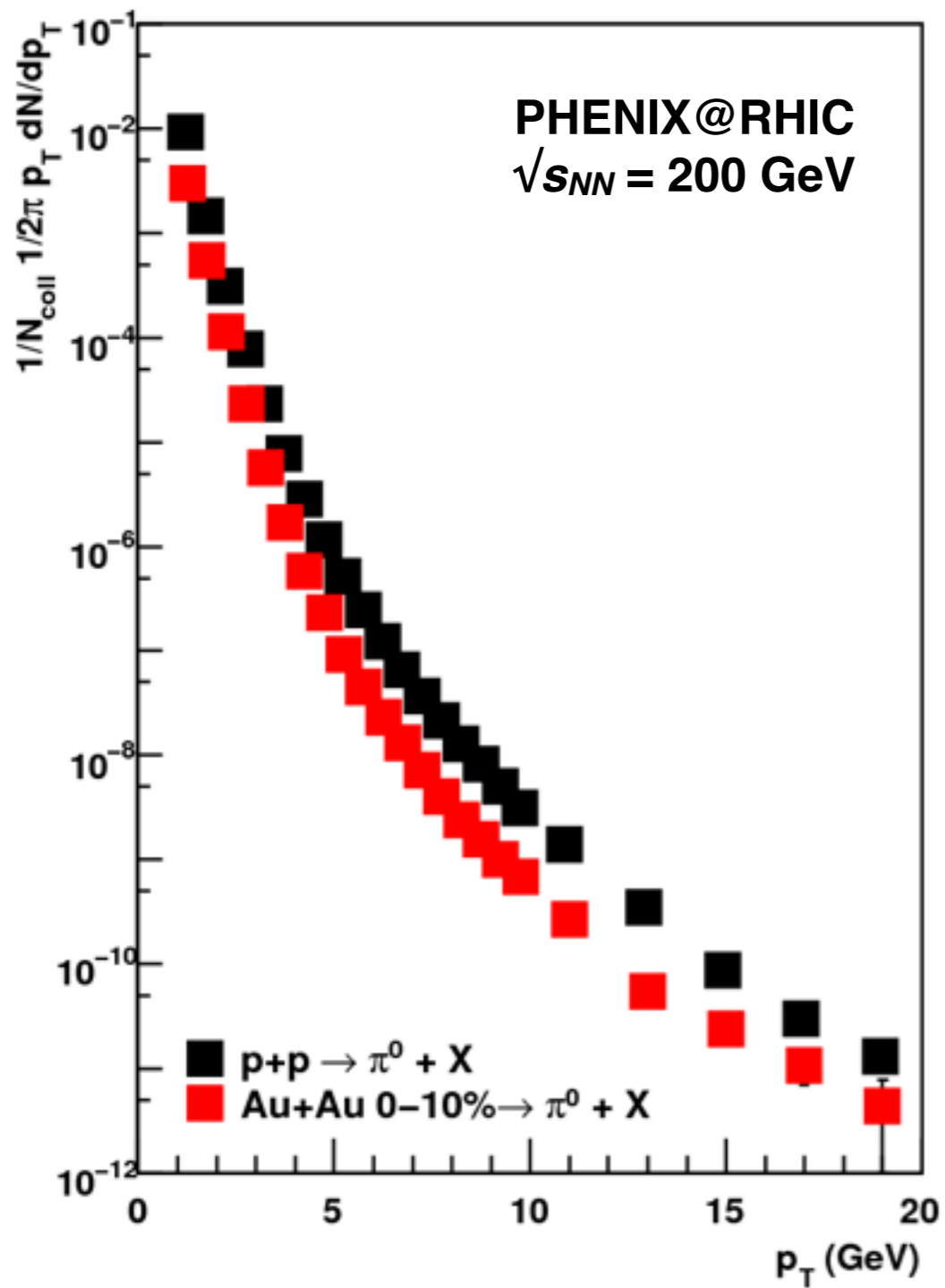
$$R_{AA} = \frac{dN/dp_T|_{A+A}}{N_{coll} dN/dp_T|_{p+p}}$$

Scaled by N_{coll}

Direct γ in A+A scales with N_{coll}

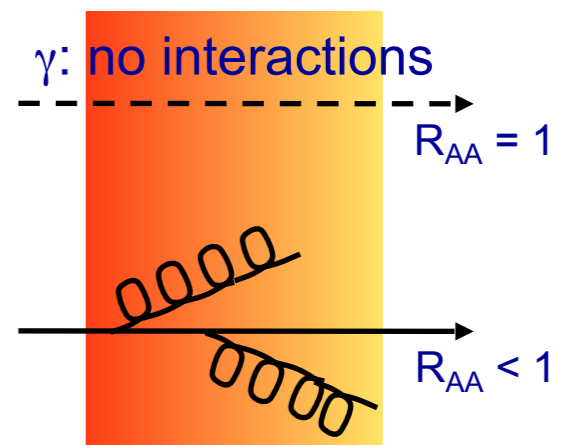
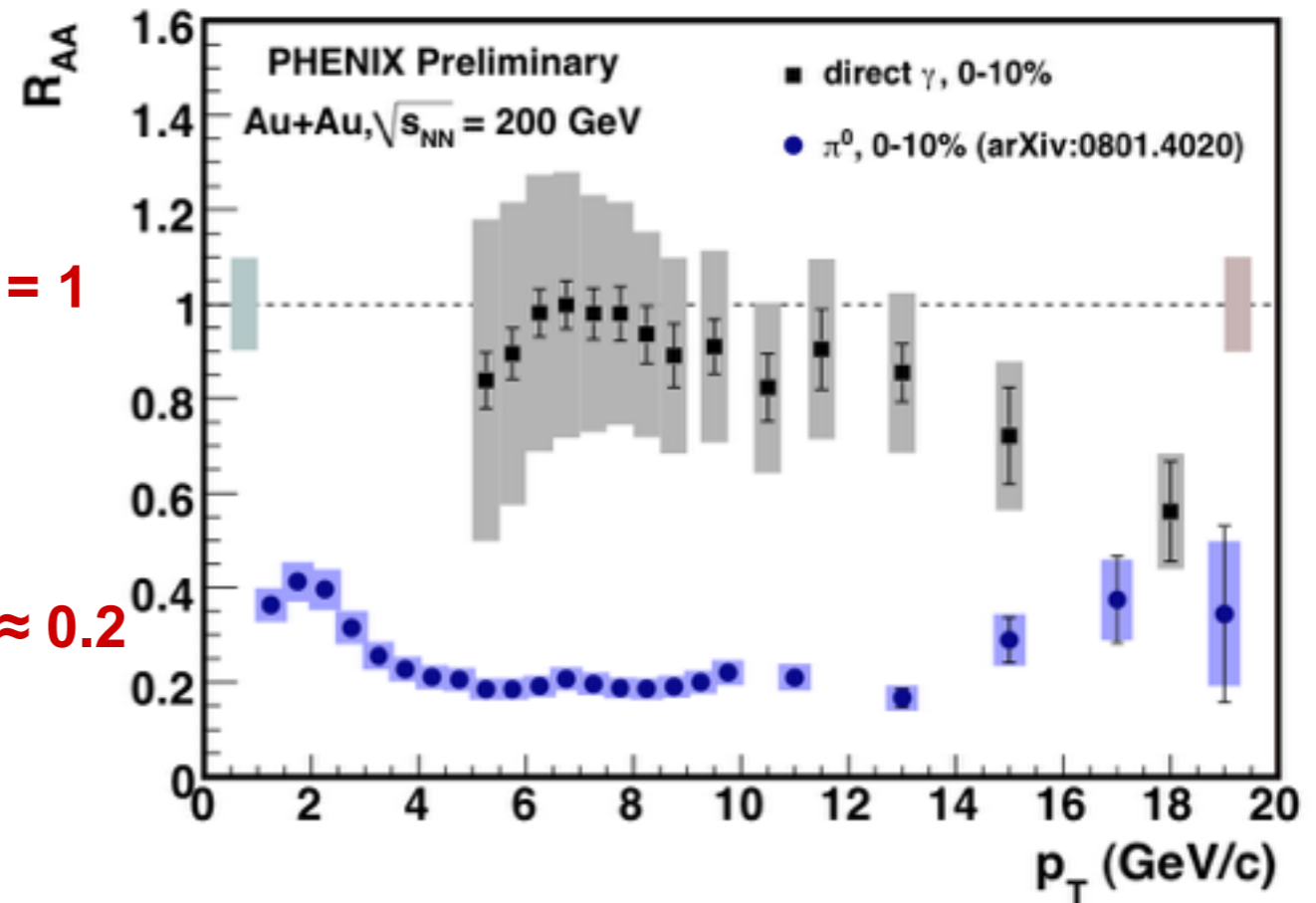
A+A initial production is incoherent superposition of p+p for hard probes

$\pi^0 R_{AA}$ – high- p_T suppression



$\gamma: R_{AA} = 1$

$\pi^0: R_{AA} \approx 0.2$



Hard partons lose energy in the hot matter

Hadrons: energy loss