Twenty years of collaboration... Gerard's 65th birthday celebration

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Twenty years of collaboration...





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Outline

Quadratic eigenvalue problems are no problem Utrecht 1994-1996

Relax to the max Toulouse 2002-2004

The subspaces are shrinking Delft 2005-



Quadratic eigenvalue problems are no problem

UTRECHT



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- At that moment, there was very little research on the numerical solution of nonlinear eigenvalue problems (this has changed completely...).
- ► I asked Gerard if it was possible to use JD for this problem.
- The next day he gave me a manuscript of ten pages ...



JD for nonlinear problem $F(\lambda)\mathbf{x} = \mathbf{0}$

a. Start:

Choose initial subspace ${\bf V}$

- b. Repeat:
 - (1) Compute desired eigenpair (θ, y) from $\mathbf{V}^*F(\theta)\mathbf{V}y = 0$, $\|y\| = 0$
 - (2) Compute $\mathbf{u} \leftarrow \mathbf{V}y$, $\mathbf{r} \leftarrow F(\theta)\mathbf{u}$
 - (3) Stop if satisfied
 - $(4) \ \mathbf{w} \leftarrow F'(\theta) \mathbf{u}$
 - (5) Solve $(\mathbf{I} \frac{\mathbf{u} \, \mathbf{w}^*}{\mathbf{w}^* \mathbf{u}}) F(\theta) (\mathbf{I} \mathbf{u} \, \mathbf{u}^*) \mathbf{t} = -\mathbf{r}$
 - (6) Expand $\mathbf{V}: \mathbf{V} \leftarrow \mathsf{Ortho}([\mathbf{V} | \mathbf{t}]).$



Application: room acoustics

Propagation of sound:

$$\frac{1}{c^2}\frac{\partial^2 p}{\partial t^2} = \bigtriangleup p \ in \ \Omega$$

 $c\!\!:$ speed of sound, $p\!\!:$ acoustic pressure

Boundary conditions:

- Reflecting wall: $\frac{\partial p}{\partial n}=0$
- Open wall (pressure release): p=0
- Absorbing wall: $\frac{\partial p}{\partial n} = -\frac{1}{z}\frac{\partial p}{\partial t}$
- \boldsymbol{z} is the (complex) impedance, may depend on frequency



Analytical eigenproblem

Assume solution

$$p = \bar{p}e^{\lambda t}$$

with λ an analytical eigenvalue. Substitute in wave equation:

$$\frac{\lambda^2}{c^2}\bar{p} = \triangle \bar{p}$$

B.c. for absorbing wall:

$$\frac{\partial \bar{p}}{\partial n} = -\frac{\lambda}{z(\lambda)}\bar{p}$$

Eigenfrequencies: $f = Im(\lambda)/2\pi$



Discretization with FEM

Discretisation with the Finite Element method gives:

$$\mathbf{K}\mathbf{p} + \lambda \mathbf{C}(\lambda)\mathbf{p} + \lambda^2 \mathbf{M}\mathbf{p} = \mathbf{0}$$

- p: Discretization of \bar{p}
- λ : Algebraic eigenvalue
- K: Stiffness matrix, discretization of $-\triangle$
- ► C: Damping matrix, discretization of impedance condition
- M: Mass matrix, discretization of $\frac{1}{c^2}$



Example cube

- Domain: Ω is $4\times 4\times 4~m^3$
- 5 reflecting walls
- 1 absorbing wall
- Soundspeed c=340m/s
- Impedance: z = (0.2 1.5i)c(does not depend on frequency \rightarrow quadratic eigenproblem)
- Discretisation: linear tetrahedral elements
- Known analytical eigenvalue: $\lambda = -5.19 + 217.5i$



Results cube

- Grid: $64 \times 64 \times 64 \times 5$ elements \rightarrow 274,625 grid points
- Goal: eigenvalue closest to -5.19 + 217.5i
- Selection of Ritz value closest to analytical eigenvalue
- Correction equation 'solved' with 30 GMRES steps
- Restart after 20 JD-iterations to save memory

Calculations performed on Cray T3D using 64 processors. Eigenvalue -5.20 + 217.5i converged after 33 JD-iterations.

Computing time: 93.4 seconds (20 years ago!!!)



Relax to the max

TOULOUSE



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Relaxation





Relaxation strategies

Krylov methods allow the matrix-vector products to be computed to lower accuracy when the process starts to converge.

This was a hot topic at Cerfacs and Jasper van den Eshof and Gerard were also working on it in Utrecht.



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A simple relaxation strategy is to choose

$$\eta_j = \frac{\epsilon}{\|\mathbf{r}_j\|}$$

No relaxation



Relaxation



Idea: use nested Krylov method

- Outer loop
 - Aim for desired accuracy
 - Use relaxation
- Inner loop (preconditioner)
 - Aim for low accuracy
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Idea: use nested Krylov method

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Almost all matrix-vector products can be performed at low accuracy!



Relax to the max



Example: ocean circulation



Discrete ocean problem

Global system:

$$\left(\begin{array}{cc} r\mathbf{L} - \mathbf{C} & A\bar{\mathbf{L}} \\ -\bar{\mathbf{L}}^H & \mathbf{M} \end{array}\right) \left(\begin{array}{c} \bar{\psi} \\ \bar{\zeta} \end{array}\right) = \left(\begin{array}{c} \mathbf{f} \\ \mathbf{0} \end{array}\right)$$

Schur complement for the vorticity:

$$(\mathbf{M} + A\bar{\mathbf{L}}^H(r\mathbf{L} - \mathbf{C})^{-1}\bar{\mathbf{L}})\bar{\boldsymbol{\zeta}} = \bar{\mathbf{L}}^H(r\mathbf{L} - \mathbf{C})^{-1}\mathbf{f}$$

- ► Action of (rL C)⁻¹ is approximated by preconditioned Bi-CGSTAB.
- Accuracy of Bi-CGSTAB, and therefore of the multiplication with the Schur complement, can be controlled.



Relax to the max



TUDelft

The subspaces are shrinking

DELFT



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The subspaces are shrinking (using IDR)





Two theorems

Theorem (IDR)

Let $\mathbf{P} = [\mathbf{p}_1, \dots, \mathbf{p}_s]$ be an $n \times s$ matrix and let (μ_j) be a sequence in \mathbb{C} . With $\mathcal{G}_0 \equiv \mathcal{C}^n$, define,

$$\mathcal{G}_{j+1} \equiv (\mu_j \mathbf{I} - \mathbf{A})(\mathcal{G}_j \cap \mathbf{P}^{\perp}) \qquad (j = 0, 1, \ldots).$$

If \mathbf{P}^{\perp} does not contain an eigenvector of \mathbf{A} , then, for all $j = 0, 1, \ldots$, we have that 1) $\mathcal{G}_{j+1} \subset \mathcal{G}_j$, and 2) dim $\mathcal{G}_{j+1} < \dim \mathcal{G}_j$ unless $\mathcal{G}_j = \mathbf{0}$.



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Theorem (BLOCK-KRYLOV)

Let \mathbf{P} , (μ_j) and \mathcal{G}_j be as above. Consider the polynomial p of degree j given by $p(\lambda) \equiv \prod_{i=1}^{j} (\mu_i - \lambda) \quad (\lambda \in \mathbb{C})$. Then

$$\mathcal{G}_j = \{ p_j(\mathbf{A}) \mathbf{v} \mid \mathbf{v} \perp \mathcal{K}_j(\mathbf{A}^H, \mathbf{P}) \}.$$



IDRStab

Can other polynomials than products of linear factors

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IDRStab

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be used (cf. BiCGSTAB versus BiCGstab(ℓ))? Yes, by combining several dimension reductions steps it is possible to construct higher order factors.

The name of this method is IDRStab. (SISC 2010)



This is the moment Gerard figured it out





Example

Convection-diffusion problem with shift on unit square:

$$-u_{xx} - u_{yy} + 1000(xu_x + yu_y) + 10u = f .$$

- Dirichlet conditions
- ► Discretised with the finite volume method on a 65 × 65 grid
- The right-hand side vector such that the solution is one.



Convergence of IDRStab





Numerical Linear Algebra course

 For several years we have taught together the Mastermath course on Numerical Linear Algebra



Gerard can teach Numerical Linear Algebra to any audience





I do my best ...





To conclude: nice memories of Japan





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Gerard can find the way in any language





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of collaboration



- of collaboration
- but also of friendship and fun



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Gerard: CONGRATULATIONS!

